

A visualization of a particle collision event simulation. It shows a central point from which numerous lines radiate outwards. The lines are primarily red, with some green and yellow lines interspersed. The lines are thicker near the center and become thinner as they extend outwards, suggesting a branching or showering process. The overall shape is roughly circular, with the lines filling the space around the center.

# Event Simulation at the Large Hadron Collider

Bryan Webber  
Cavendish Laboratory  
University of Cambridge

# Event Simulation at the Large Hadron Collider

- Monte Carlo event generation:
  - ✧ theoretical status and limitations
- Recent improvements:
  - ✧ perturbative and non-perturbative
- Overview of results:
  - ✧ W, Z, top, Higgs, BSM (+jets)

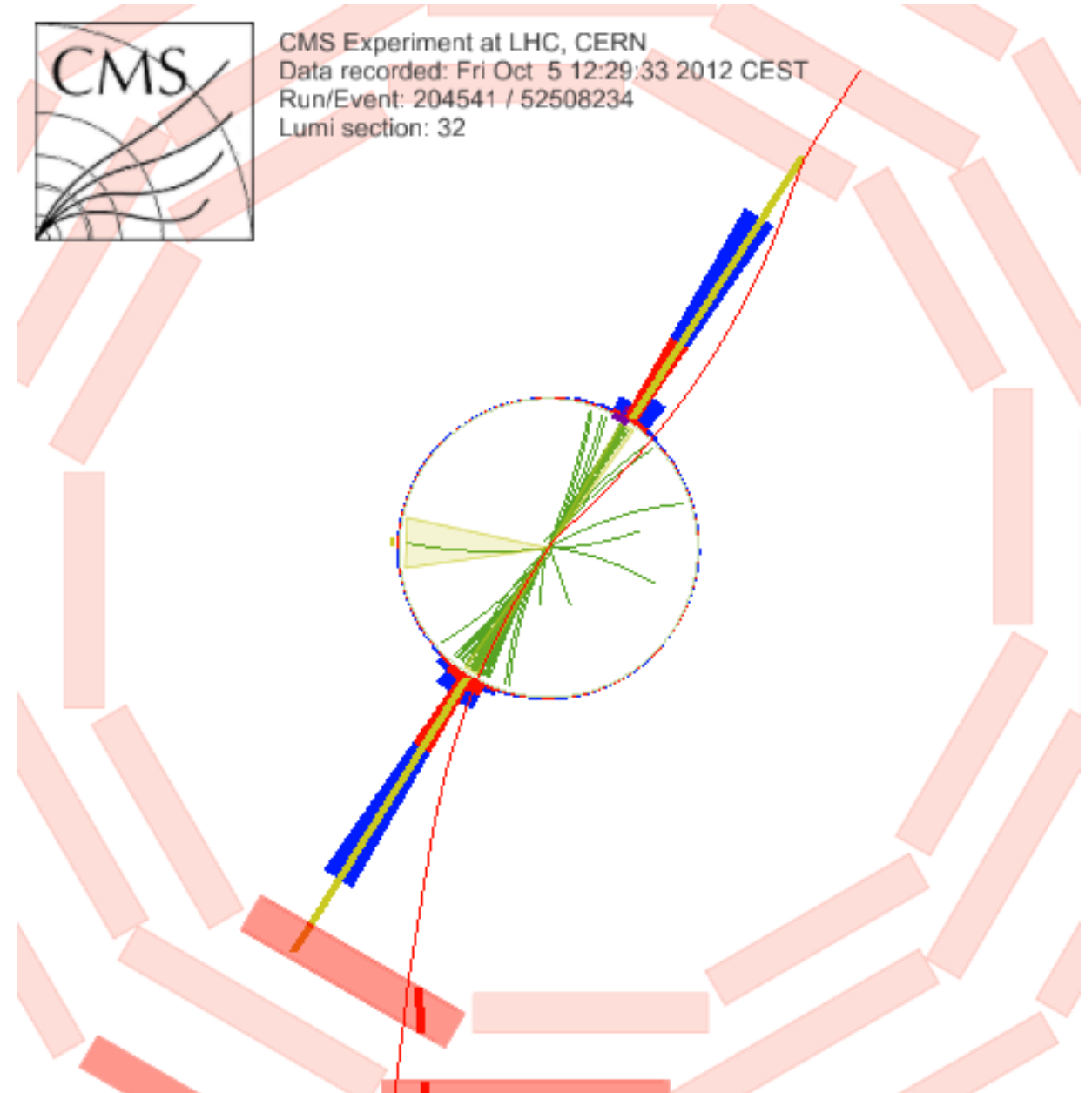
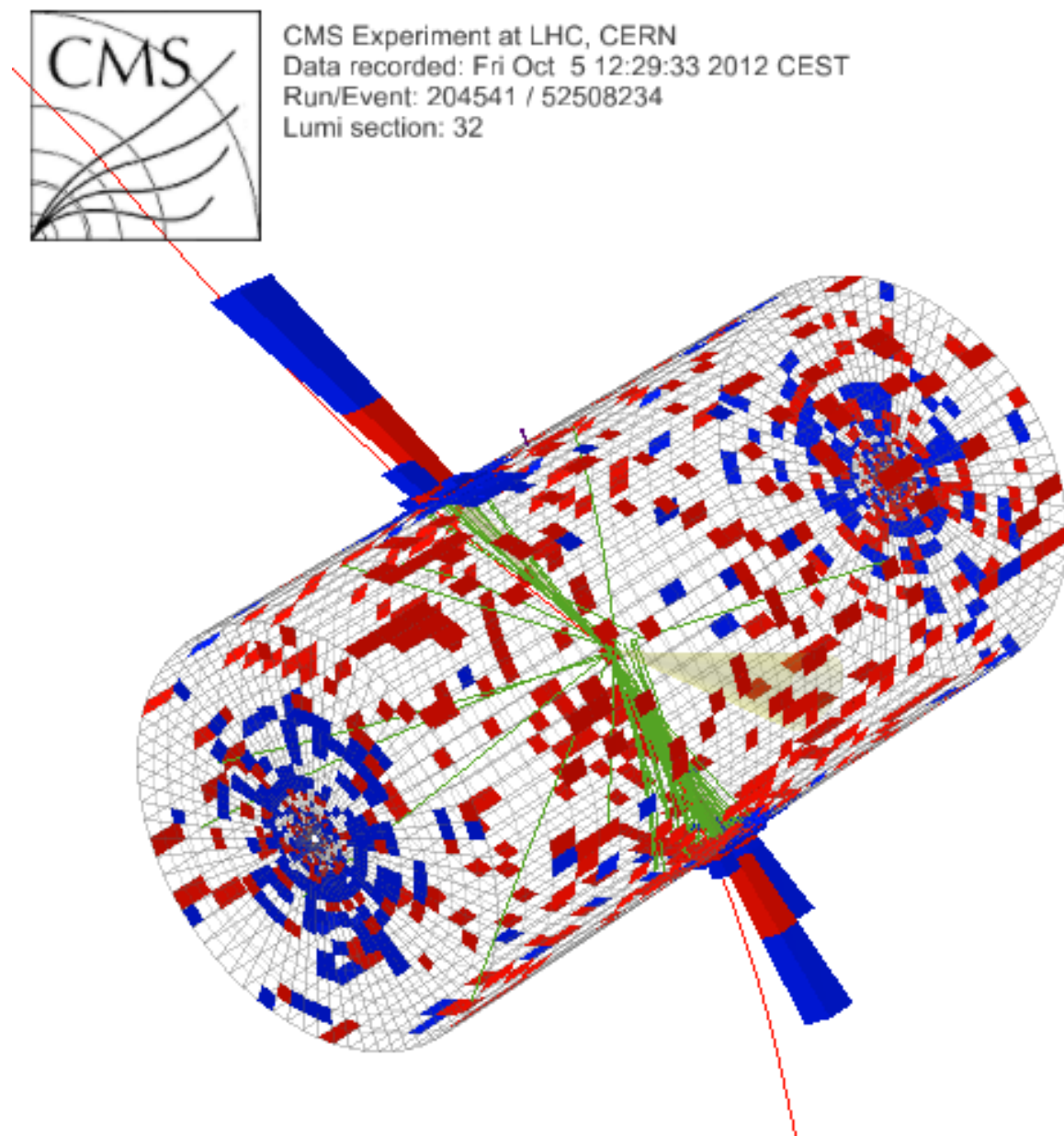
# Monte Carlo Event Generation

# Monte Carlo Event Generation

- Aim is to produce simulated (particle-level) datasets like those from real collider events
  - ✧ i.e. lists of particle identities, momenta, ...
  - ✧ simulate quantum effects by (pseudo)random numbers
- Essential for:
  - ✧ Designing new experiments and data analyses
  - ✧ Correcting for detector and selection effects
  - ✧ Testing the SM and measuring its parameters
  - ✧ Estimating new signals and their backgrounds



# A high-mass dijet event

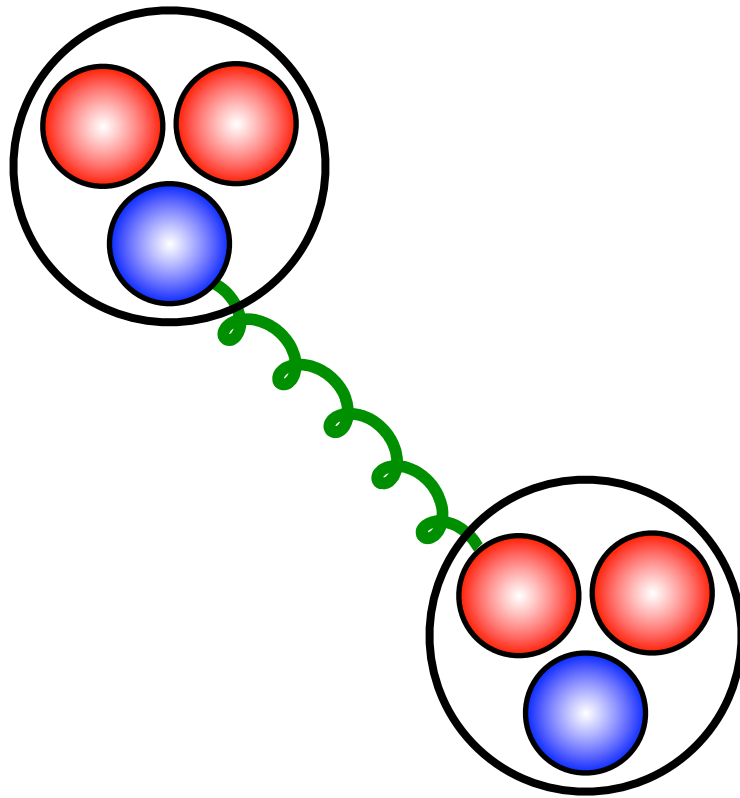


●  $M_{jj} = 5.15 \text{ TeV}$

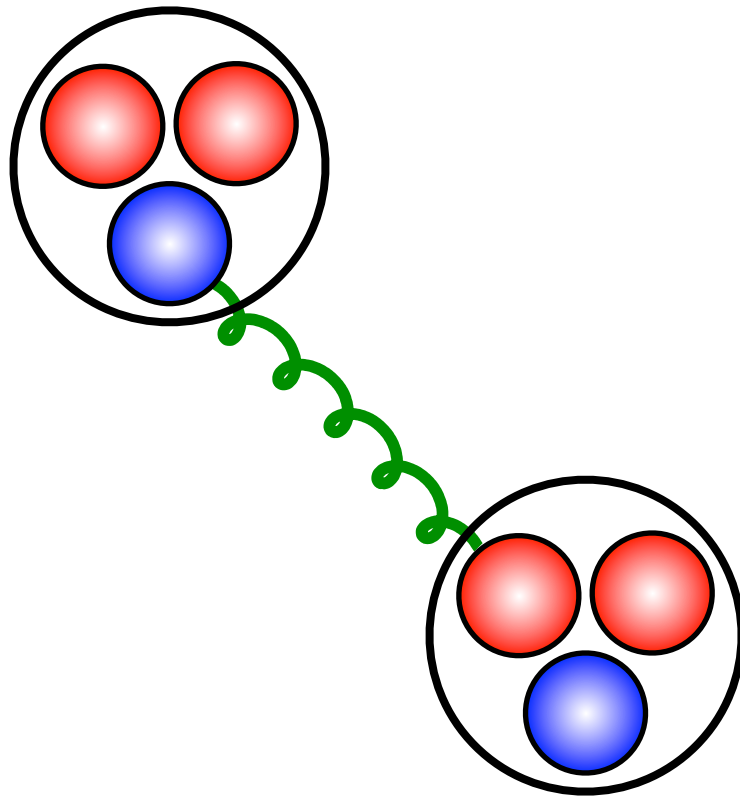
CMS PAS EXO-12-059

# LHC Dijet

# LHC Dijet

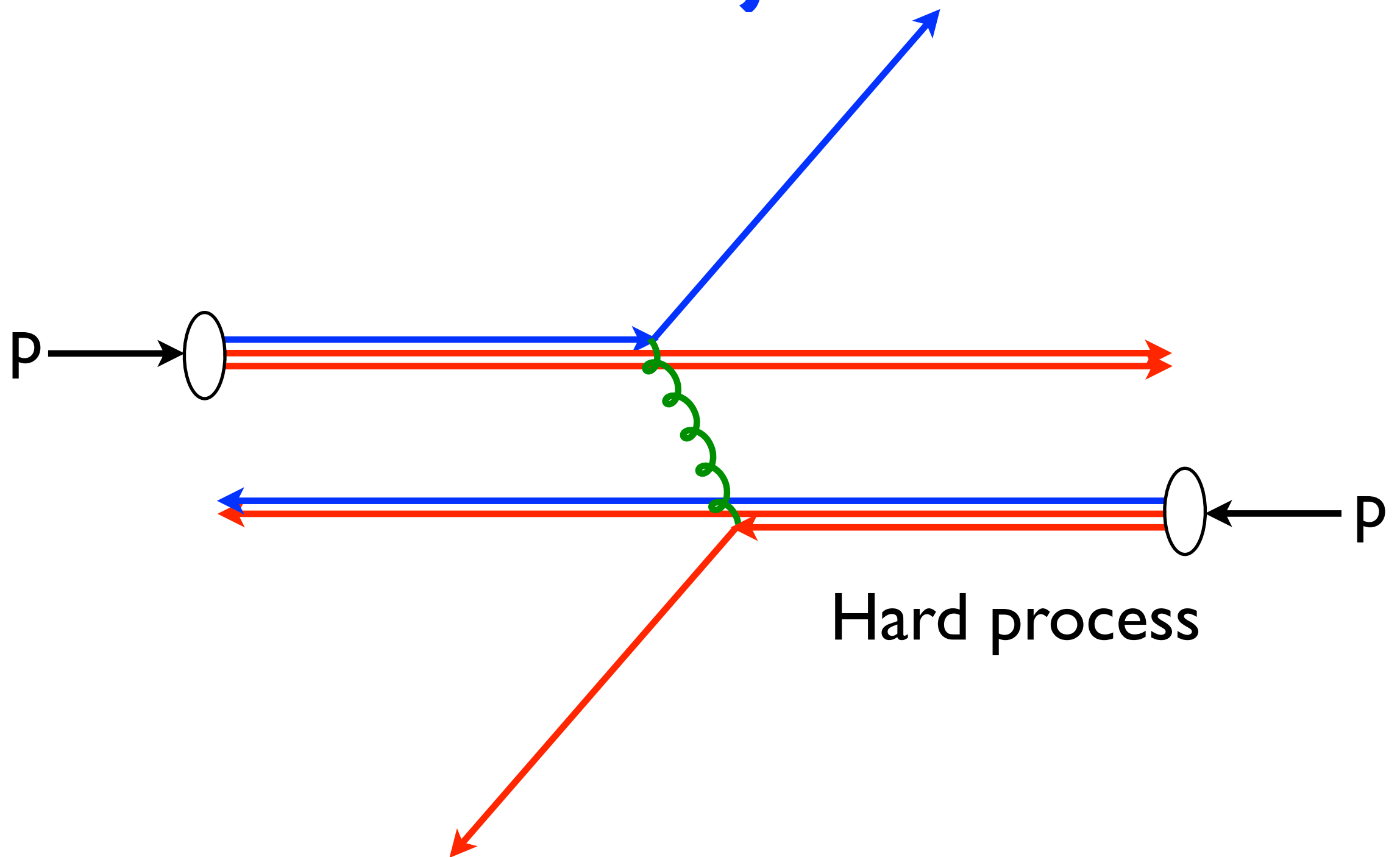


# LHC Dijet

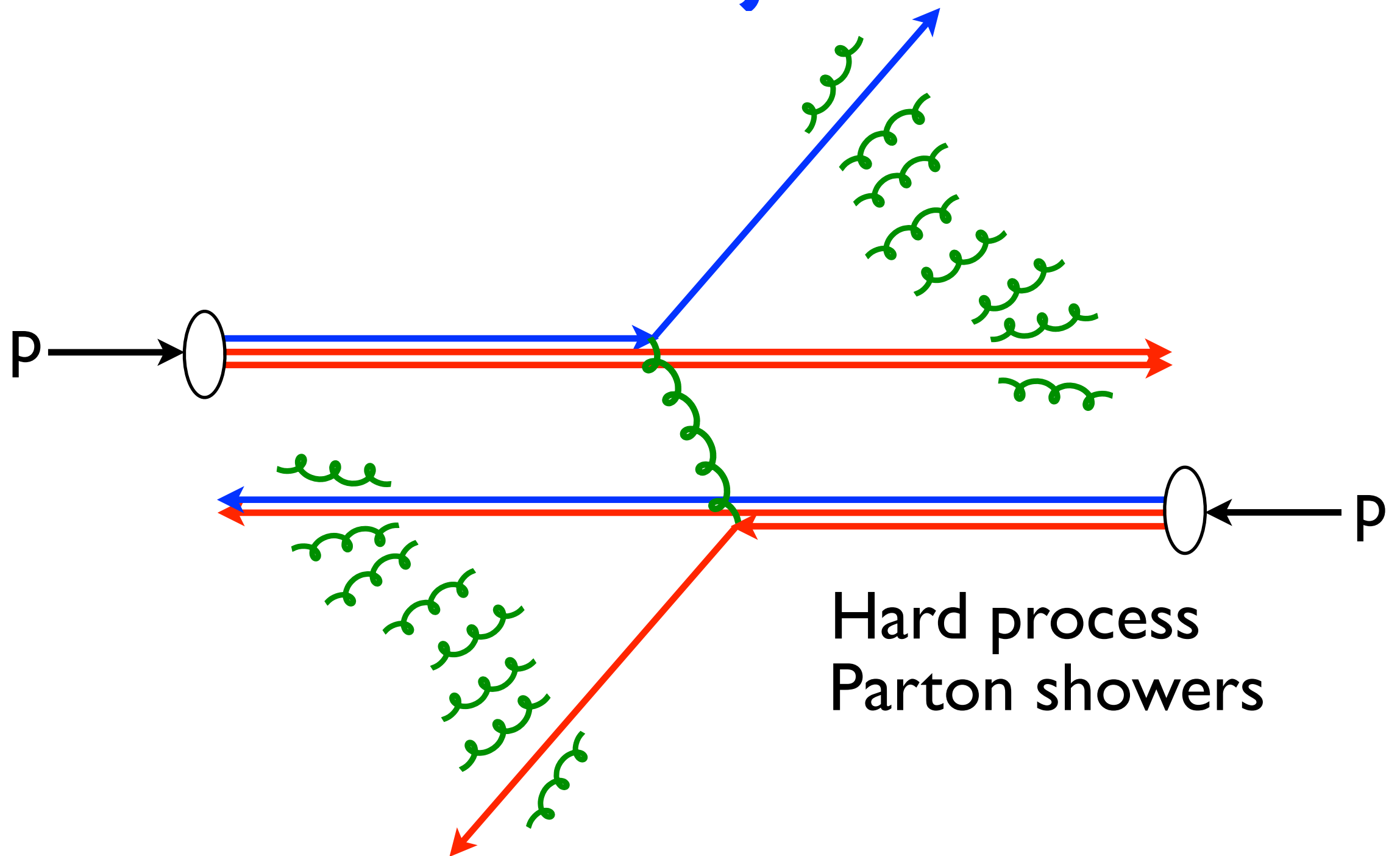


# LHC Dijet

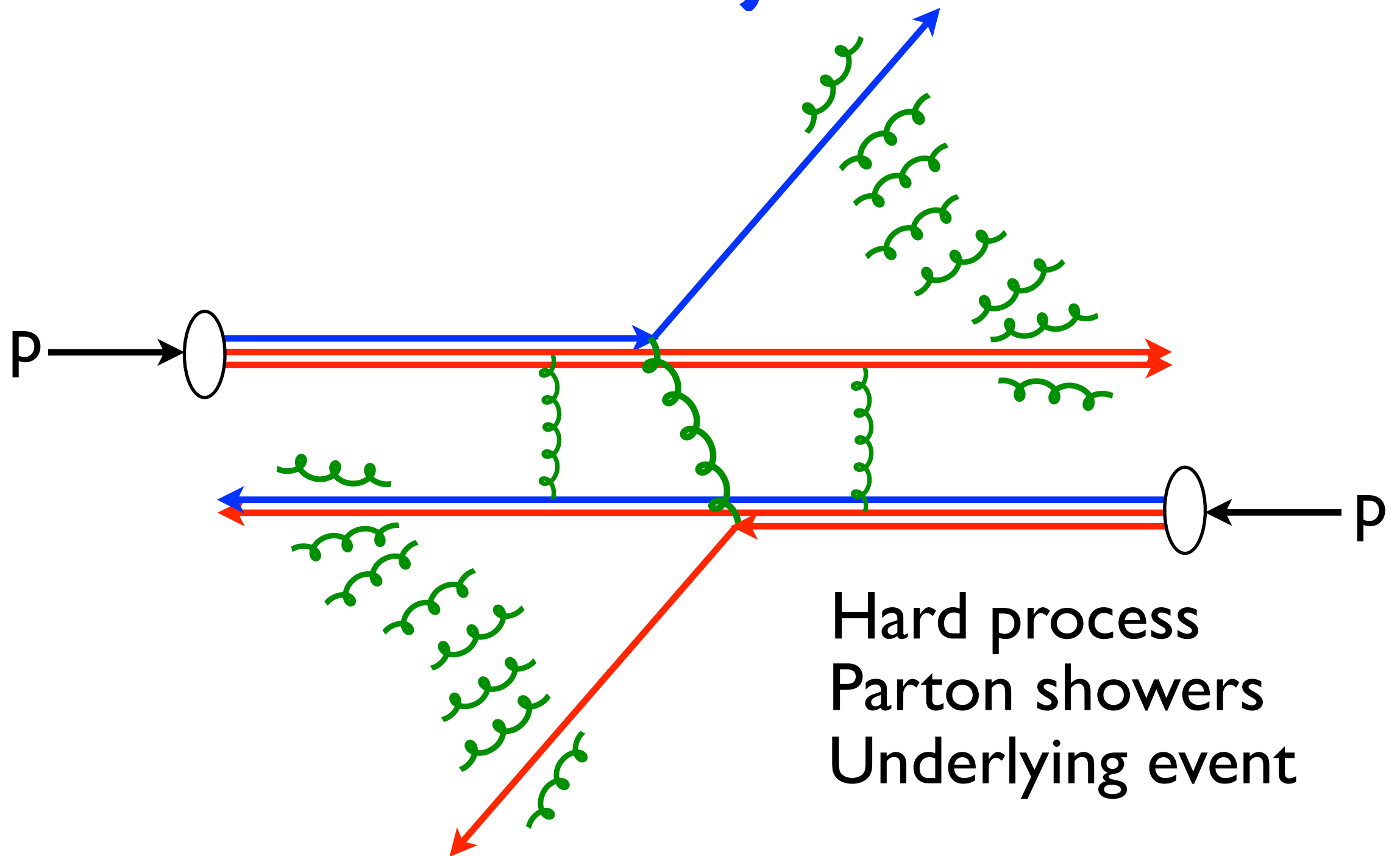
# LHC Dijet



# LHC Dijet

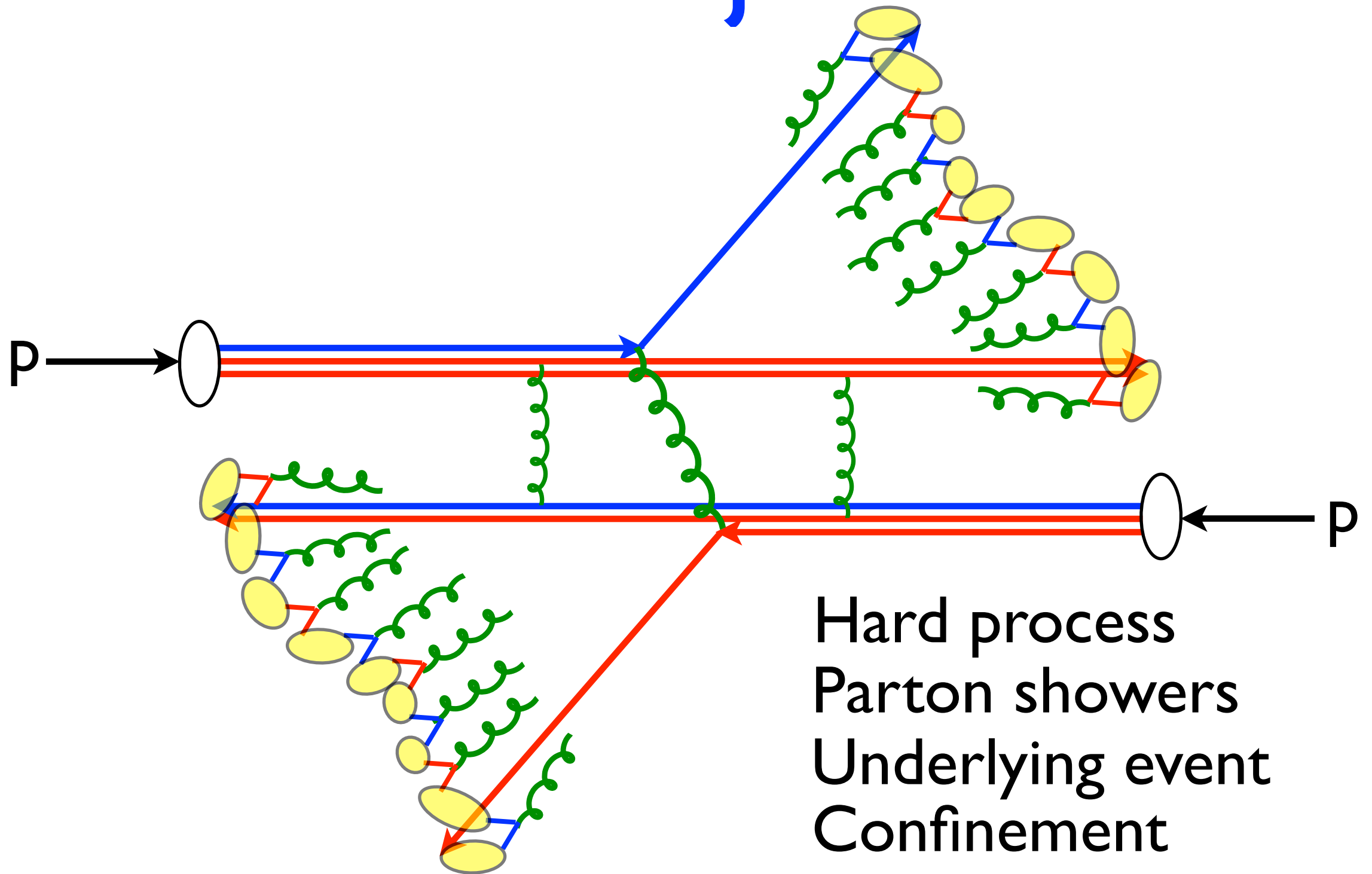


# LHC Dijet

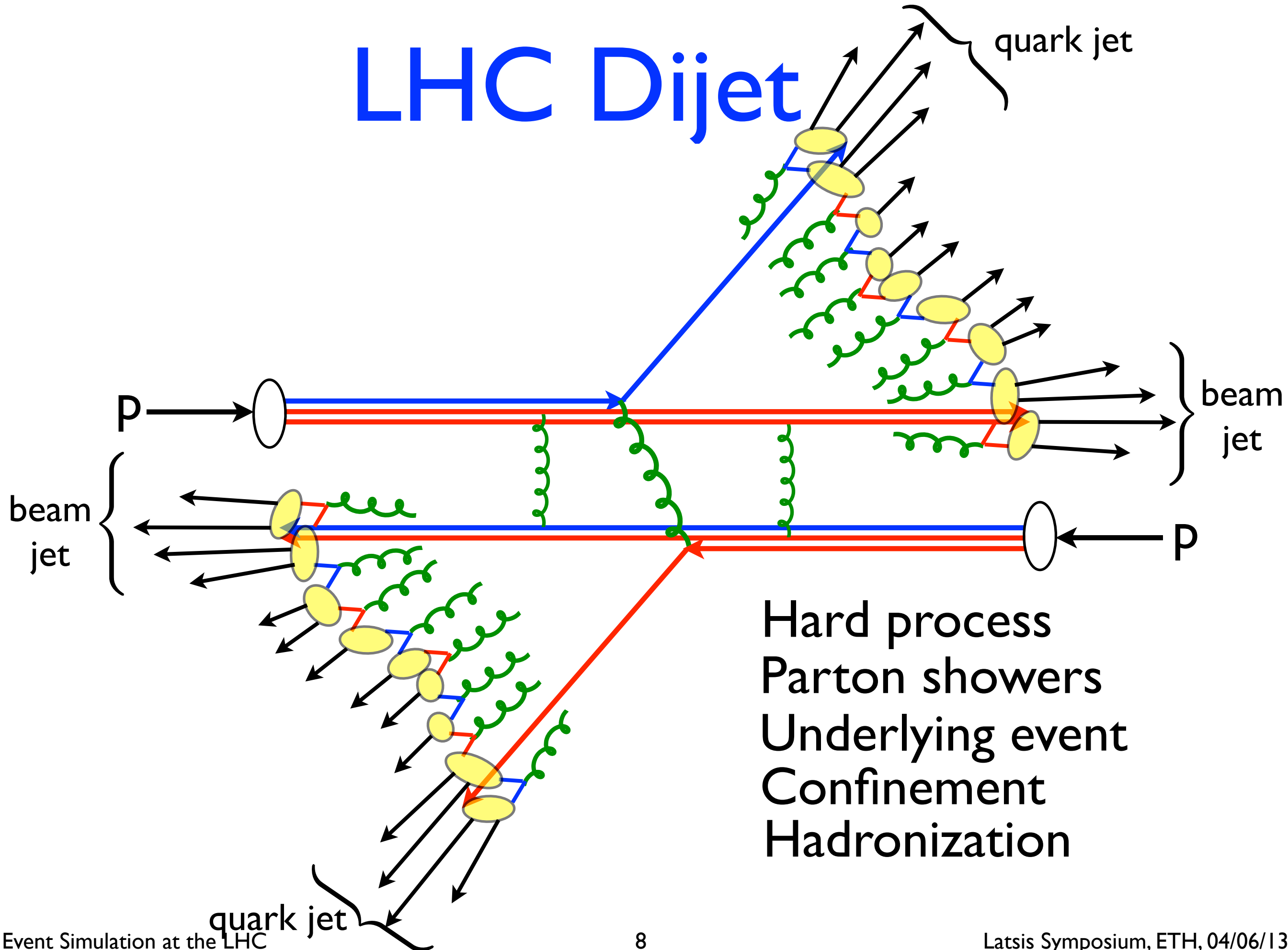




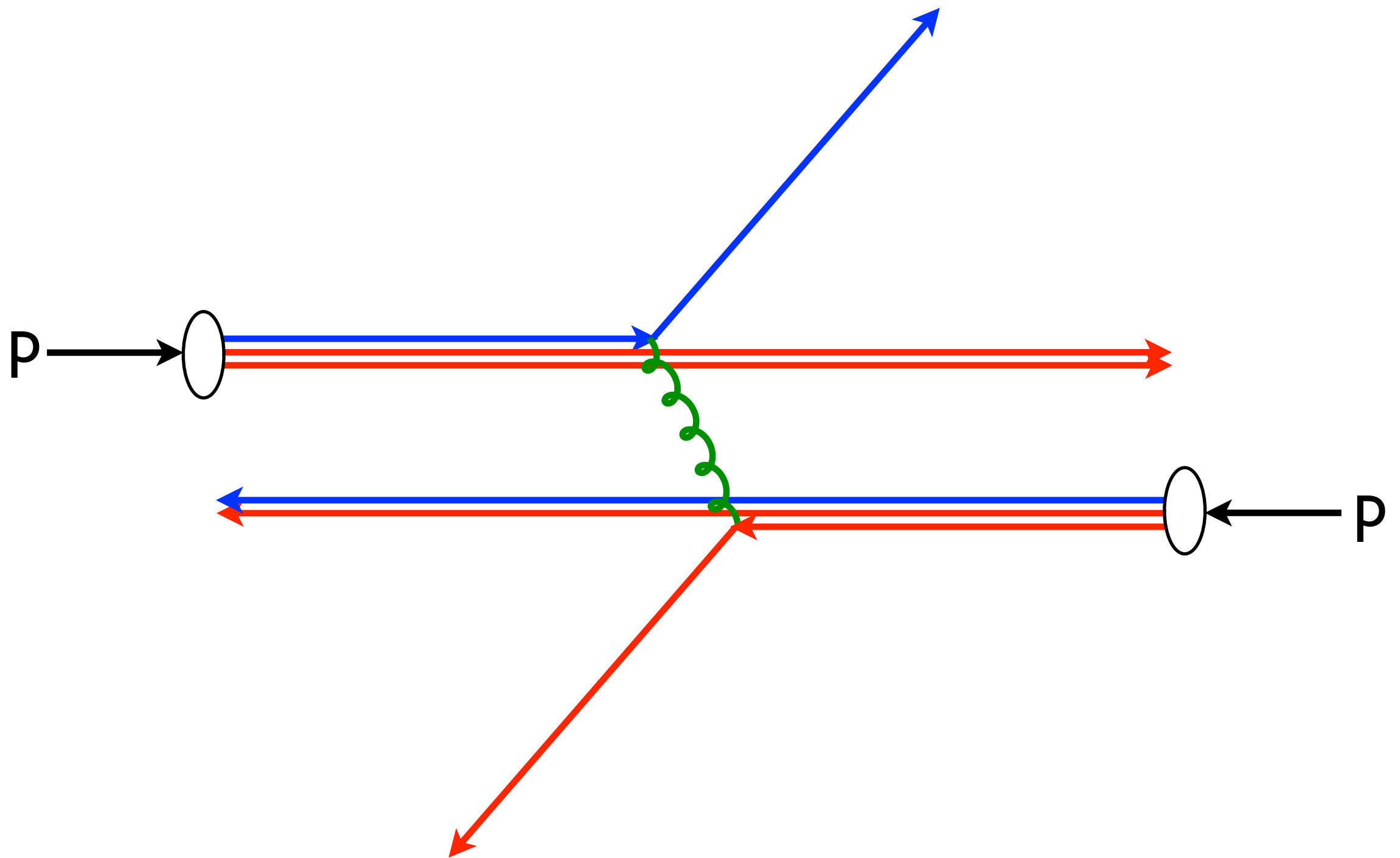
# LHC Dijet



# LHC Dijet

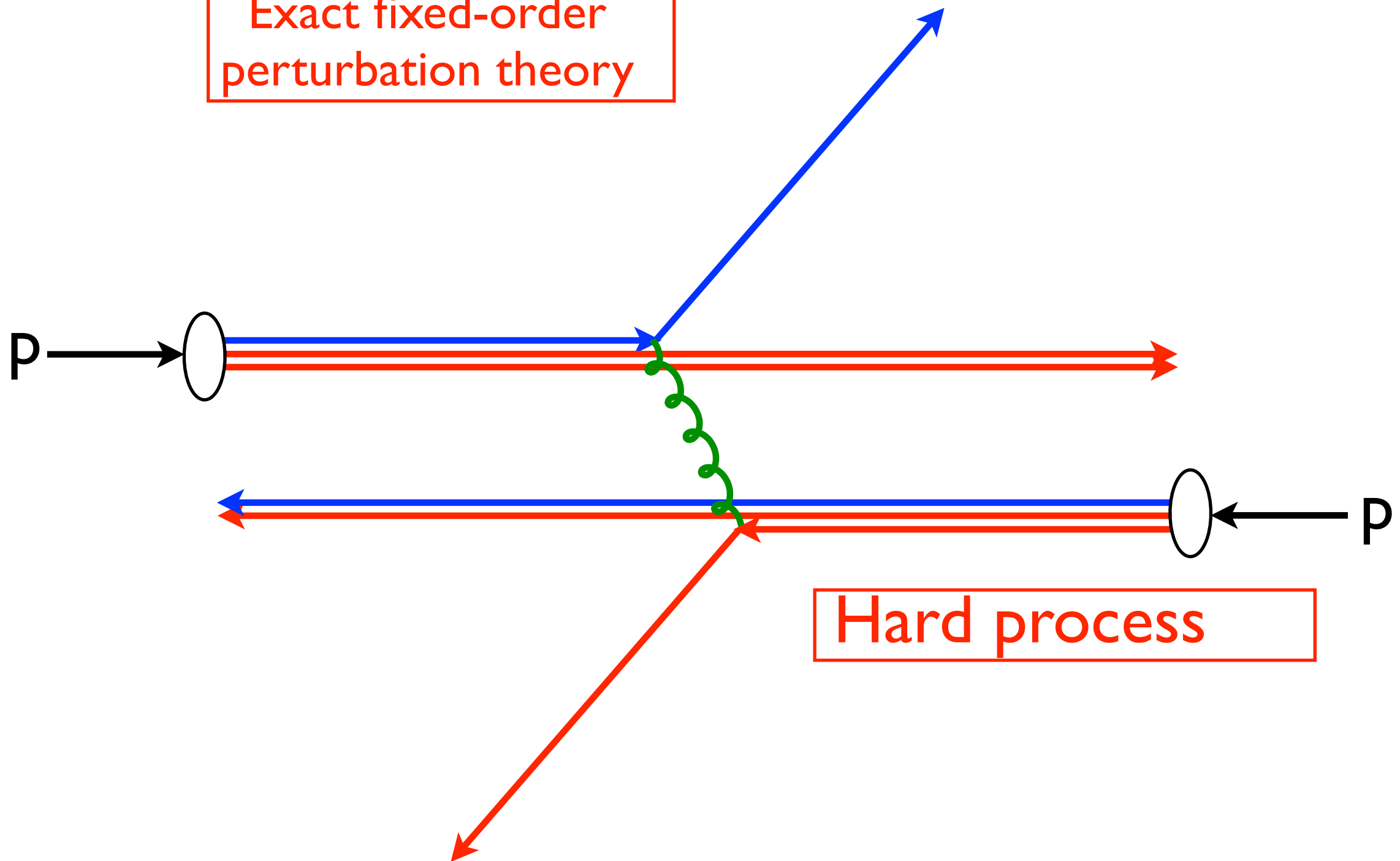


# Theoretical status

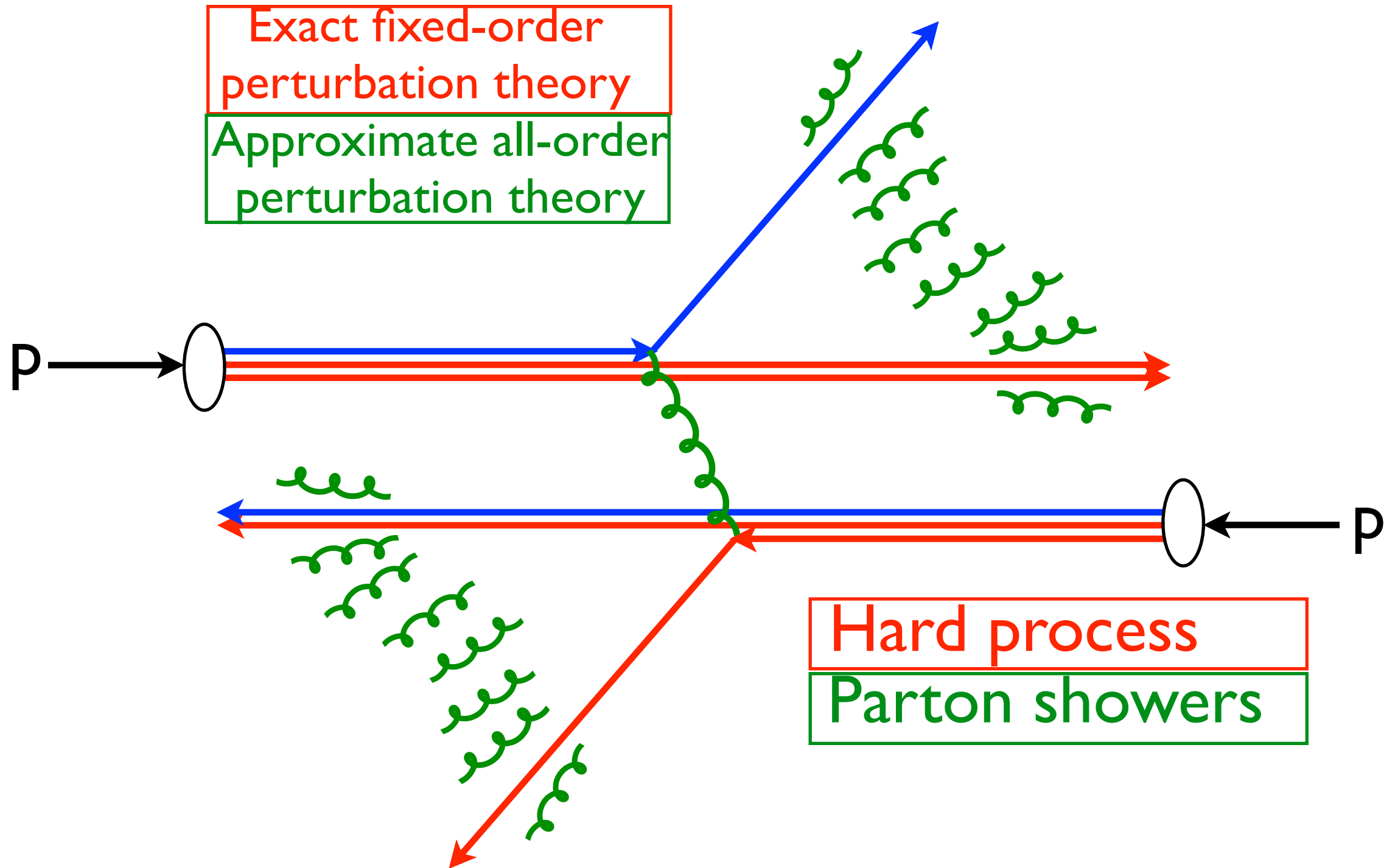


# Theoretical status

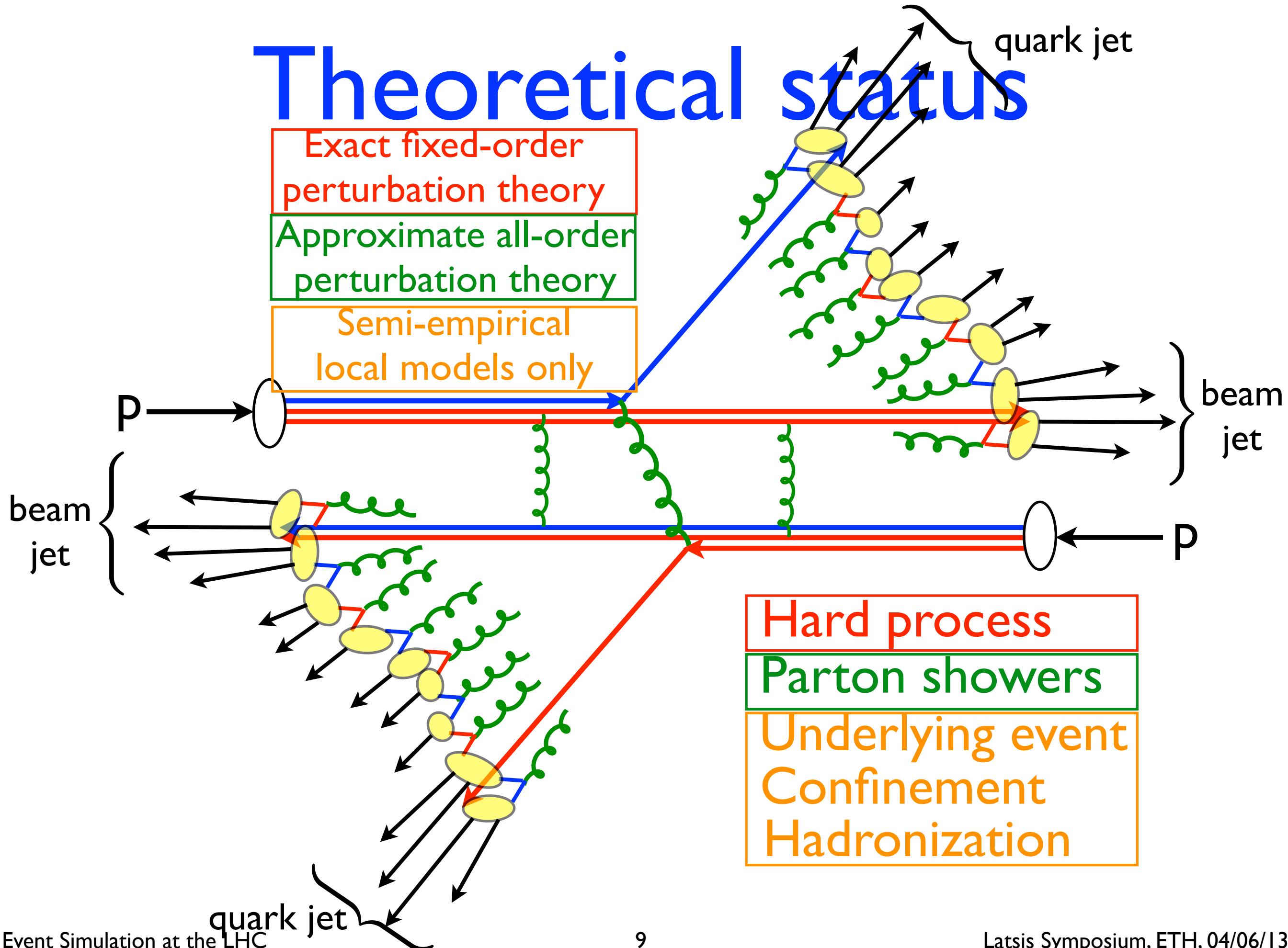
Exact fixed-order  
perturbation theory



# Theoretical status



# Theoretical status



# QCD Factorization

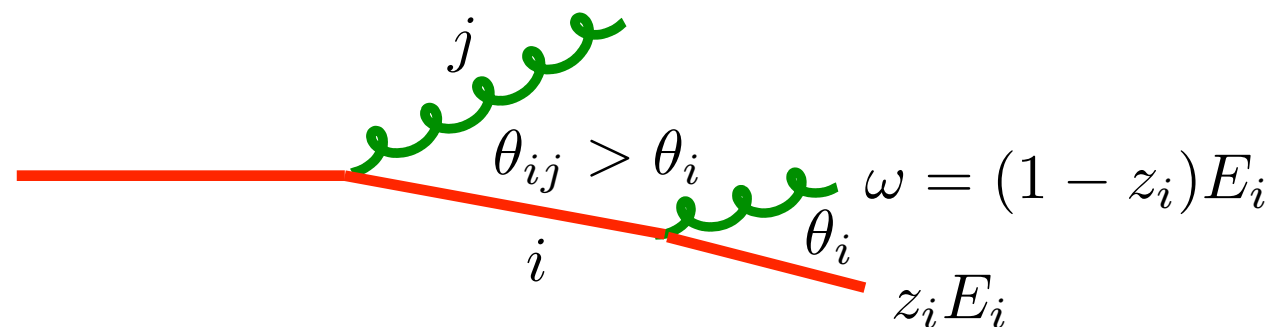
$$\sigma_{pp \rightarrow X}(E_{pp}^2) = \int_0^1 dx_1 dx_2 \underbrace{f_i(x_1, \mu^2) f_j(x_2, \mu^2)}_{\substack{\text{parton} \\ \text{distributions} \\ \text{at scale } \mu^2}} \underbrace{\hat{\sigma}_{ij \rightarrow X}(x_1 x_2 E_{pp}^2, \mu^2)}_{\substack{\text{hard process} \\ \text{cross section}}}$$

momentum fractions

- Jet formation and underlying event take place over a much longer time scale, with unit probability
- Hence they cannot affect the cross section
- Scale dependences of parton distributions and hard process cross section are perturbatively calculable, and cancel order by order

# Parton Shower Approximation

- Keep only most singular parts of QCD matrix elements:
- **Collinear**  $d\sigma_{n+1} \approx \frac{\alpha_S}{2\pi} \sum_i P_{ii}(z_i, \phi_i) dz_i \frac{d\xi_i}{\xi_i} \frac{d\phi_i}{2\pi} d\sigma_n$   $\xi_i = 1 - \cos \theta_i$
- **Soft**  $d\sigma_{n+1} \approx \frac{\alpha_S}{2\pi} \sum_{i,j} (-\mathbf{T}_i \cdot \mathbf{T}_j) \frac{p_i \cdot p_j}{p_i \cdot k p_j \cdot k} \omega d\omega d\xi_i \frac{d\phi_i}{2\pi} d\sigma_n$   
 $= \frac{\alpha_S}{2\pi} \sum_{i,j} (-\mathbf{T}_i \cdot \mathbf{T}_j) \frac{\xi_{ij}}{\xi_i \xi_j} \frac{d\omega}{\omega} d\xi_i \frac{d\phi_i}{2\pi} d\sigma_n$   
 $\approx \frac{\alpha_S}{2\pi} \sum_{i,j} (-\mathbf{T}_i \cdot \mathbf{T}_j) \Theta(\xi_{ij} - \xi_i) \frac{d\omega}{\omega} \frac{d\xi_i}{\xi_i} d\sigma_n$

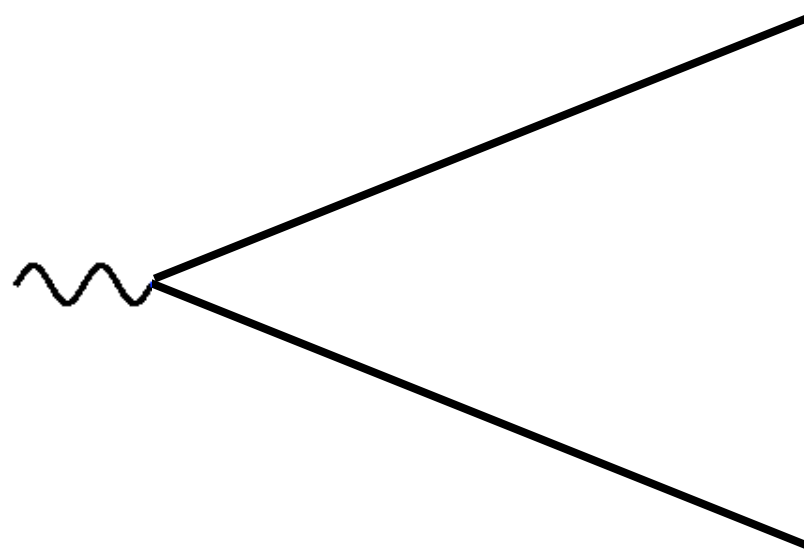


➔ Angular-ordered **parton shower** (or **dipoles**)



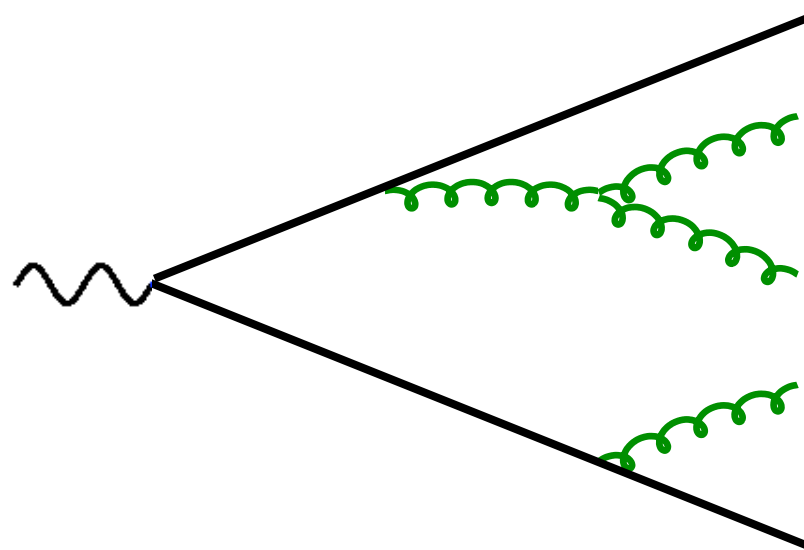
# Hadronization Models

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Colour, flavour and momentum flows are only **locally** redistributed by hadronization



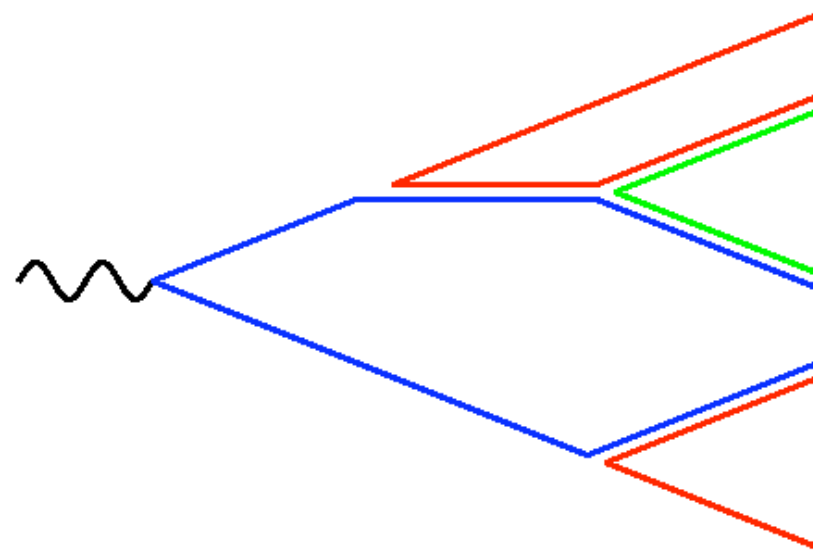
# Hadronization Models

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Colour, flavour and momentum flows are only **locally** redistributed by hadronization



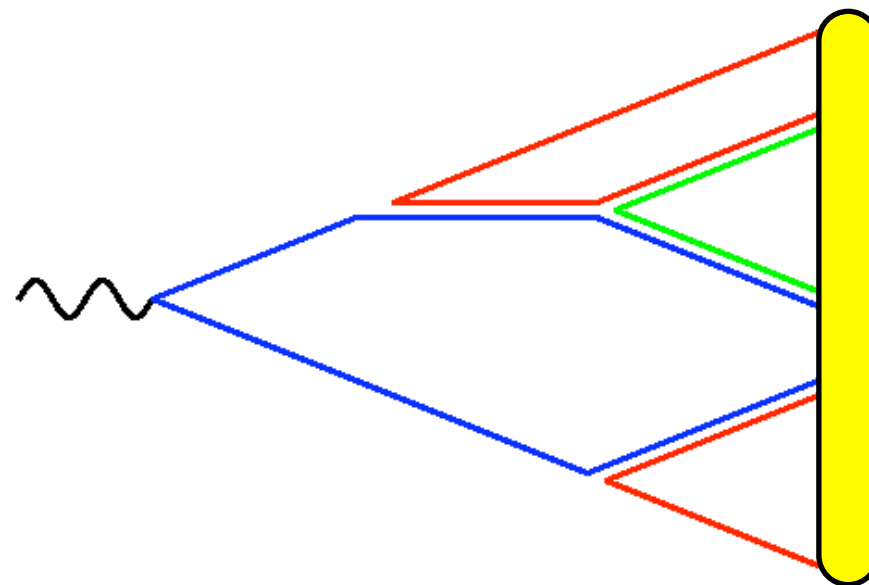
# Hadronization Models

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Colour, flavour and momentum flows are only **locally** redistributed by hadronization



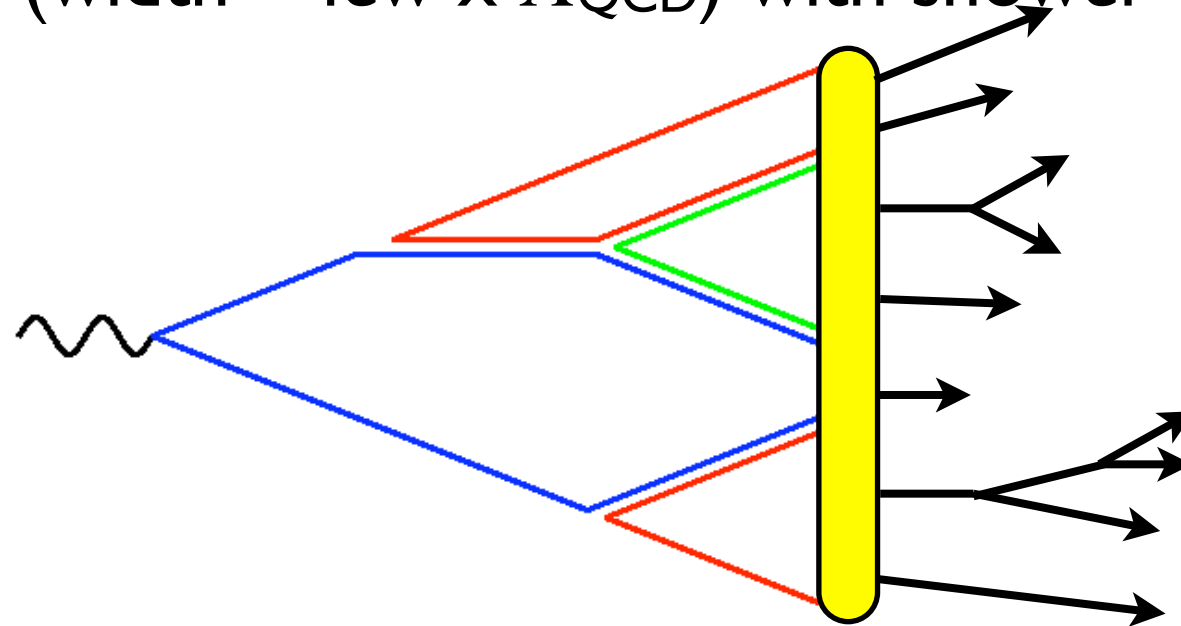
# String Hadronization Model

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Colour flow dictates how to connect hadronic string (width  $\sim \text{few} \times \Lambda_{\text{QCD}}$ ) with shower



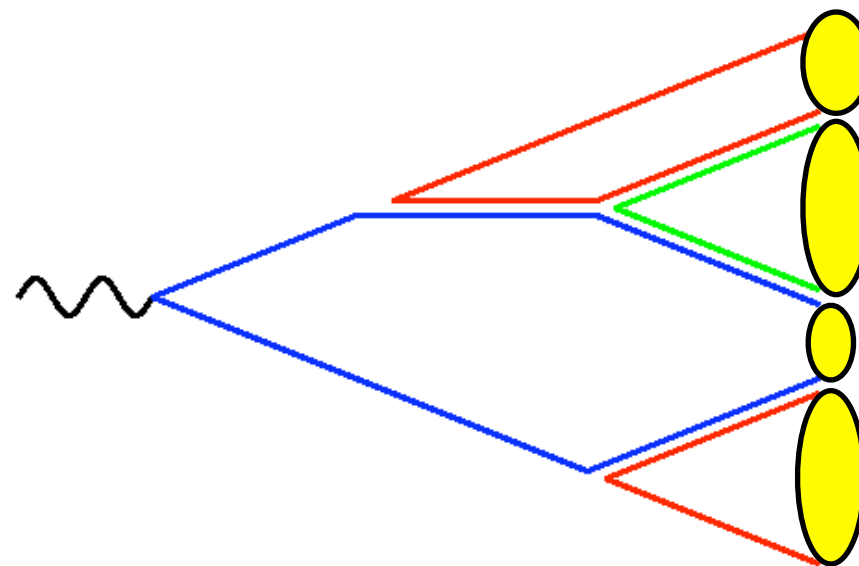
# String Hadronization Model

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Colour flow dictates how to connect hadronic string (width  $\sim \text{few} \times \Lambda_{\text{QCD}}$ ) with shower



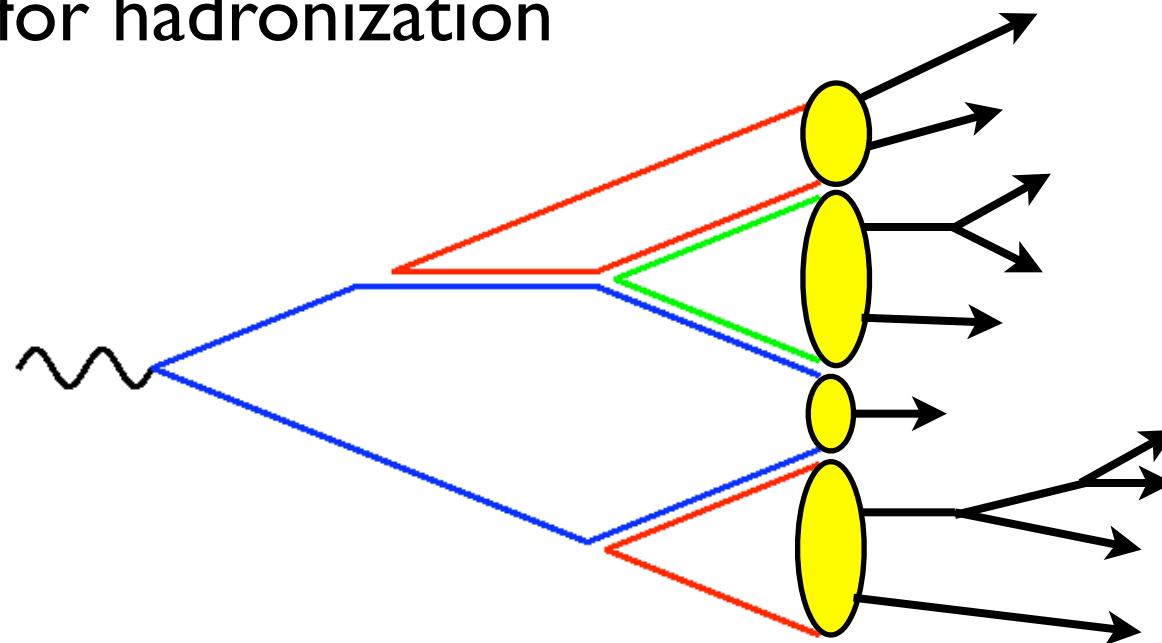
# Cluster Hadronization Model

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Decay of preconfinement clusters provides a direct basis for hadronization

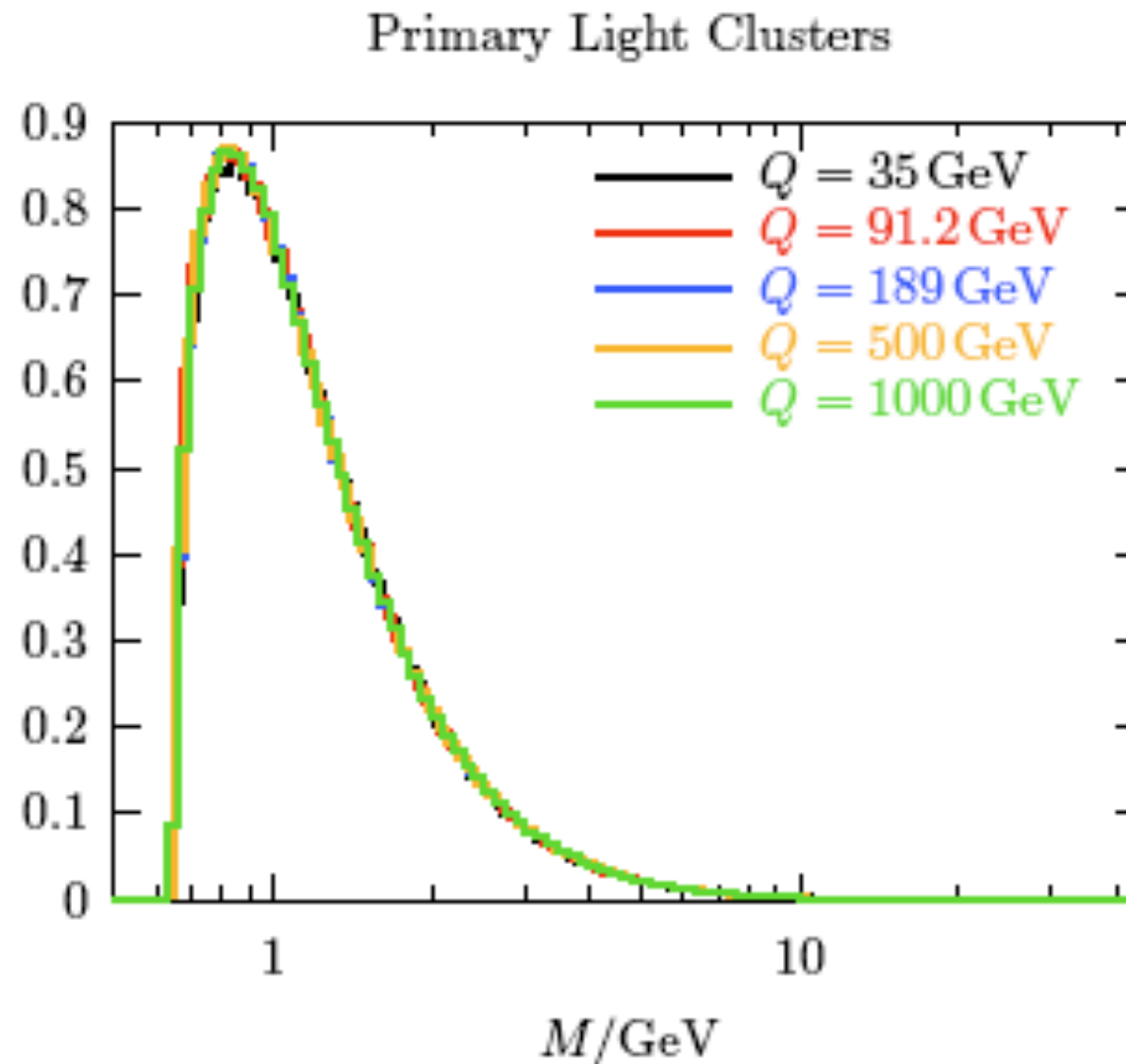


# Cluster Hadronization Model

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $Q_0 \sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Decay of preconfinement clusters provides a direct basis for hadronization



# Cluster Hadronization Model



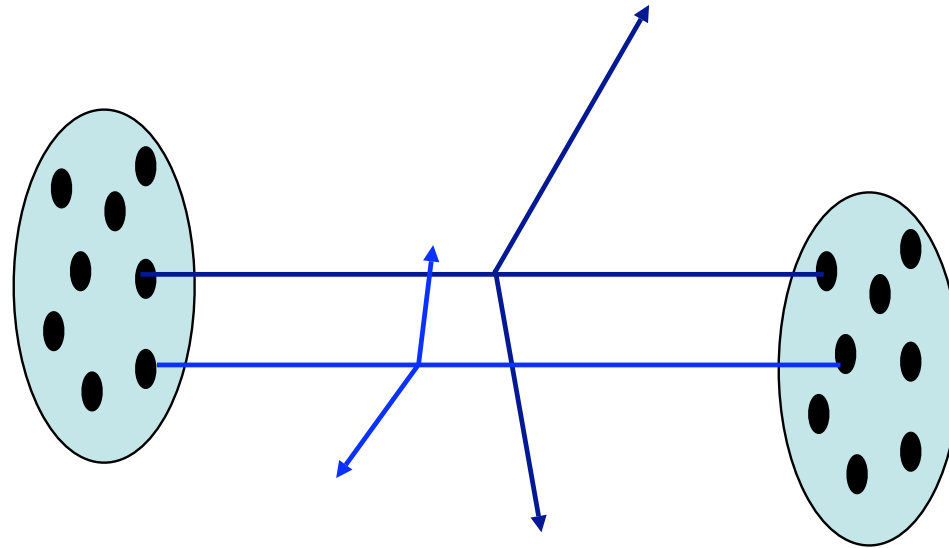
- Mass distribution of preconfined clusters is universal
- Phase-space decay model for most clusters
- High-mass tail decays anisotropically (string-like)



# Hadronization Status

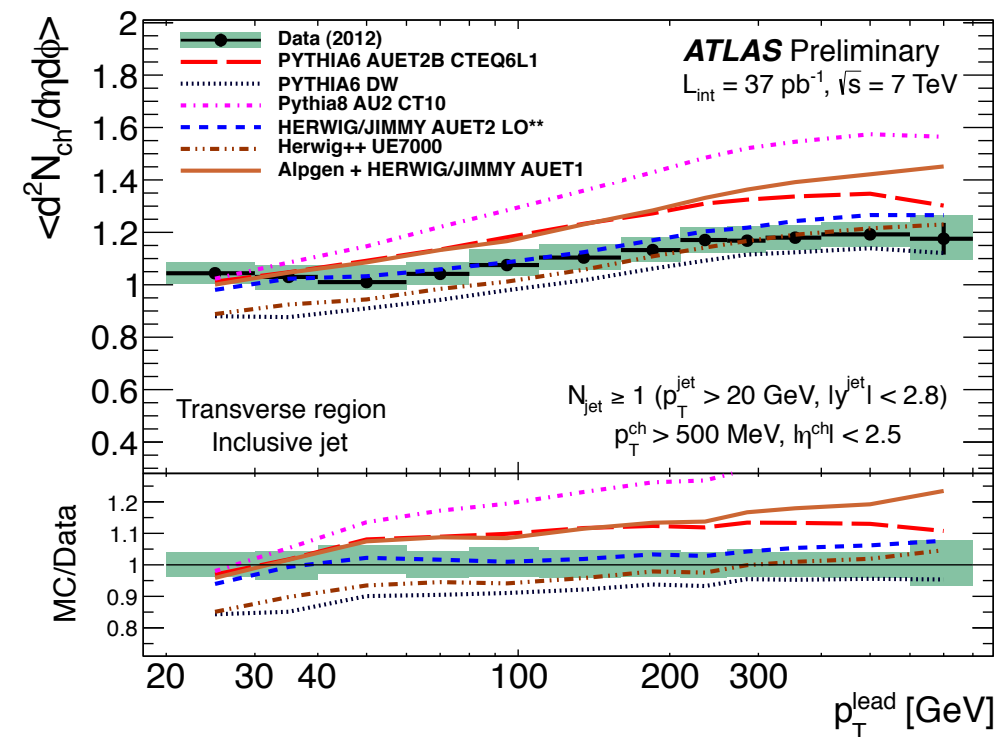
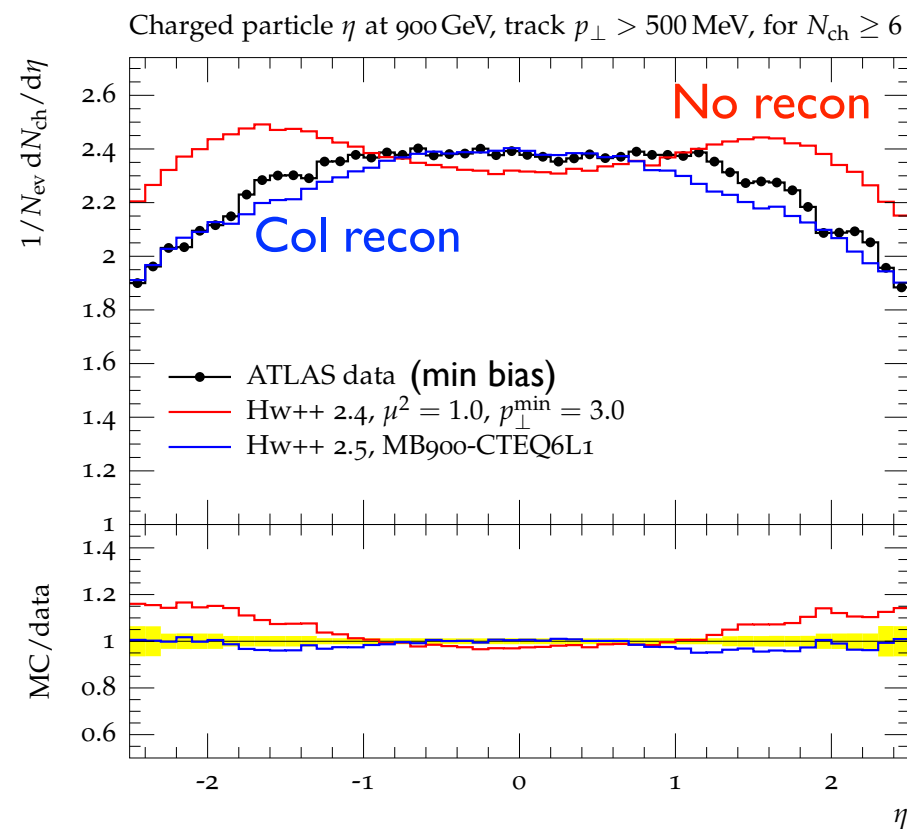
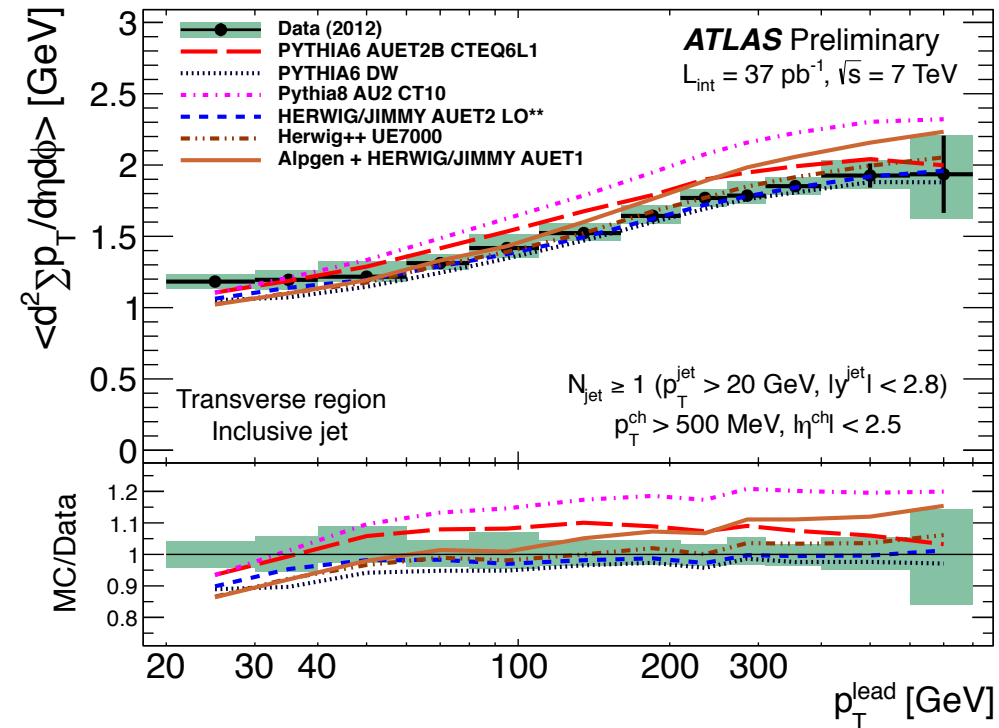
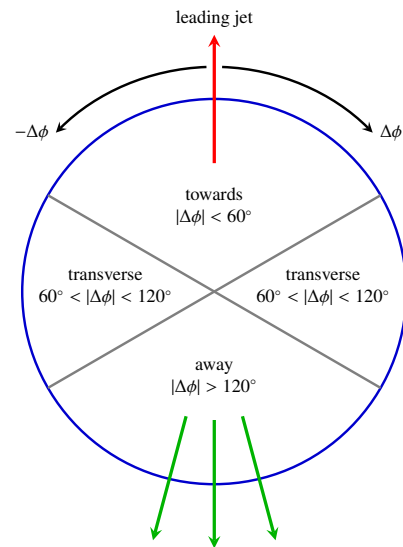
- No fundamental progress since 1980s
  - ✿ Available non-perturbative methods (lattice, AdS/QCD, ...) are inapplicable
- Less important in some respects in LHC era
  - ✿ Jets, leptons and photons are observed objects, not hadrons
- But still important for detector effects
  - ✿ Jet response, heavy-flavour tagging, lepton and photon isolation, ...

# Underlying Event



- Multiple parton interactions in same collision
  - ✦ Depends on density profile of proton
- Assume QCD 2-to-2 secondary collisions
  - ✦ Need cutoff at low  $p_T$
- Need to model colour flow
  - ✦ Colour reconnections are necessary

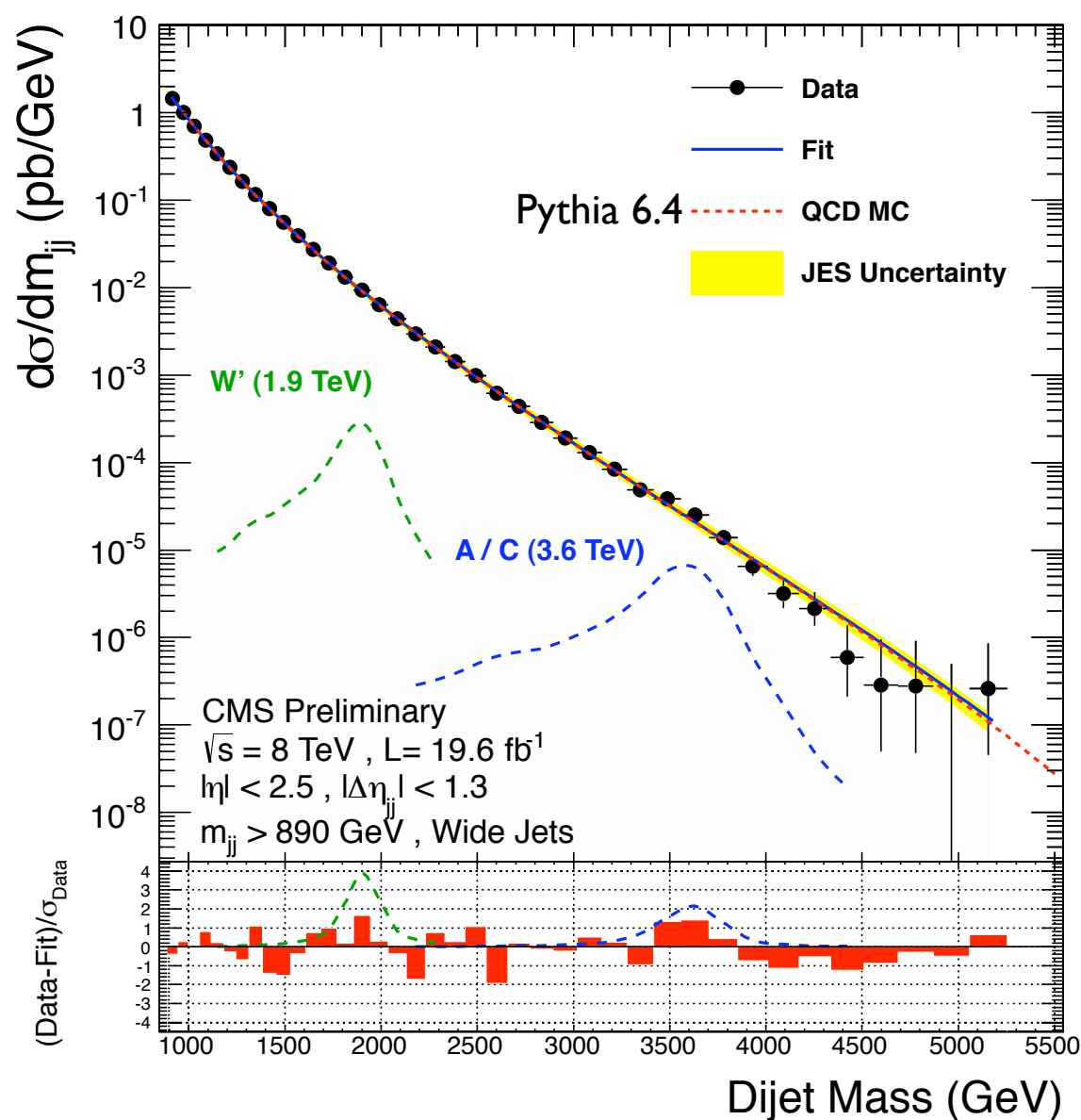
# Underlying Event



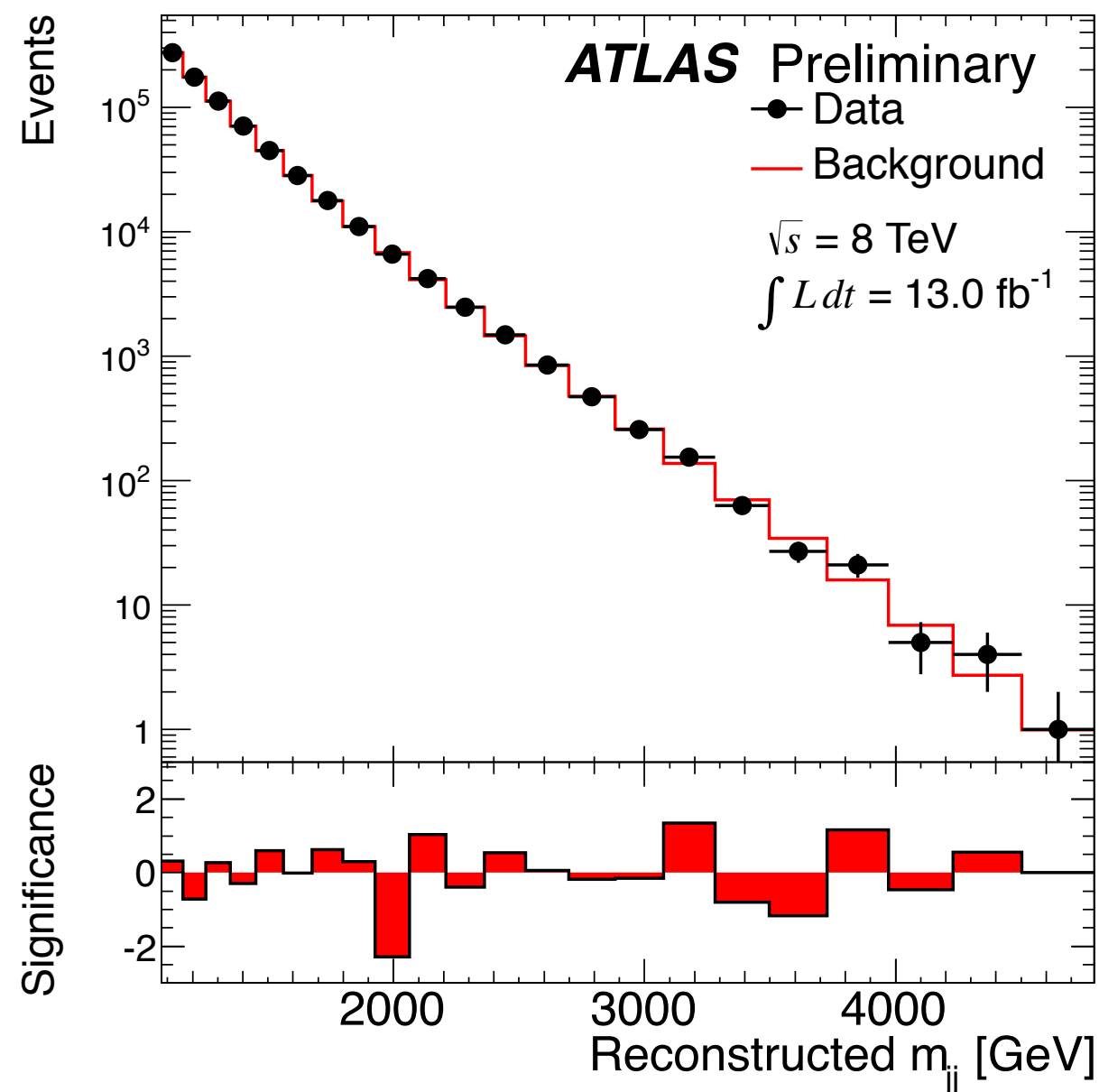
Gieseke, Röhr, Siódmok, arXiv:1206.2205

ATLAS CONF-2012-164

# Dijet Mass Distribution



CMS PAS EXO-12-059



ATLAS CONF-2012-148

- No sign of deviation from Standard Model (yet)

# MC Event Generators

## ● HERWIG

<http://projects.hepforge.org/herwig/>

- ➔ Angular-ordered parton shower, cluster hadronization
- ➔ v6 Fortran; Herwig++

## ● PYTHIA

<http://www.thep.lu.se/~torbjorn/Pythia.html>

- ➔ Dipole-type parton shower, string hadronization
- ➔ v6 Fortran; v8 C++

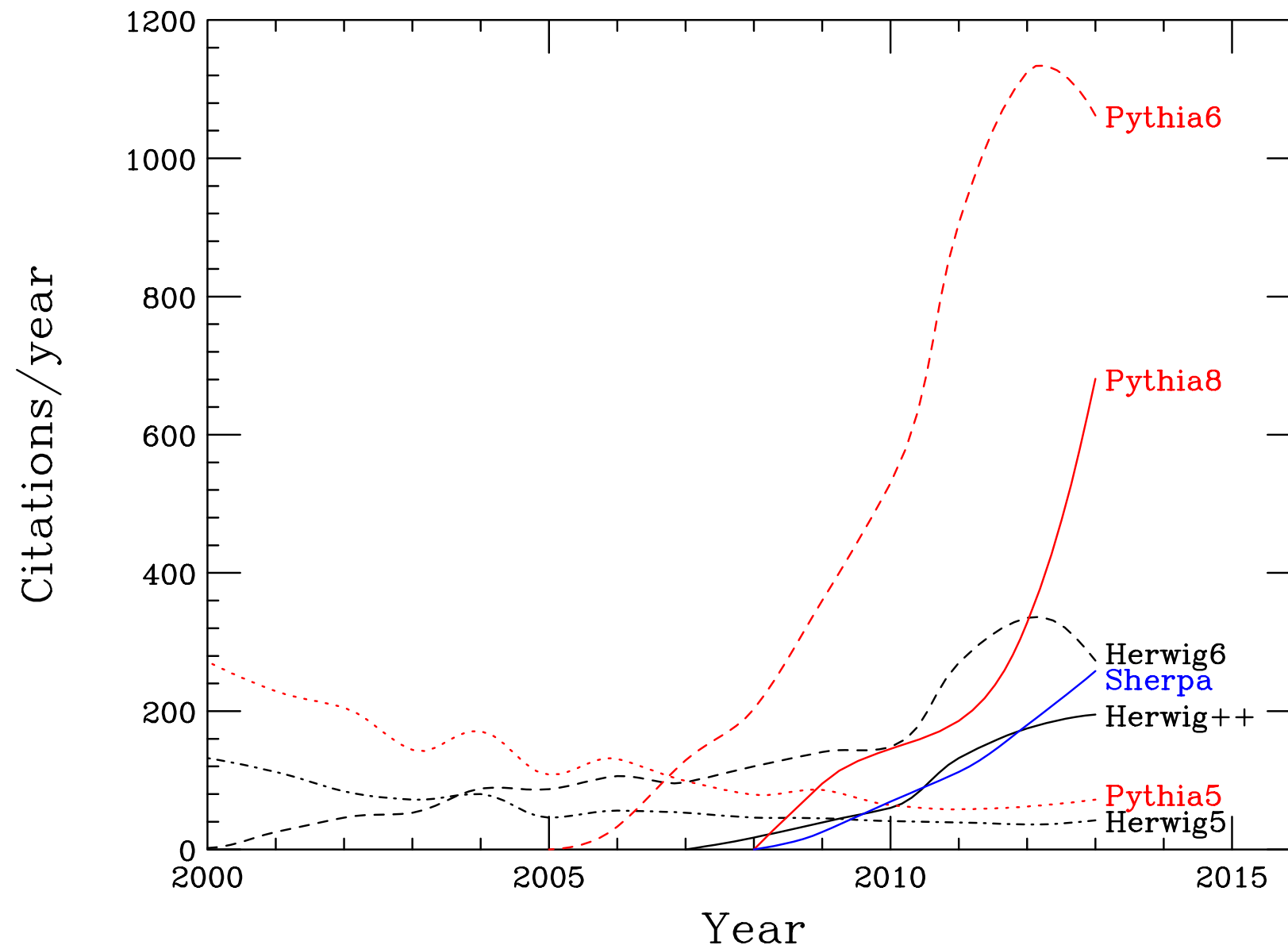
## ● SHERPA

<http://projects.hepforge.org/sherpa/>

- ➔ Dipole-type parton shower, cluster hadronization
- ➔ C++

“General-purpose event generators for LHC physics”,  
A Buckley et al., arXiv:1101.2599, Phys. Rept. 504(2011)145

# Generator Citations



- Most-cited article only for each version
- 2013 is extrapolation (Jan to mid-May x3)

# Other relevant software

## (with apologies for omissions)

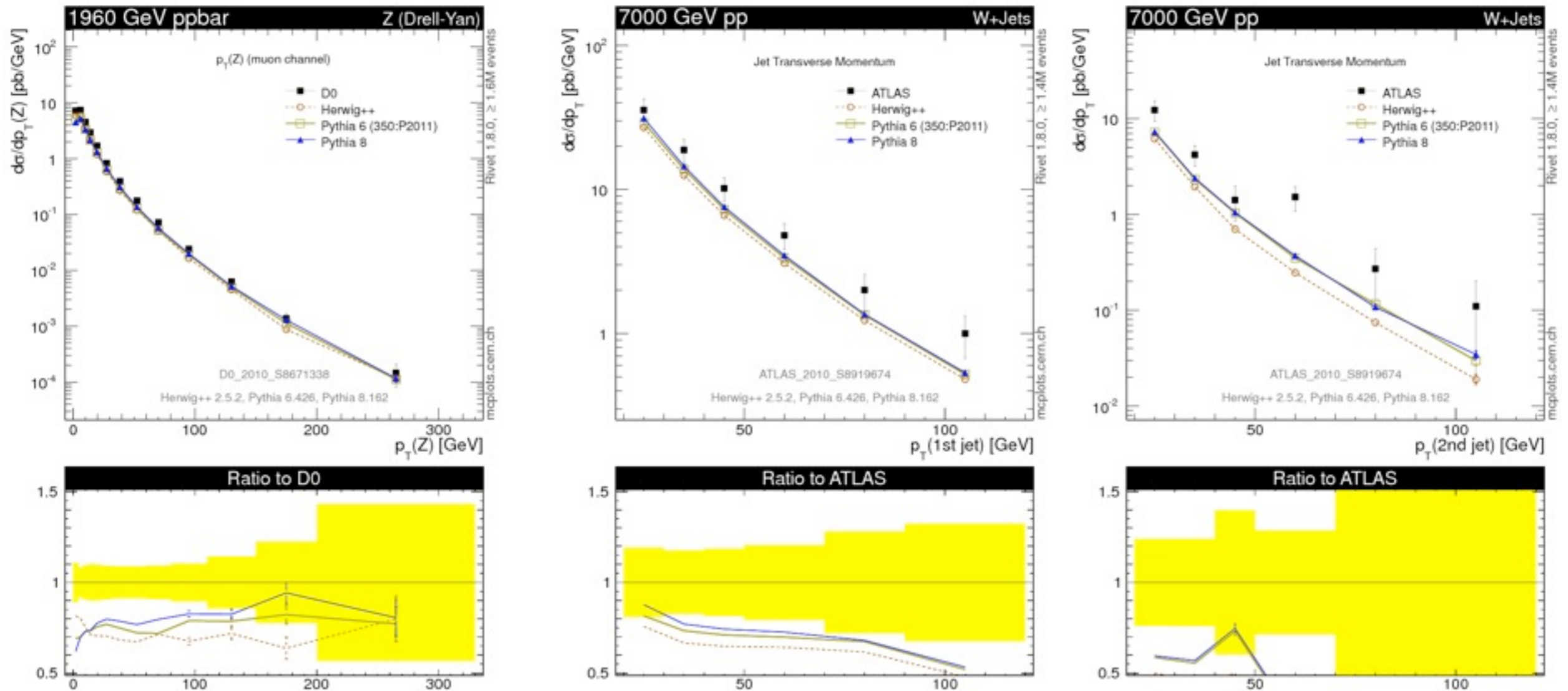
- Other event/shower generators: PhoJet, Ariadne, Dipsy, Cascade, Vincia
- Matrix-element generators: MadGraph/MadEvent, CompHep, CalcHep, Helac, Whizard, Sherpa, GoSam, aMC@NLO
- Matrix element libraries: AlpGen, POWHEG BOX, MCFM, NLOjet++, VBFNLO, BlackHat, Rocket
- Special BSM scenarios: Prospino, Charybdis, TrueNoir
- Mass spectra and decays: SOFTSUSY, SPHENO, HDecay, SDecay
- Feynman rule generators: FeynRules
- PDF libraries: LHAPDF
- Resummed ( $p_{\perp}$ ) spectra: ResBos
- Approximate loops: LoopSim
- Jet finders: anti- $k_{\perp}$  and FastJet
- Analysis packages: Rivet, Professor, MCPLOTS
- Detector simulation: GEANT, Delphes
- Constraints (from cosmology etc): DarkSUSY, MicrOmegas
- Standards: PDF identity codes, LHA, LHEF, SLHA, Binoth LHA, HepMC

Sjöstrand, Nobel Symposium, May 2013

# Parton Shower Monte Carlo

<http://mcplots.cern.ch/>

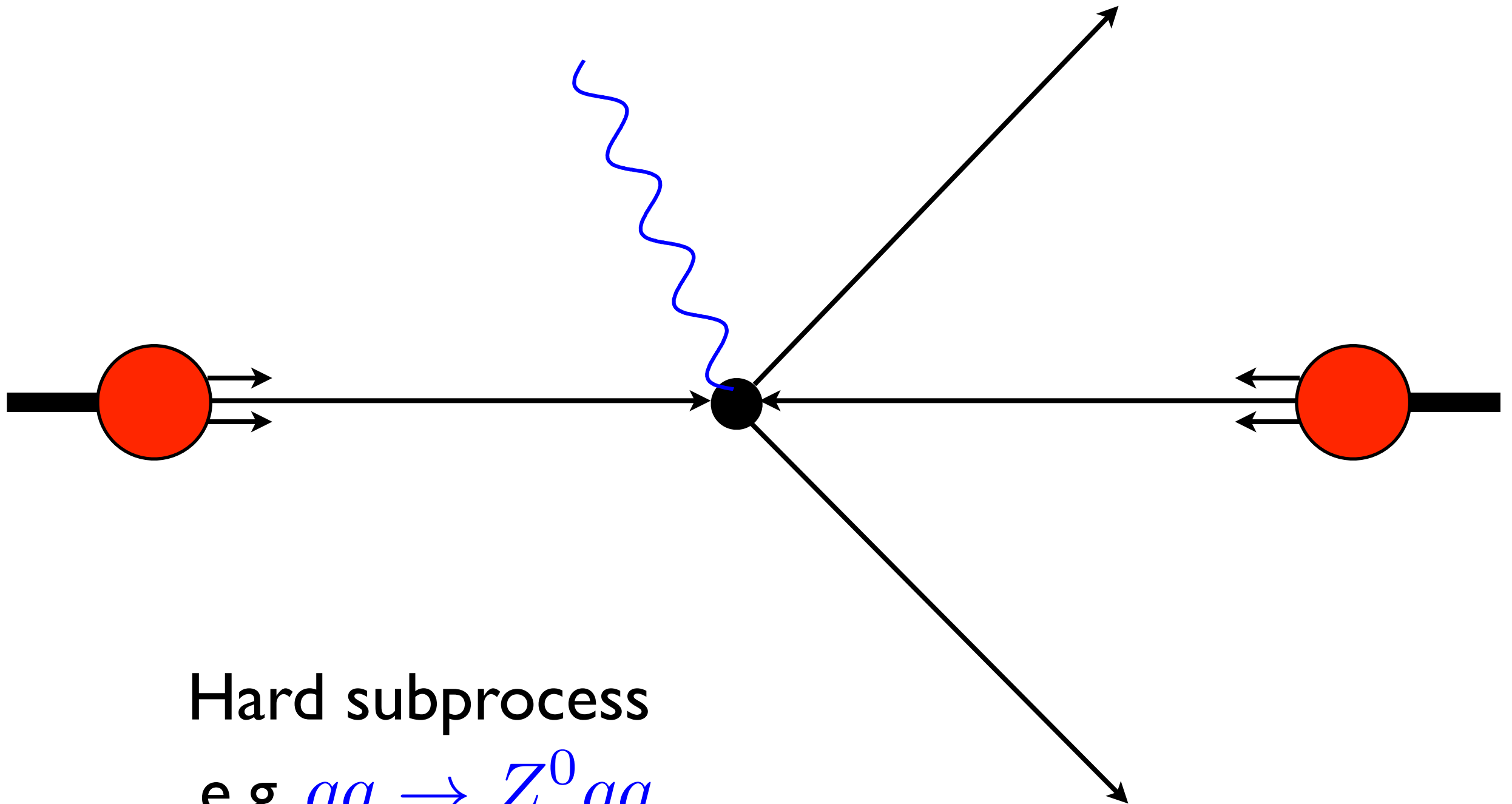
- Hard subprocess:  $q\bar{q} \rightarrow Z^0/W^\pm$



- Leading-order (LO) normalization  $\Rightarrow$  need next-to-LO (NLO)
- Worse for high  $p_T$  and/or extra jets  $\Rightarrow$  need multijet merging



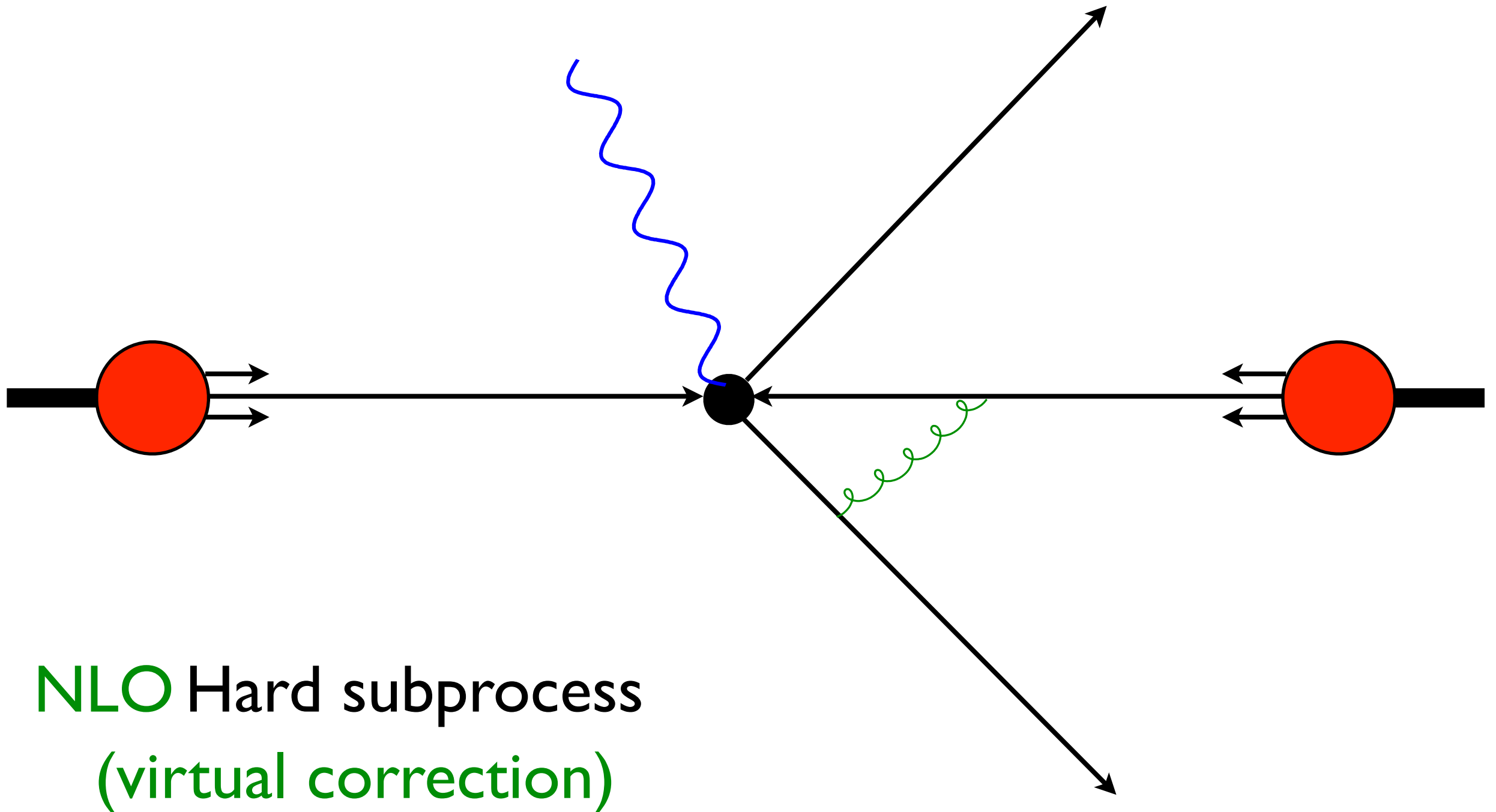
# Improving Event Simulation



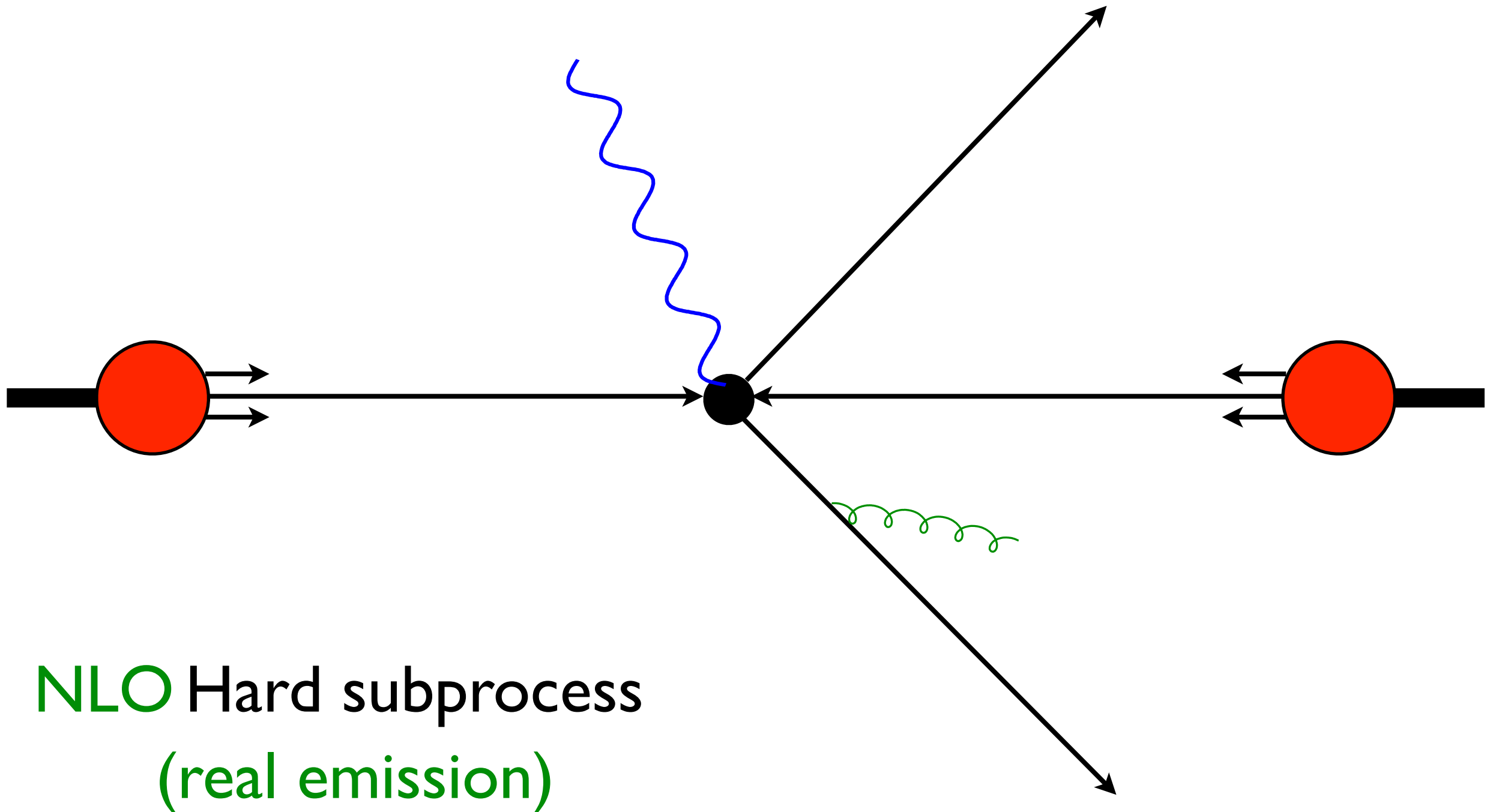
Hard subprocess

e.g.  $qq \rightarrow Z^0 qq$

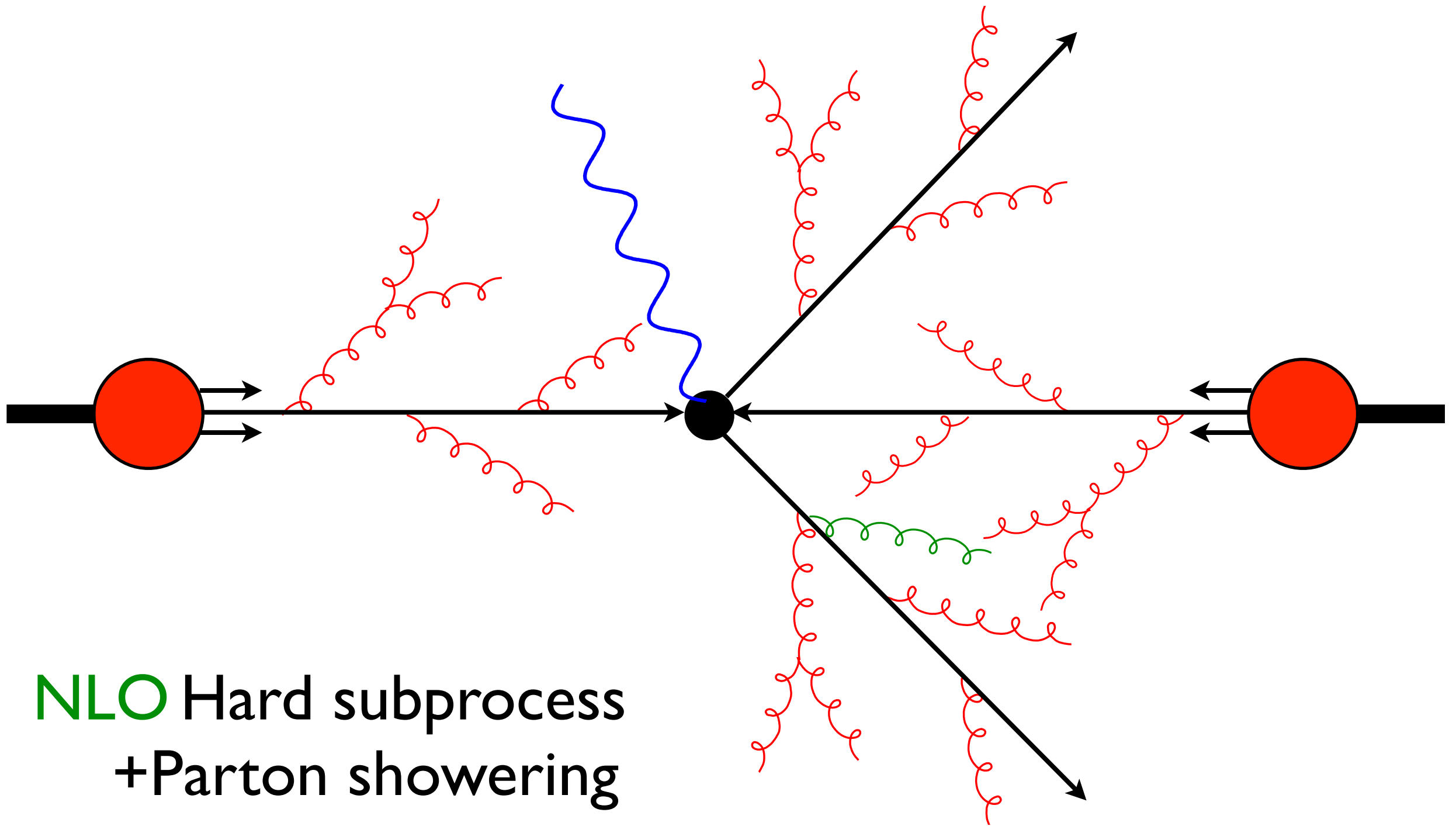
# Improving Event Simulation



# Improving Event Simulation

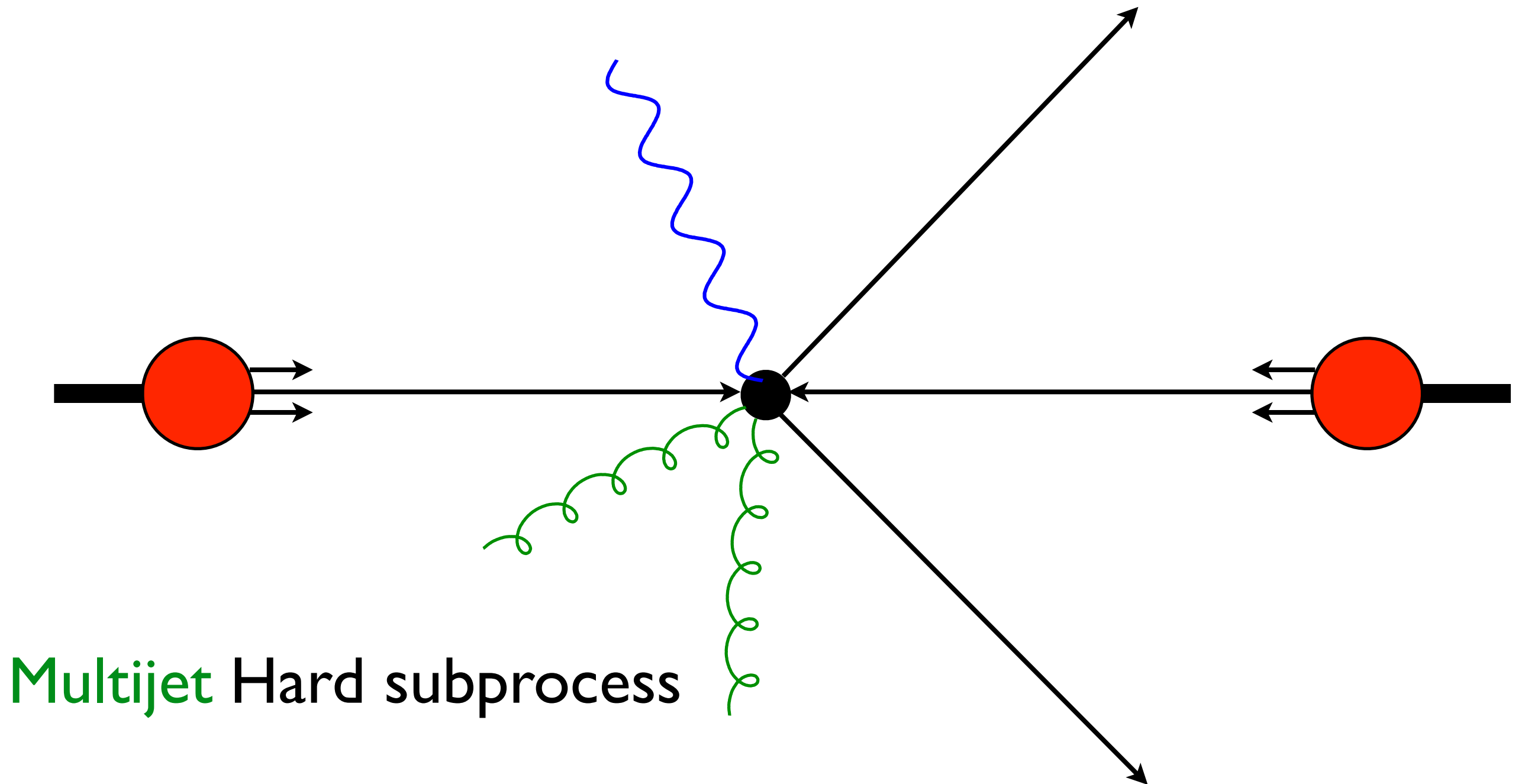


# Improving Event Simulation

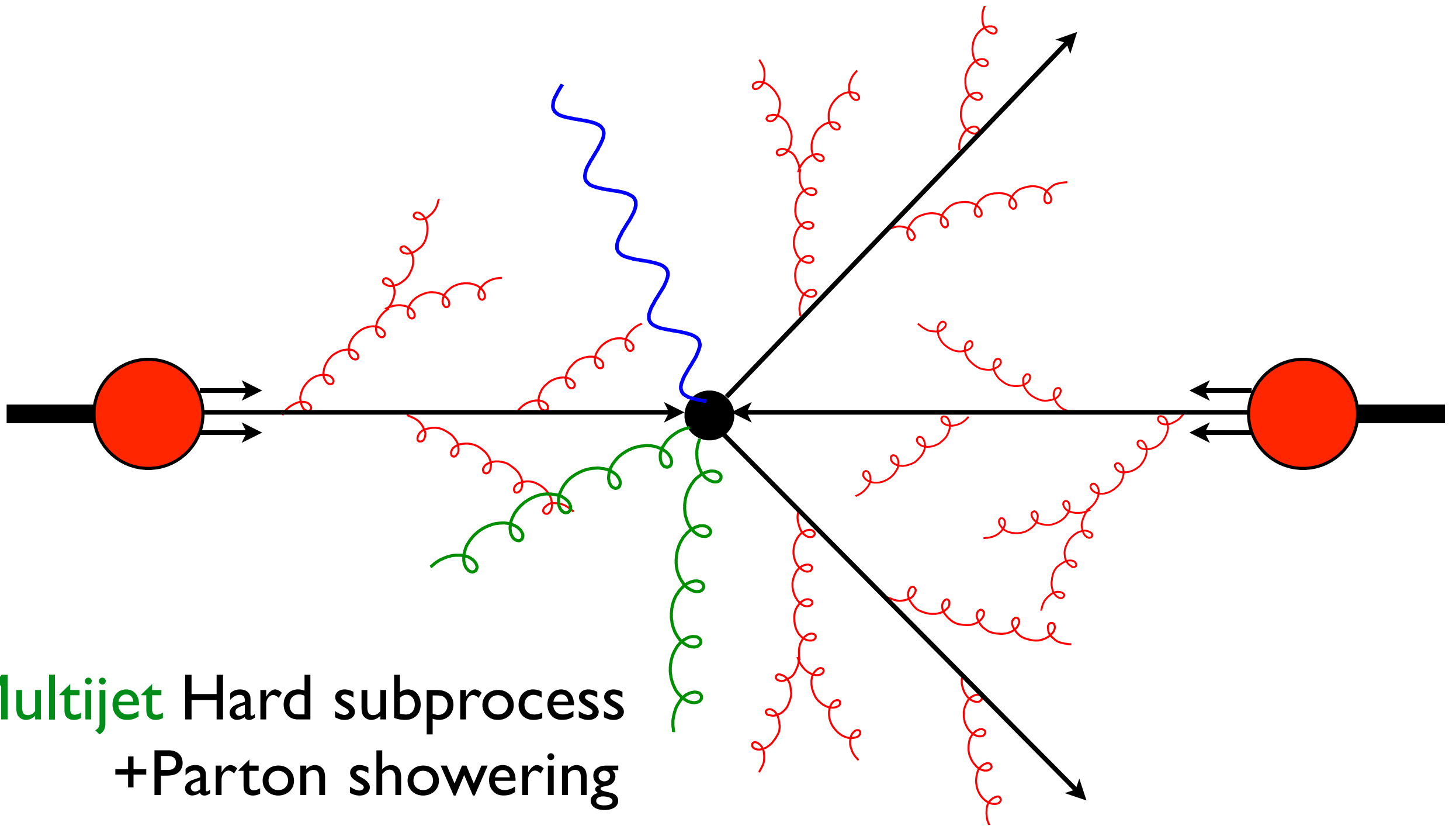


**NLO** Hard subprocess  
+ Parton showering  
= Double counting??

# Improving Event Simulation



# Improving Event Simulation



**Multijet** Hard subprocess  
+ Parton showering  
= Double counting??

# Matching & Merging

- Two rather different objectives:
- Matching parton showers to NLO matrix elements, without double counting
  - ✧ MC@NLO Frixione, BW, 2002
  - ✧ POWHEG Nason, 2004
- Merging parton showers with LO n-jet matrix elements, minimizing jet resolution dependence
  - ✧ CKKW Catani, Krauss, Kühn, BW, 2001
  - ✧ Dipole Lönnblad, 2001
  - ✧ MLM merging Mangano, 2002

# MC@NLO matching

S Frixione & BW, JHEP 06(2002)029

- Compute parton shower contributions (real and virtual) at NLO
  - ✧ Generator-dependent
- Subtract these from exact NLO
  - ✧ Cancels divergences of exact NLO!
- Generate modified no-emission (LO+virtual) and real-emission hard process configurations
  - ✧ Some may have negative weight
- Pass these through parton shower etc.
  - ✧ Only shower-generated terms beyond NLO



# MC@NLO matching

S Frixione & BW, JHEP 06(2002)029

finite virtual

divergent

$$\begin{aligned} d\sigma_{\text{NLO}} &= \left[ B(\Phi_B) + V(\Phi_B) - \int \sum_i C_i(\Phi_B, \Phi_R) d\Phi_R \right] d\Phi_B + R(\Phi_B, \Phi_R) d\Phi_B d\Phi_R \\ &\equiv \left[ B + V - \int C d\Phi_R \right] d\Phi_B + R d\Phi_B d\Phi_R \end{aligned}$$

$$\begin{aligned} d\sigma_{\text{MC}} &= B(\Phi_B) d\Phi_B \left[ \Delta_{\text{MC}}(0) + \frac{R_{\text{MC}}(\Phi_B, \Phi_R)}{B(\Phi_B)} \Delta_{\text{MC}}(k_T(\Phi_B, \Phi_R)) d\Phi_R \right] \\ &\equiv B d\Phi_B [\Delta_{\text{MC}}(0) + (R_{\text{MC}}/B) \Delta_{\text{MC}}(k_T) d\Phi_R] \end{aligned}$$

$$\begin{aligned} d\sigma_{\text{MC@NLO}} &= \left[ B + V + \int (R_{\text{MC}} - C) d\Phi_R \right] d\Phi_B [\Delta_{\text{MC}}(0) + (R_{\text{MC}}/B) \Delta_{\text{MC}}(k_T) d\Phi_R] \\ &\quad + (R - R_{\text{MC}}) \Delta_{\text{MC}}(k_T) d\Phi_B d\Phi_R \end{aligned}$$

finite  $\geq 0$

MC starting from no emission

MC starting from one emission

- Expanding gives NLO result

# POWHEG matching

P Nason, JHEP 11(2004)040

- POsitive Weight Hardest Emission Generator
- Use exact real-emission matrix element to generate hardest (highest relative  $p_T$ ) emission configurations
  - ✧ No-emission probability implicitly modified
  - ✧ (Almost) eliminates negative weights
  - ✧ Some uncontrolled terms generated beyond NLO
- Pass configurations through parton shower etc

# POWHEG matching

P Nason, JHEP 11(2004)040

$$d\sigma_{\text{MC}} = B(\Phi_B) d\Phi_B \left[ \Delta_{\text{MC}}(0) + \frac{R_{\text{MC}}(\Phi_B, \Phi_R)}{B(\Phi_B)} \Delta_{\text{MC}}(k_T(\Phi_B, \Phi_R)) d\Phi_R \right]$$

$$d\sigma_{\text{PH}} = \bar{B}(\Phi_B) d\Phi_B \left[ \Delta_R(0) + \frac{R(\Phi_B, \Phi_R)}{B(\Phi_B)} \Delta_R(k_T(\Phi_B, \Phi_R)) d\Phi_R \right]$$

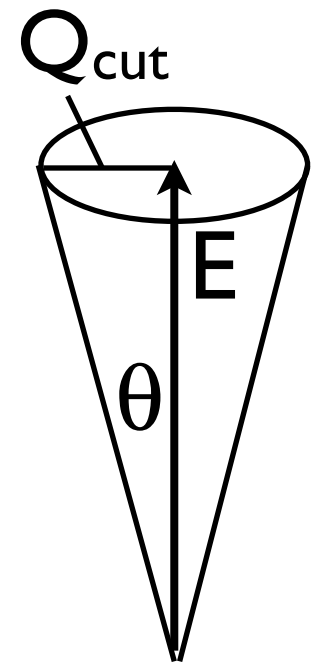
$$\bar{B}(\Phi_B) = B(\Phi_B) + V(\Phi_B) + \int \left[ R(\Phi_B, \Phi_R) - \sum_i C_i(\Phi_B, \Phi_R) \right] d\Phi_R$$

$$\Delta_R(p_T) = \exp \left[ - \int d\Phi_R \frac{R(\Phi_B, \Phi_R)}{B(\Phi_B)} \theta(k_T(\Phi_B, \Phi_R) - p_T) \right]$$

- NLO with (almost) no negative weights arbitrary NNLO
- High  $p_T$  always enhanced by  $K = \bar{B}/B = 1 + \mathcal{O}(\alpha_s)$

# Multijet Merging

- Objective: merge LO n-jet matrix elements\* with parton showers such that:
  - ✧ Multijet rates for jet resolution  $> Q_{\text{cut}}$  are correct to LO (up to  $N_{\text{max}}$ )
  - ✧ Shower generates jet structure below  $Q_{\text{cut}}$  (and jets above  $N_{\text{max}}$ )
  - ✧ Leading (and next)  $Q_{\text{cut}}$  dependence cancels



\* ALPGEN or MadGraph,  $n \leq N_{\text{max}}$

CKKW: Catani et al., JHEP 11(2001)063

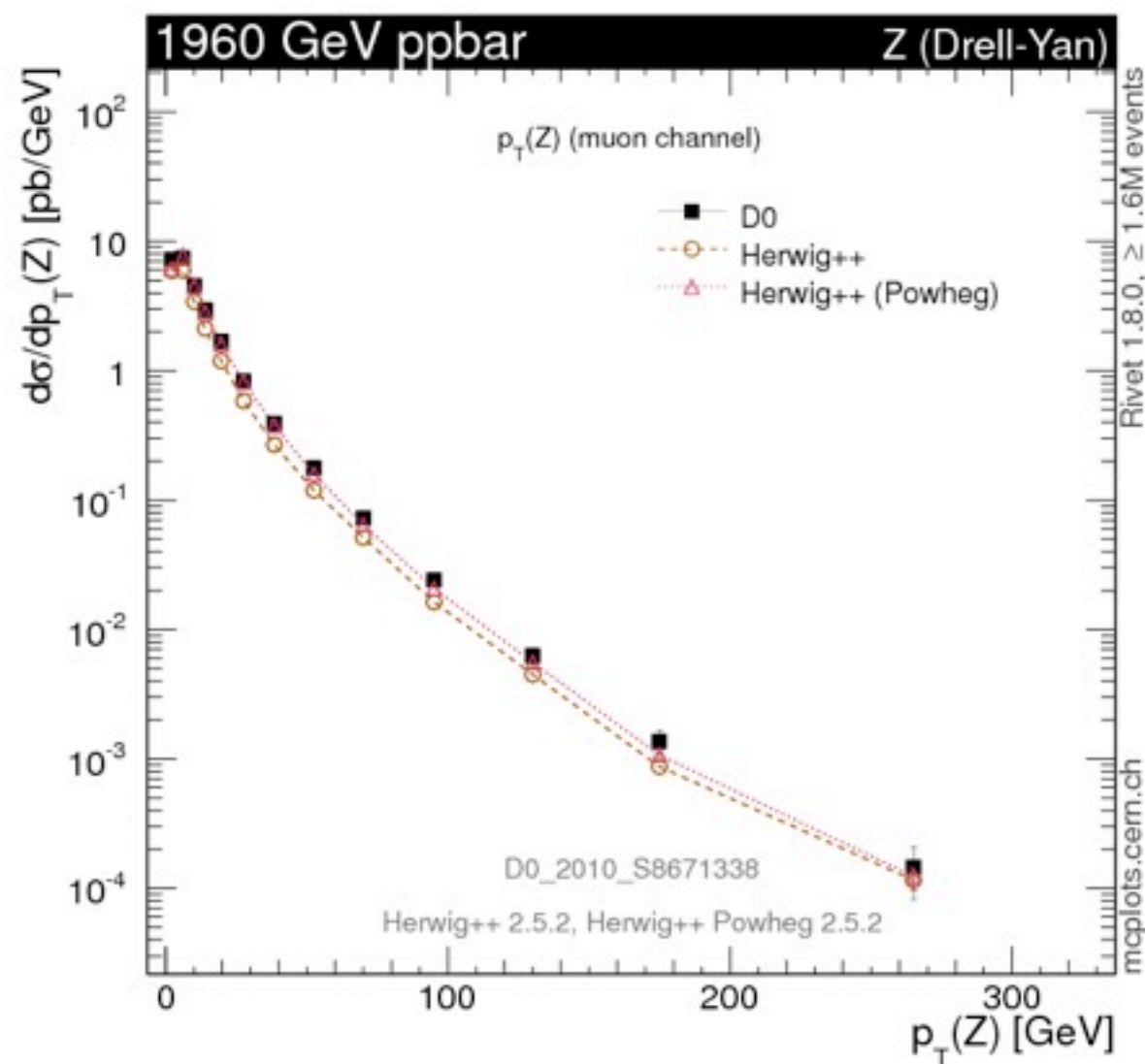
-L: Lonnblad, JHEP 05(2002)063

MLM: Mangano et al., NP B632(2002)343

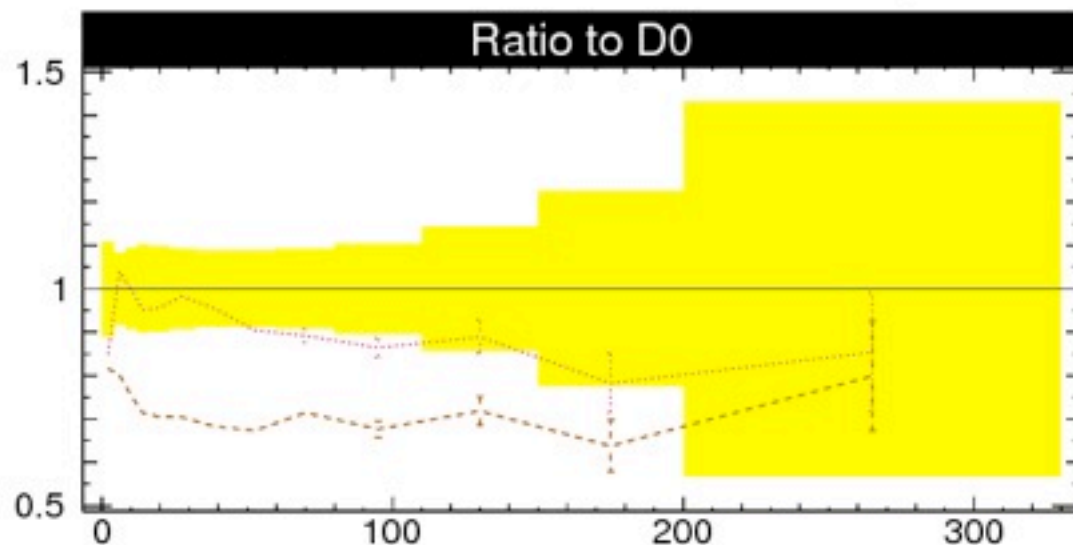
# Vector boson production

# $Z^0$ at Tevatron

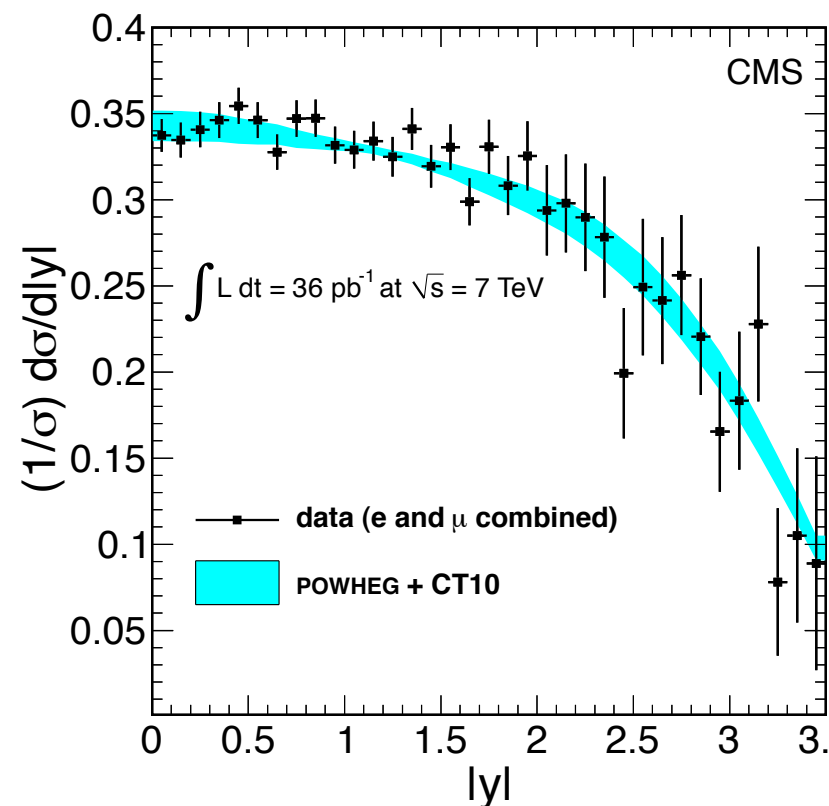
<http://mcplots.cern.ch/>



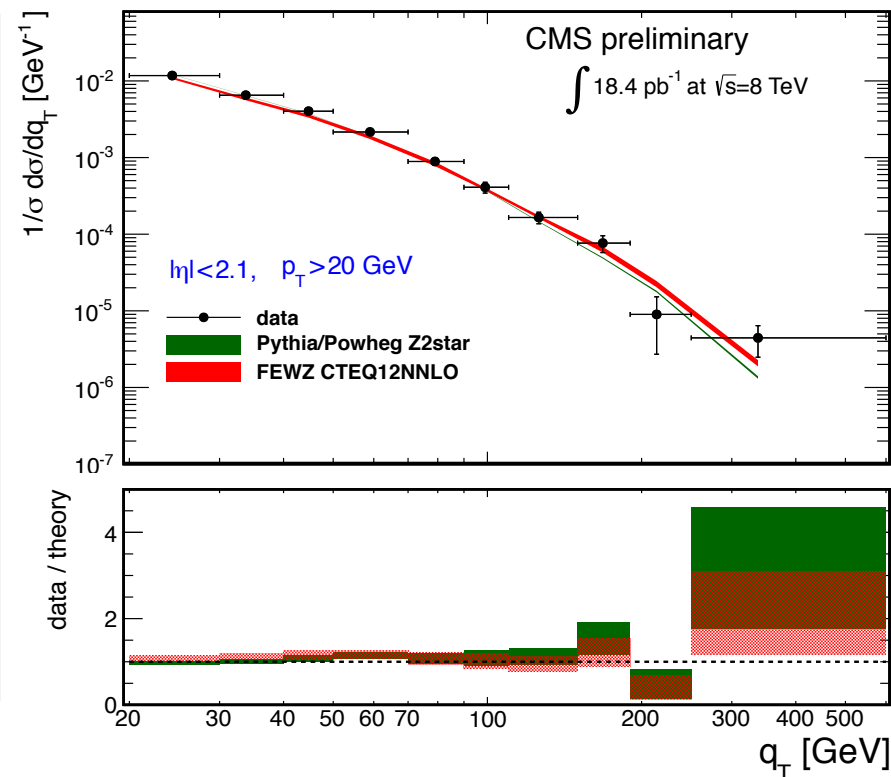
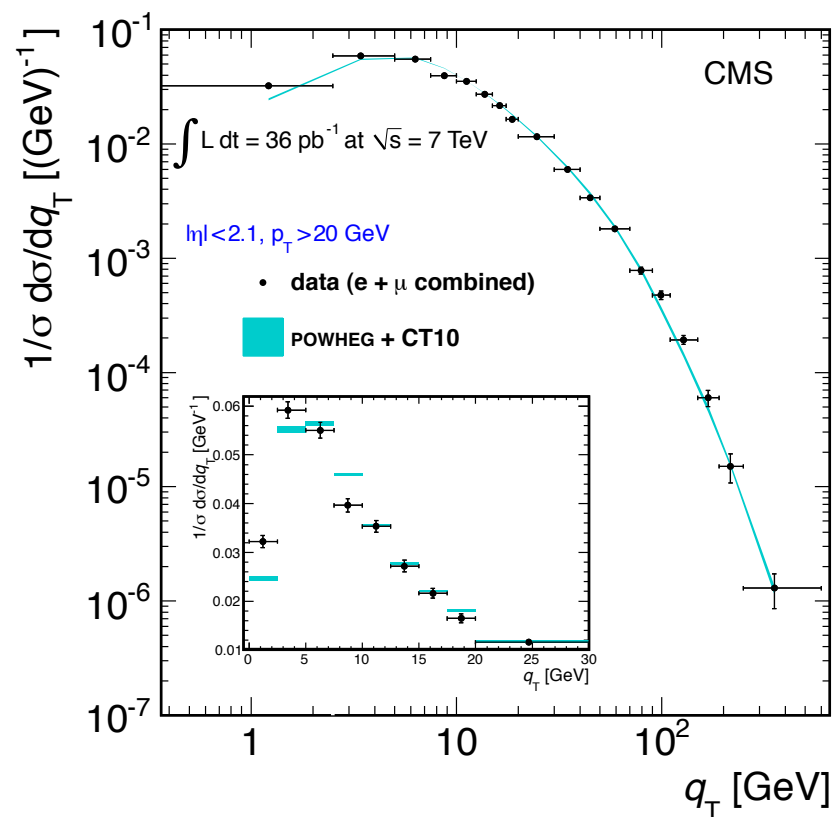
- Absolute normalization: LO too low
- POWHEG agrees with rate and distribution



# $Z^0$ at LHC



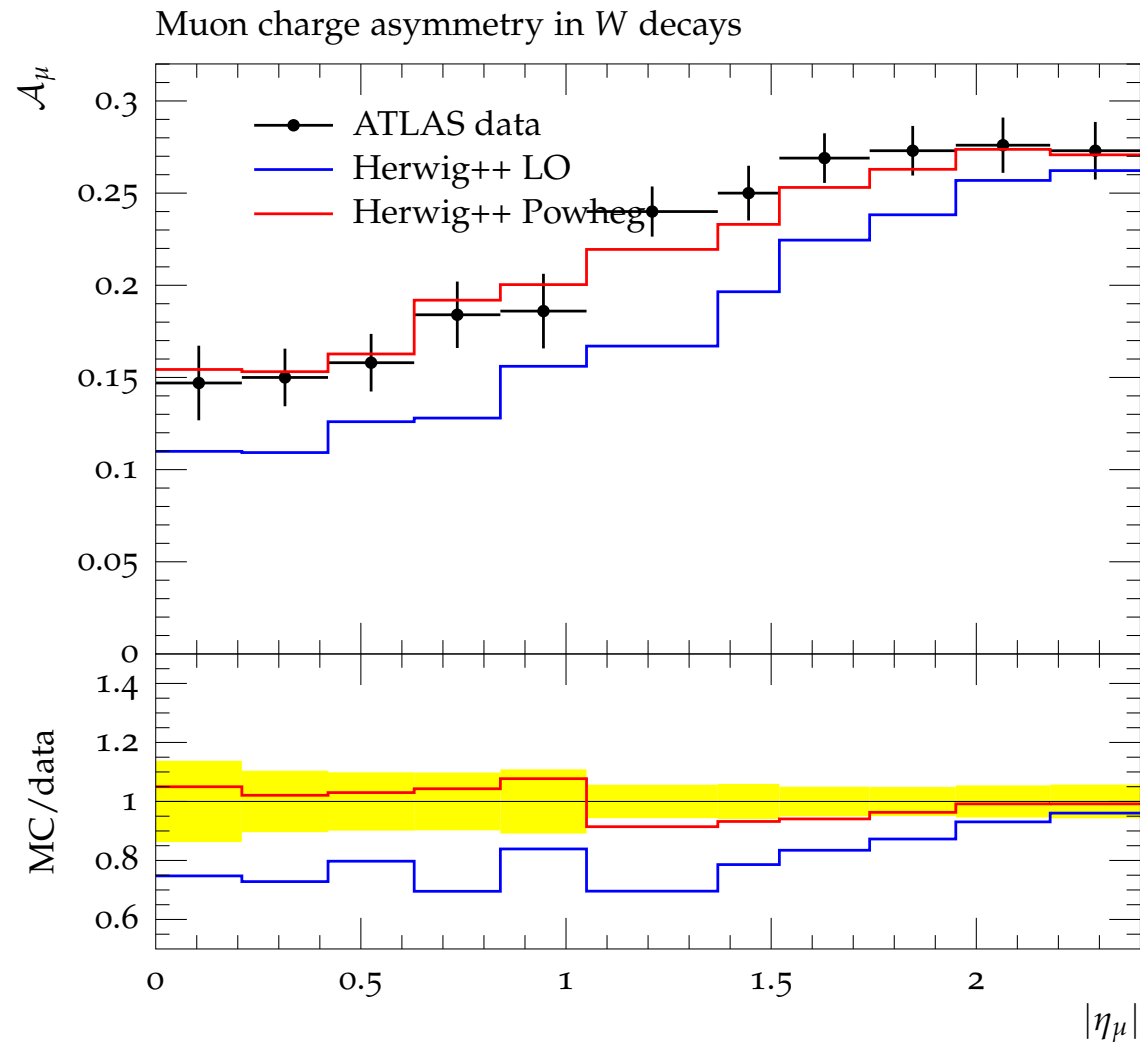
CMS, PRD85(2012)032002



CMS PAS SMP-12-025

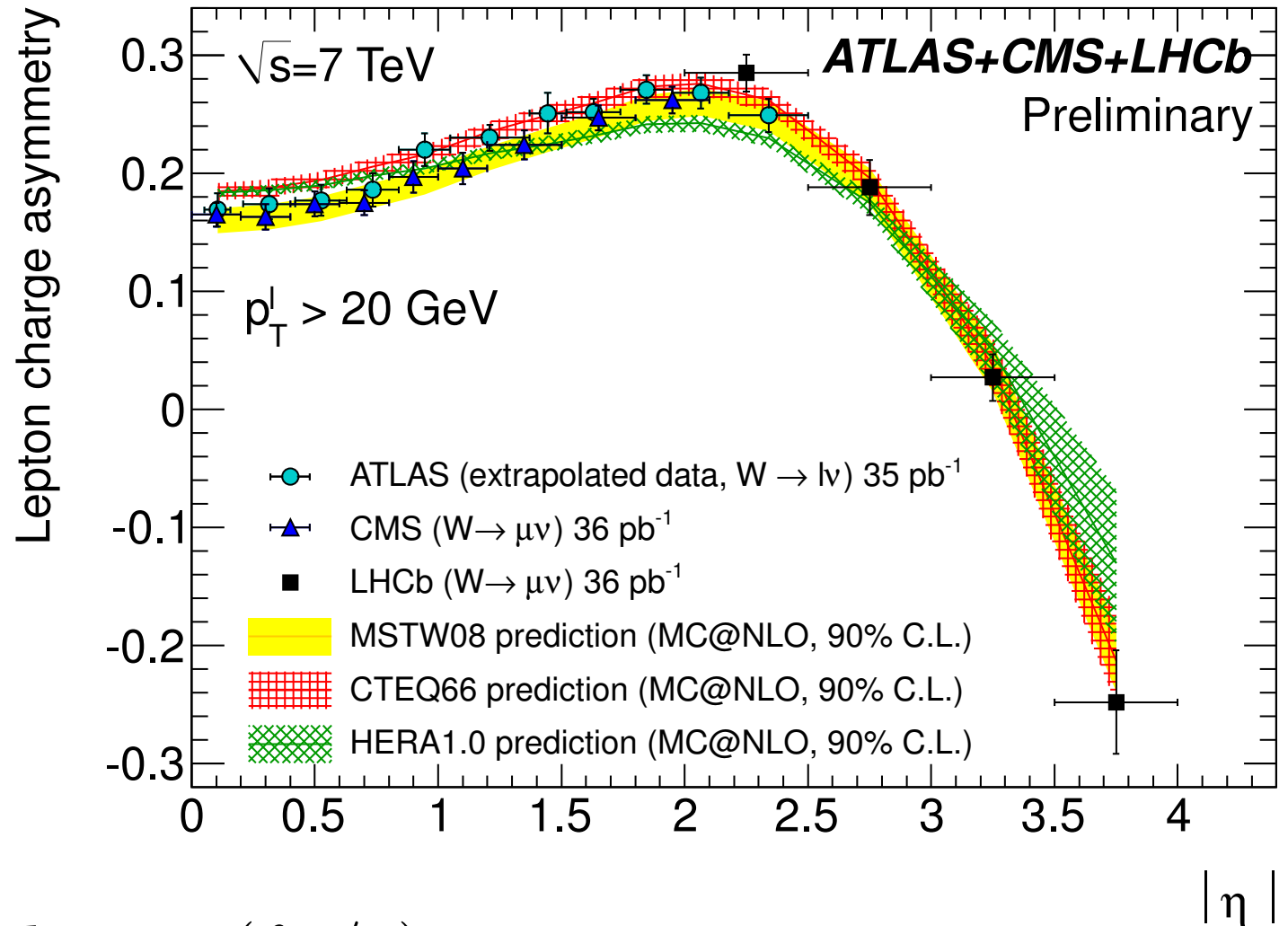
- Normalized to data
- POWHEG agrees with distribution (and NNLO)

# W asymmetry at LHC



$$A_\mu = \frac{N(\mu^+) - N(\mu^-)}{N(\mu^+) + N(\mu^-)}$$

$$\eta_\mu = \log \tan(\theta_\mu/2)$$



ATLAS-CONF-1211-129

- Asymmetry probes parton distributions

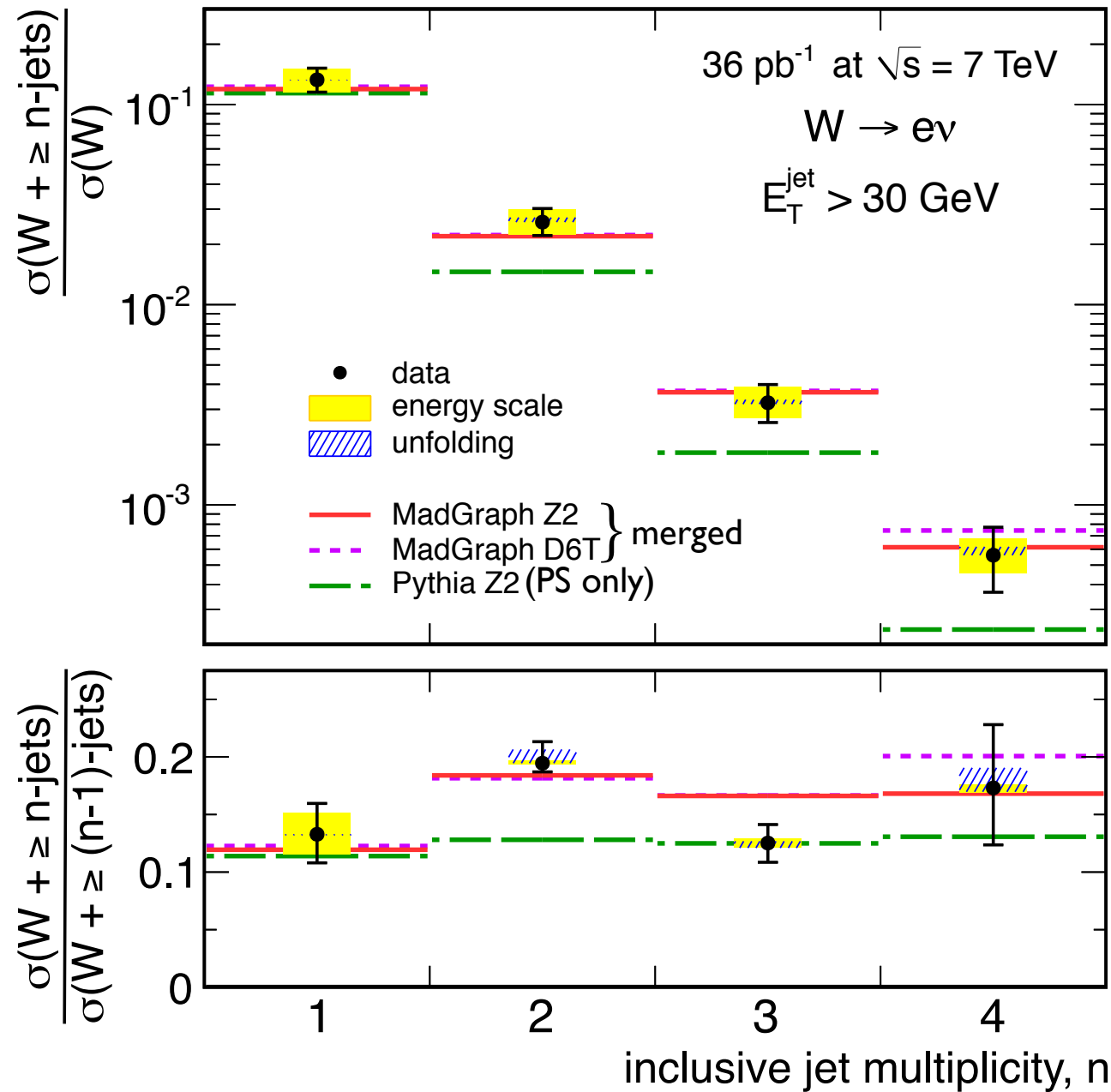
$$u\bar{d} \rightarrow W^+ \rightarrow \mu^+ \nu_\mu \quad \text{vs} \quad d\bar{u} \rightarrow W^- \rightarrow \mu^- \bar{\nu}_\mu$$



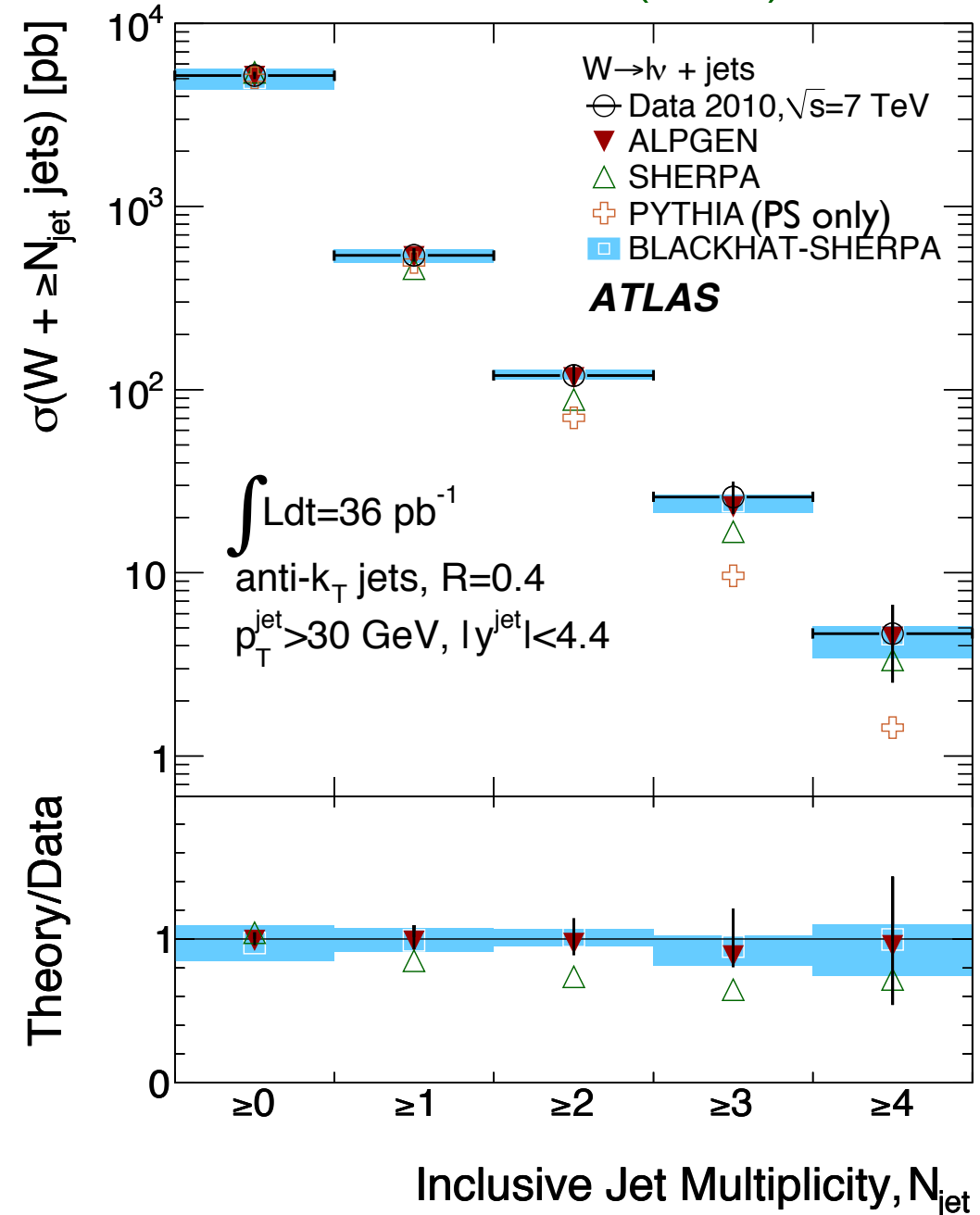
# W+jets at LHC

CMS, JHEP01(2012)010

CMS



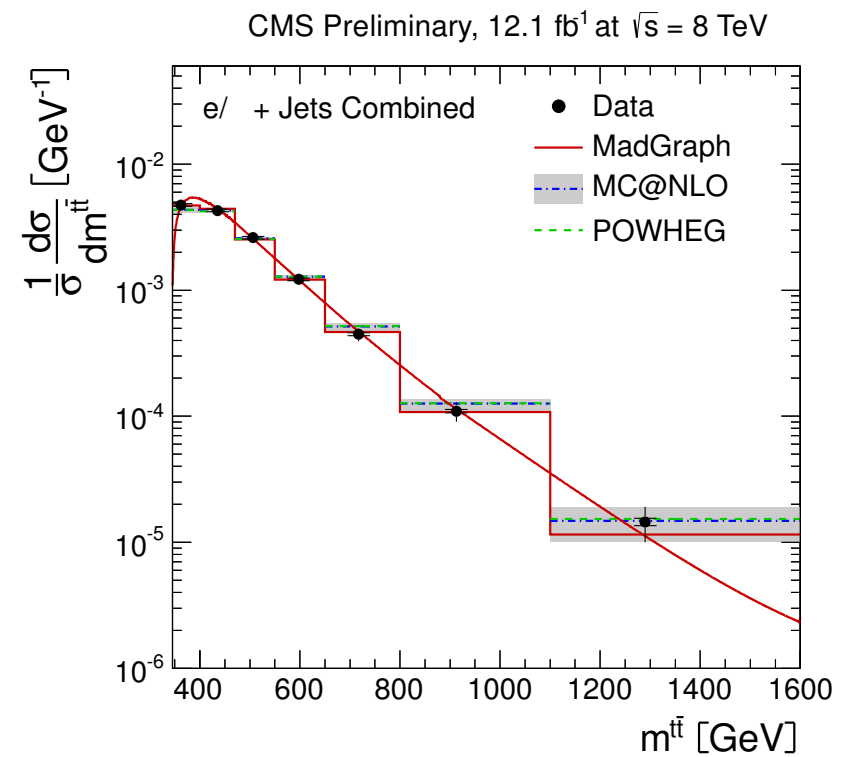
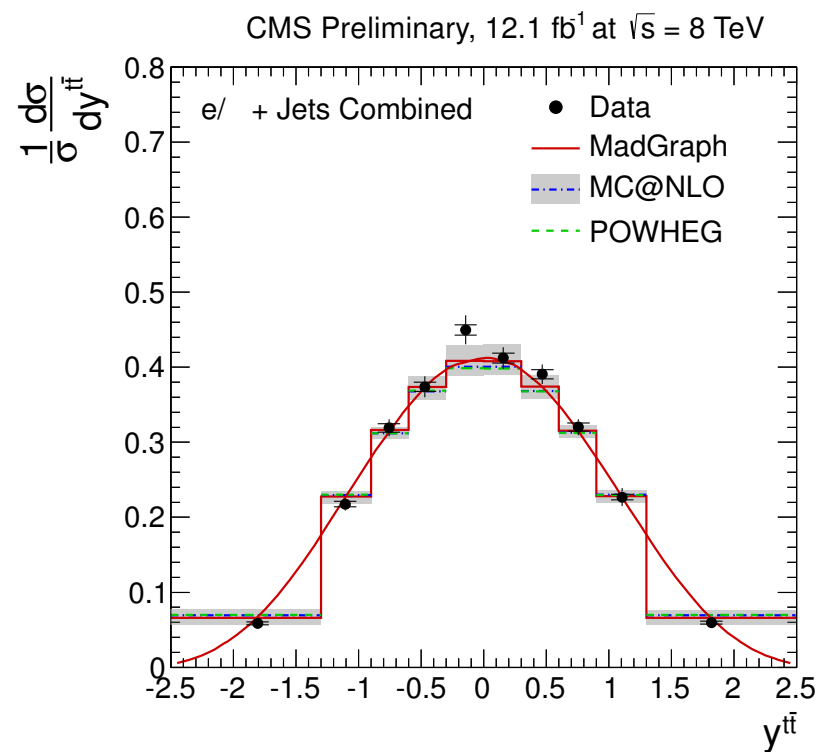
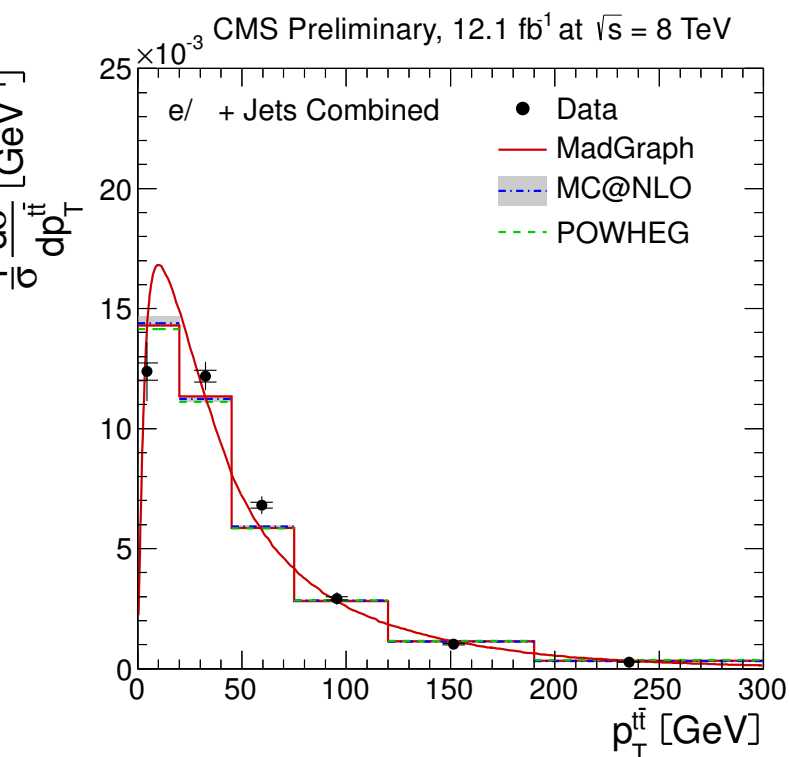
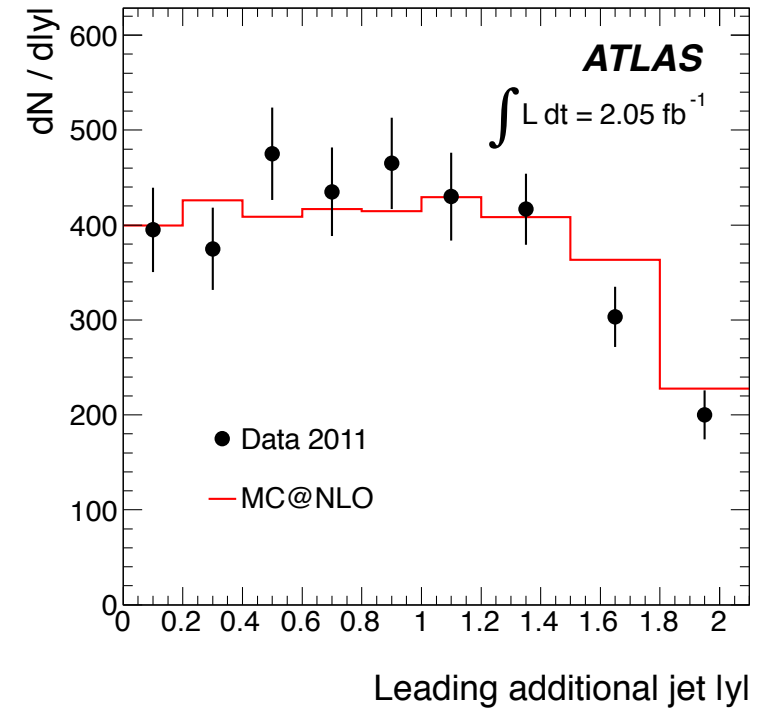
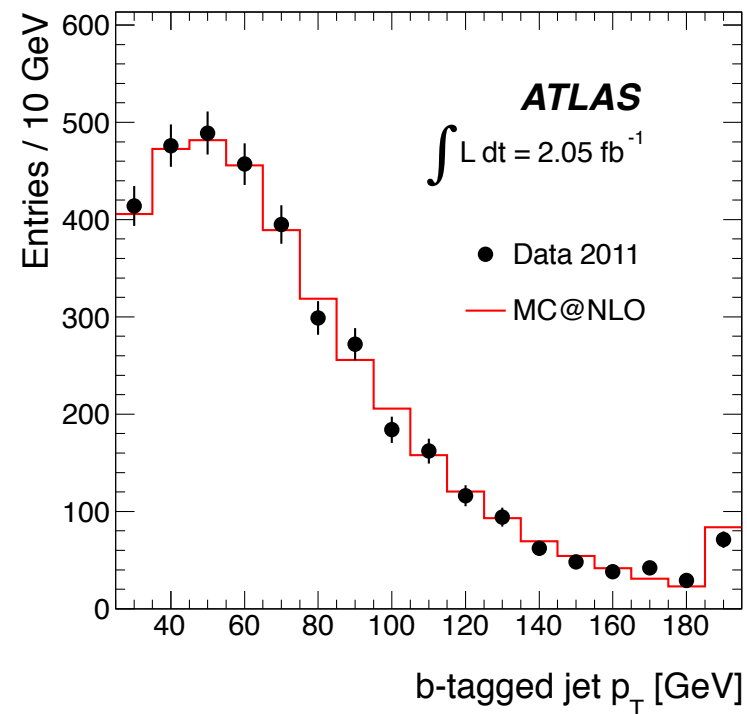
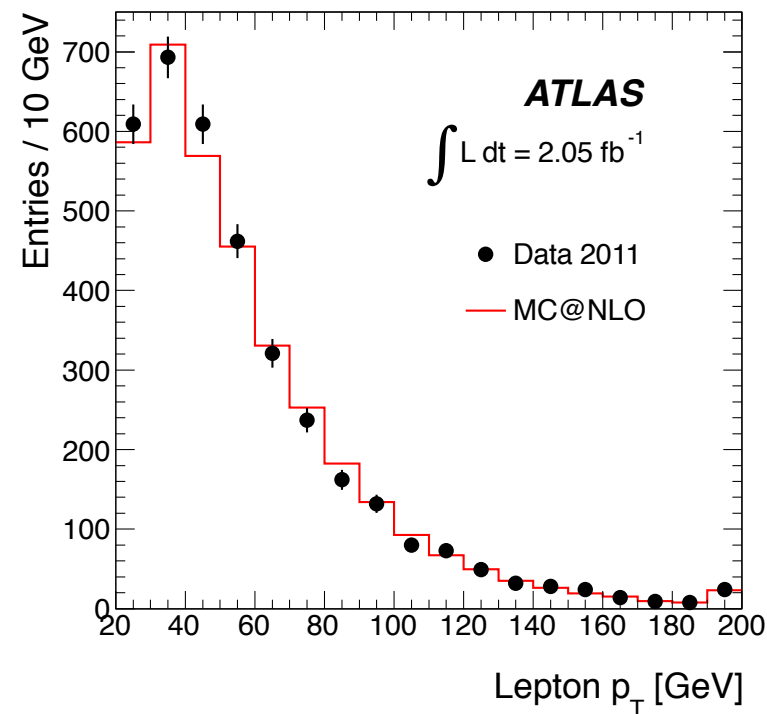
ATLAS, PRD85(2012)092002



- Very good agreement with predictions from merged simulations, while parton shower alone starts to fail for  $n_{\text{jet}} \geq 2$

# Top quark pair production

# Top quark pairs at LHC



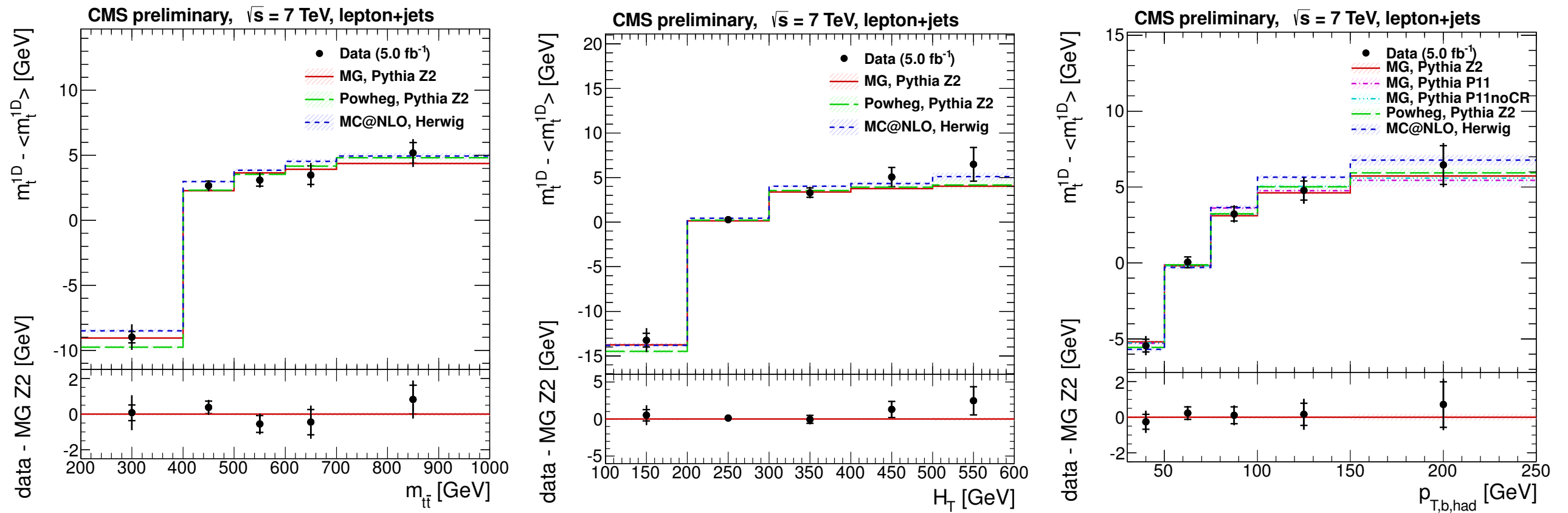
ATLAS, arXiv:1203.5015

CMS PAS TOP-12-027

Frixione, Nason, BW, JHEP 08(2003)007

Alioli, Nason, Oleari, Re, JHEP 06(2010)043

# Top mass & kinematics



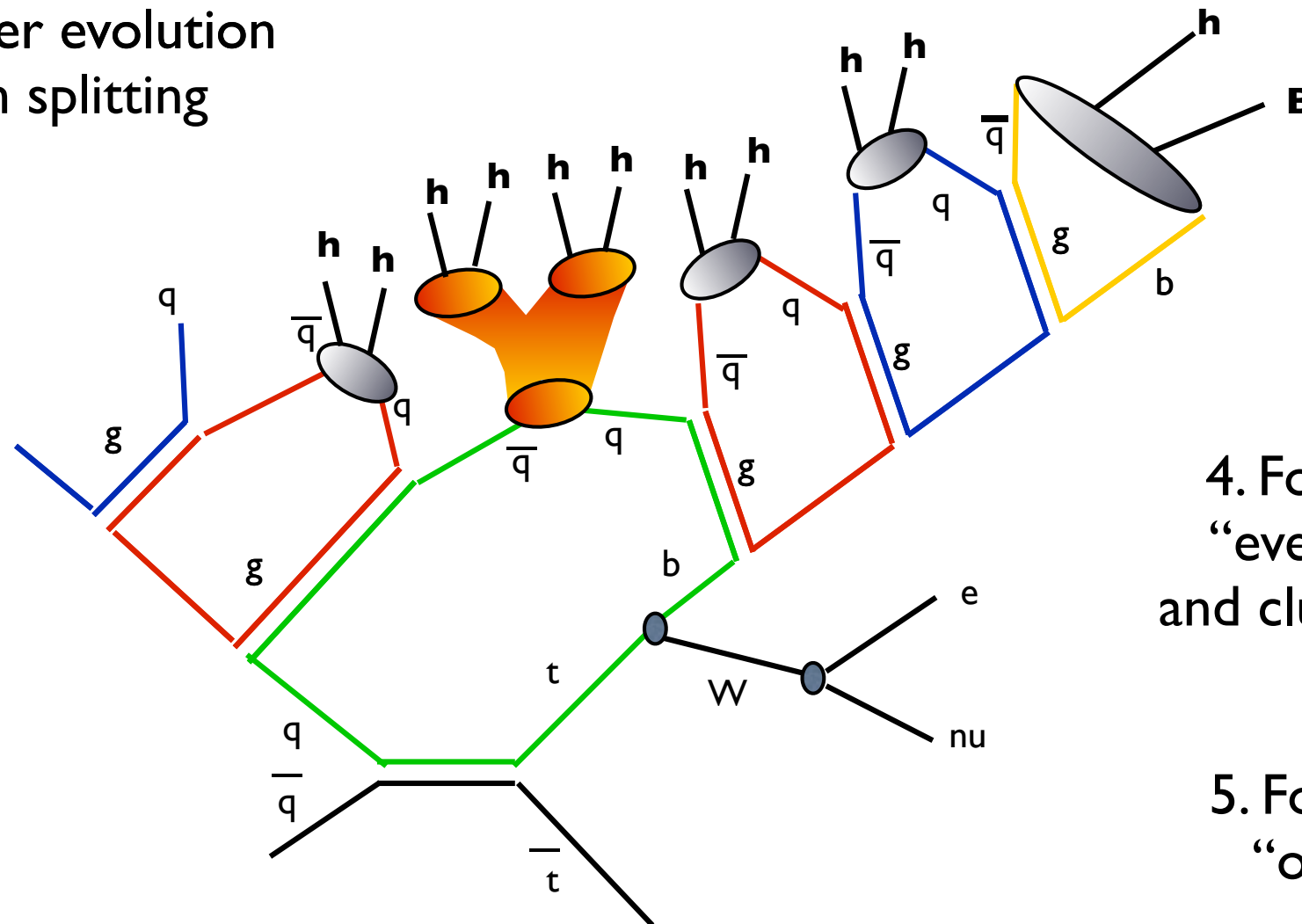
CMS PAS TOP-12-029

- Reconstructed top mass depends on kinematics
- But different generators track data well with a common input mass

# Top mass & hadronization

Mangano, Top LHC WG, July 2012

1. Hard Process
2. Shower evolution
3. Gluon splitting



4. Formation of  
“even” clusters  
and cluster decay  
to hadrons

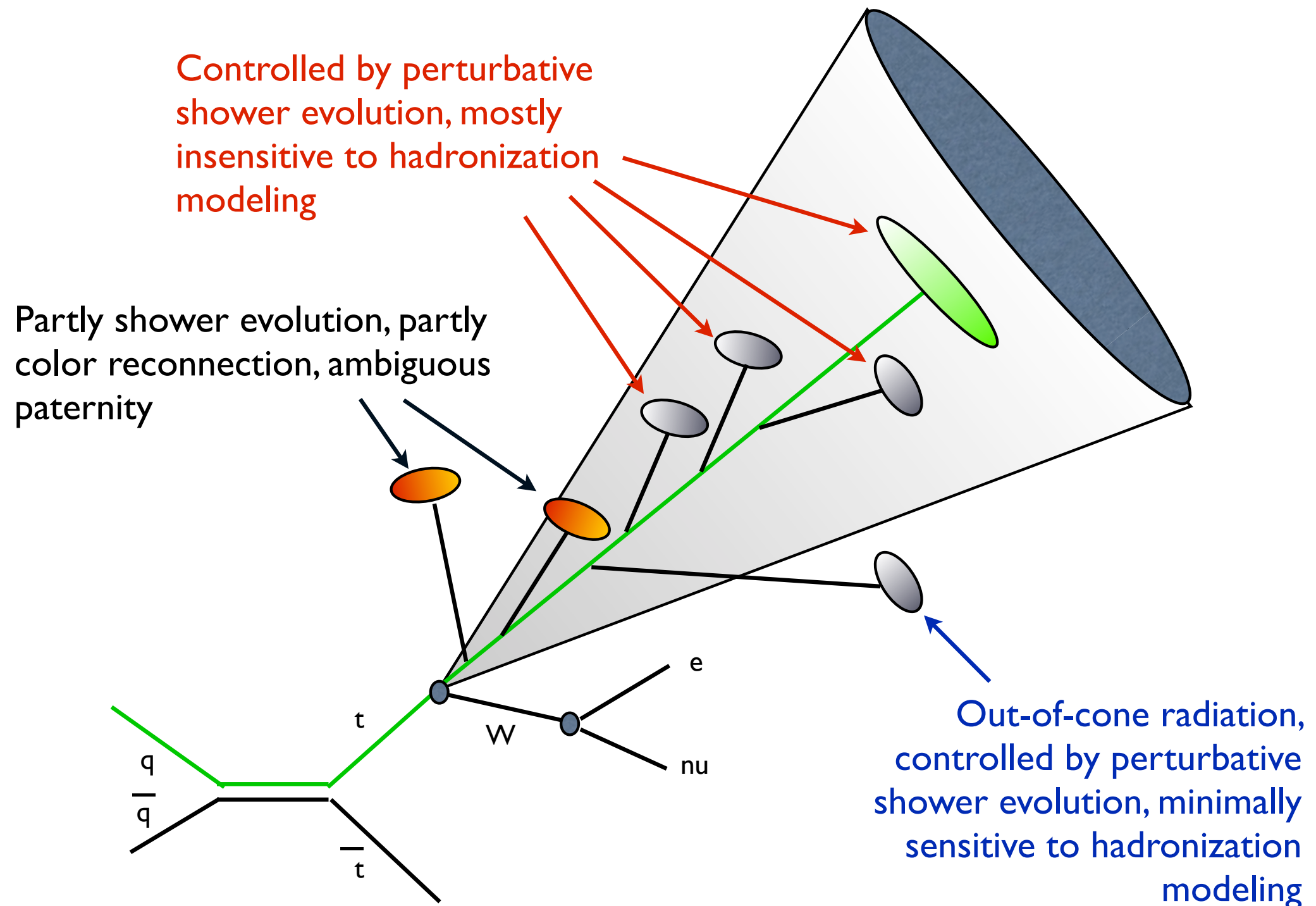
5. Formation of  
“odd” cluster

6. Decay of “odd” clusters, if  
large cluster mass, and  
decays to hadrons

- Study dependence of reconstructed mass on “odd” clusters

# Top mass & hadronization

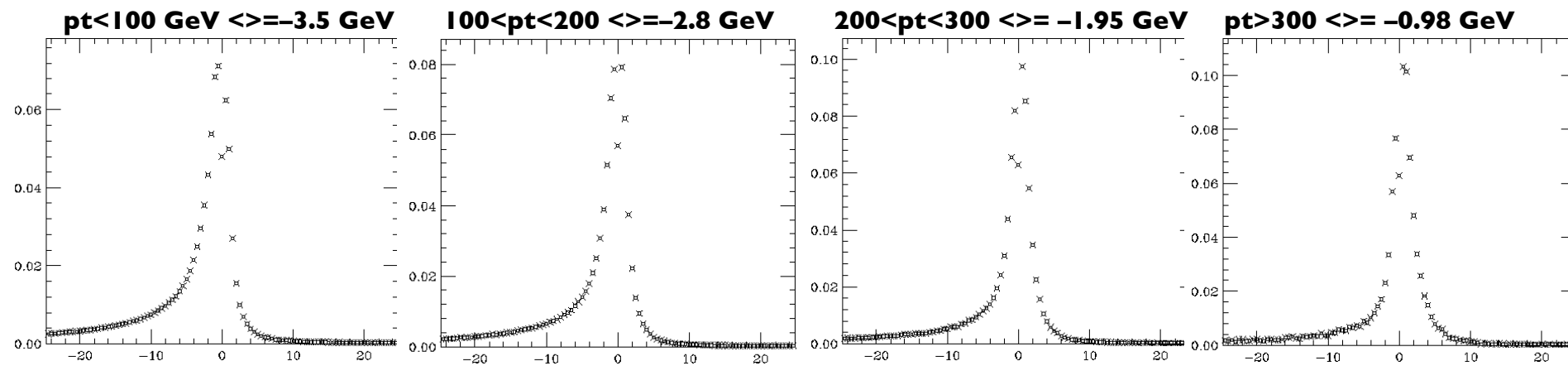
Mangano, Top LHC WG, July 2012



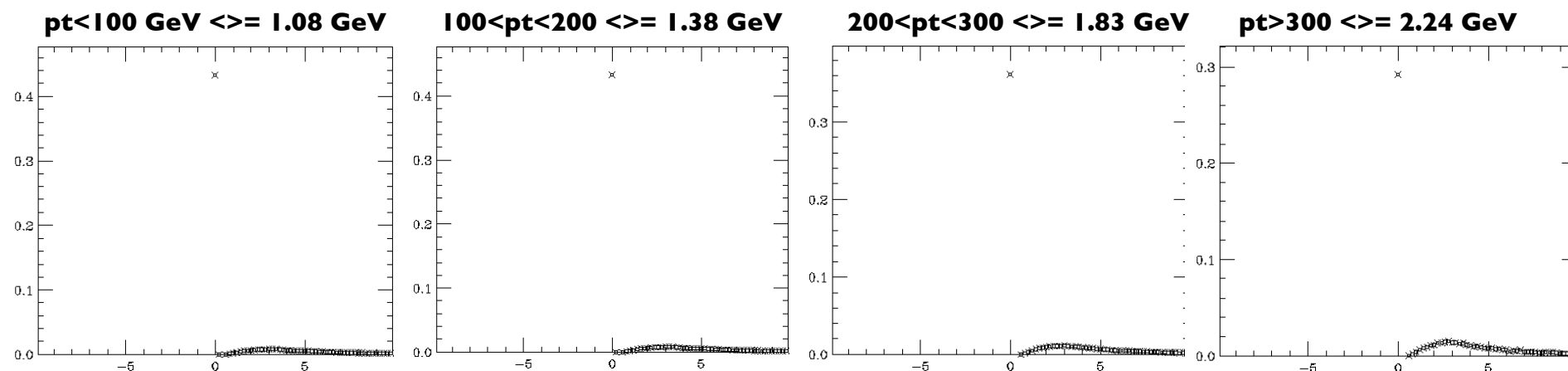
# Top mass & hadronization

**$m_{\text{top}}$  vs  $pt(\text{top})$**

**$m_{\text{top}}(\text{E+O}) - 172.5$**



**$m_{\text{top}}(\text{E+O}) - m_{\text{top}}(\text{E})$**



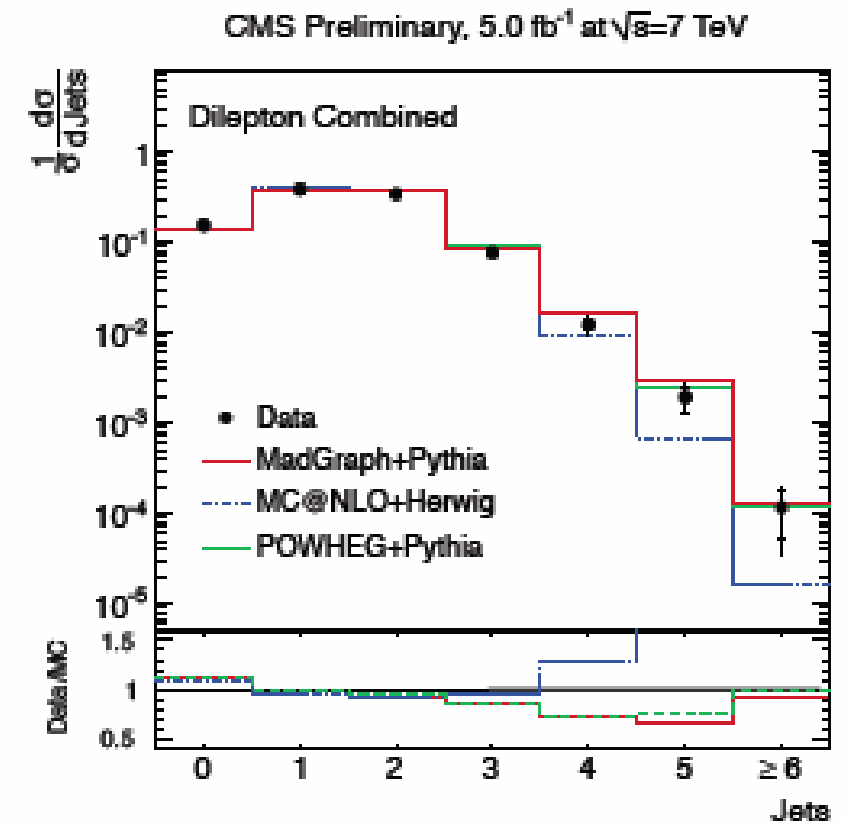
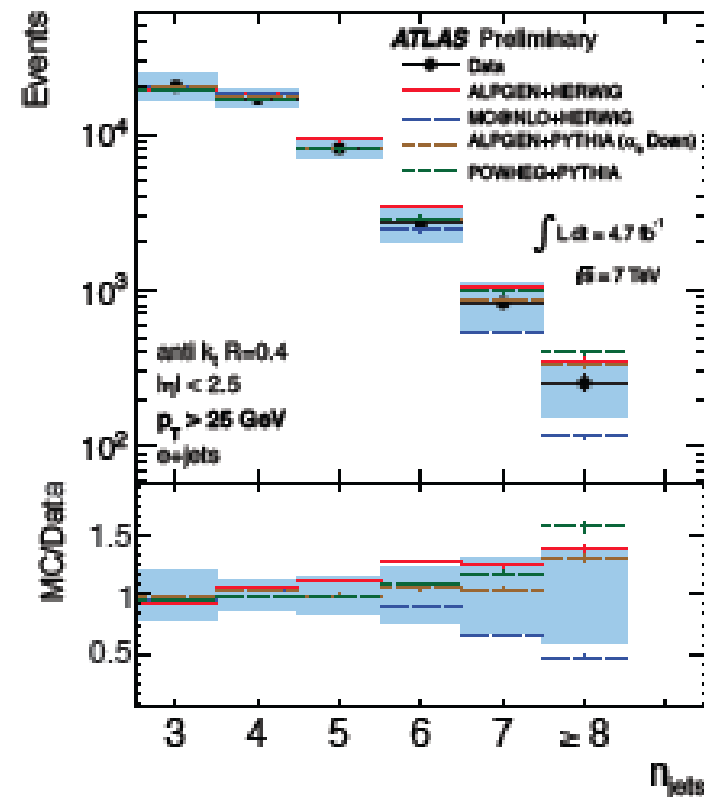
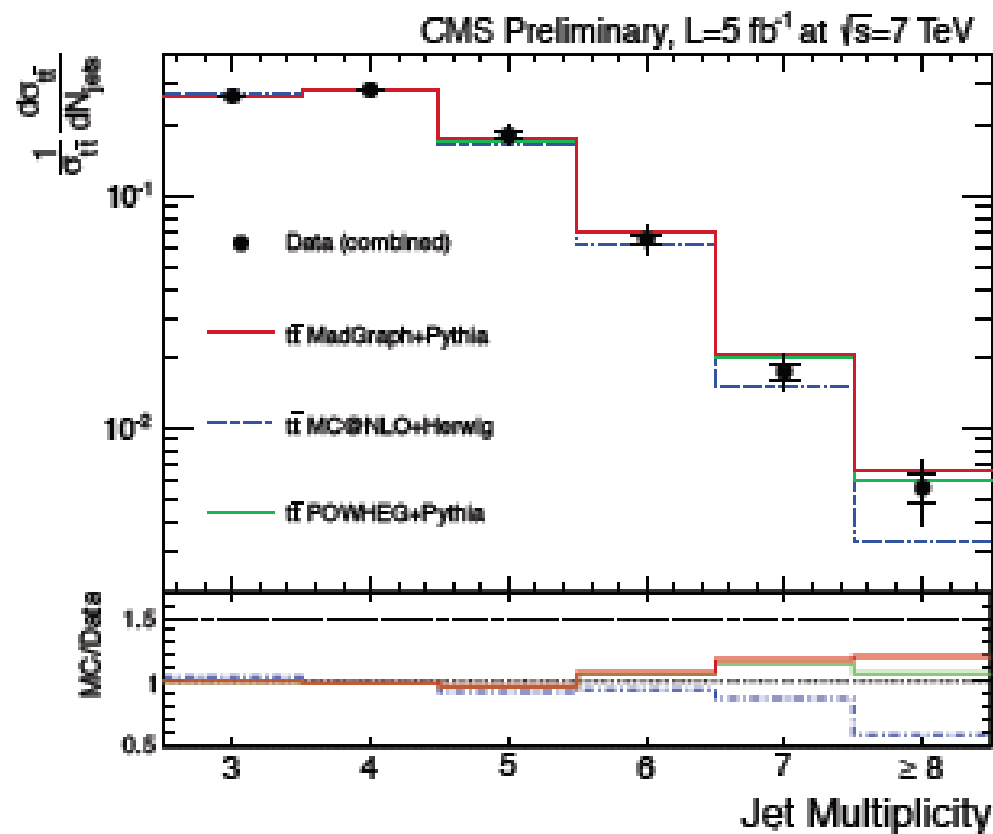
- Dependence of reconstructed mass on “odd” clusters  $\sim 1$  GeV

# Top+jets

**CMS PAS TOP-12-018** (l+jets)  
**ATLAS-CONF-2012-155** (l+jets)

$$\frac{1}{\sigma} \frac{d\sigma(N_{jets})}{dN_{jets}}$$

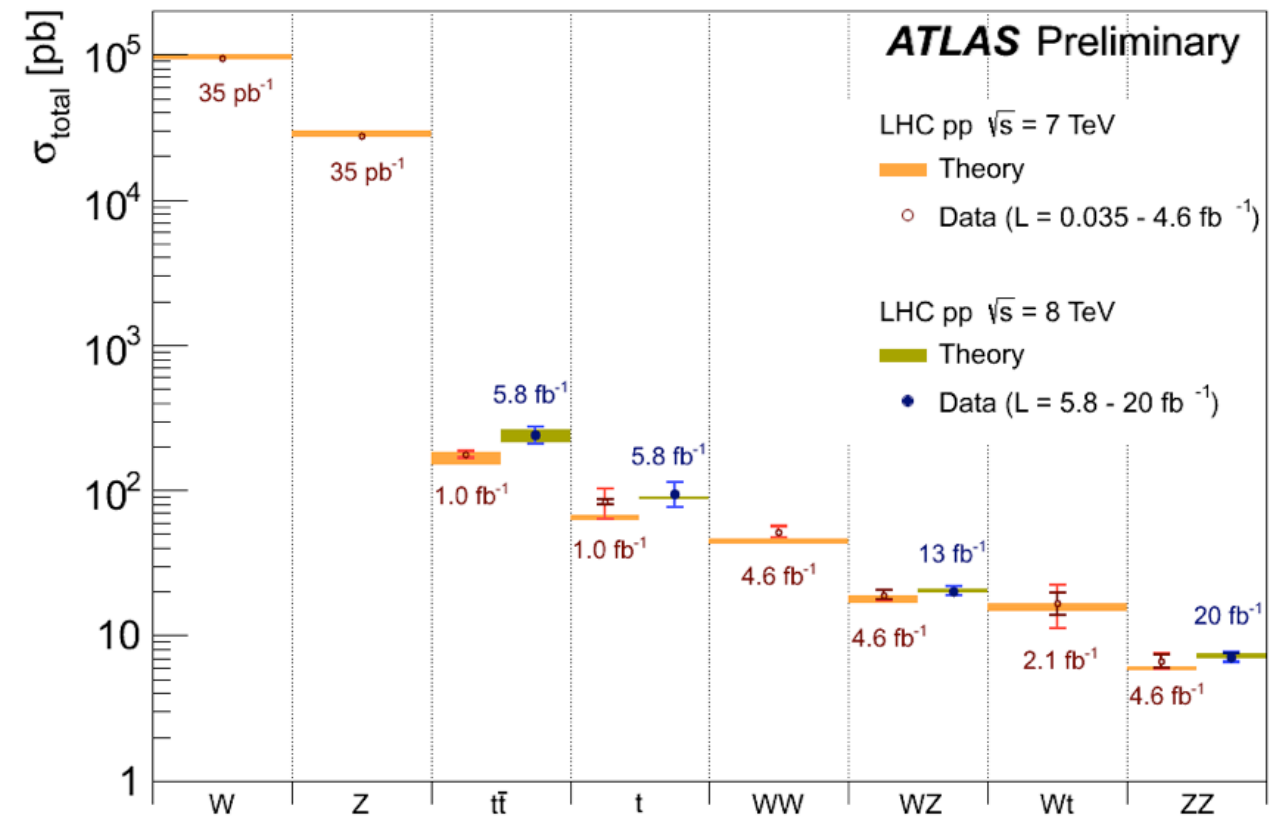
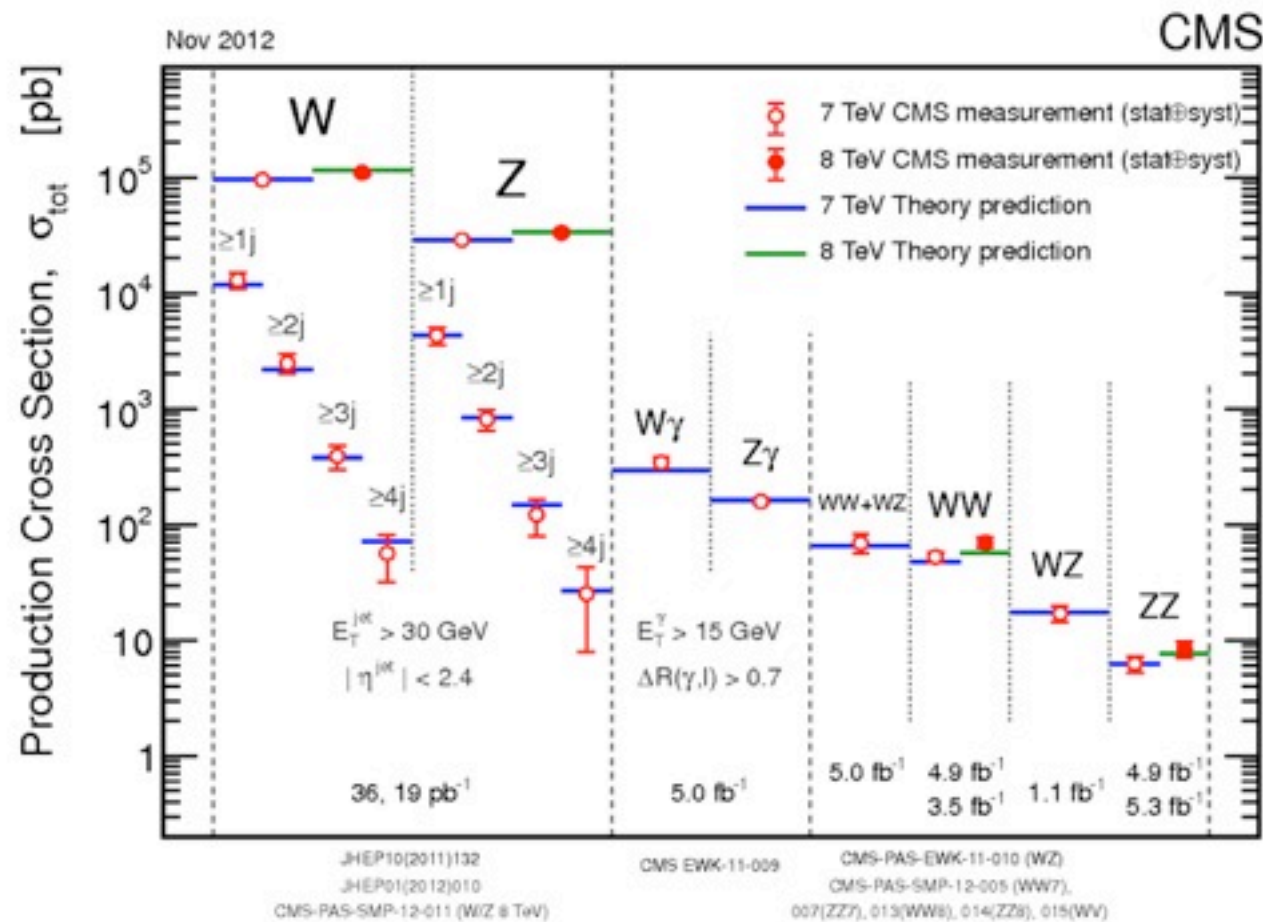
**CMS PAS TOP-12-023**  
 (dilepton)



- Matched NLO not adequate for >2 extra jets
- Merged multijets better there (for  $d\sigma/\sigma$ )



# LHC Cross Section Summary

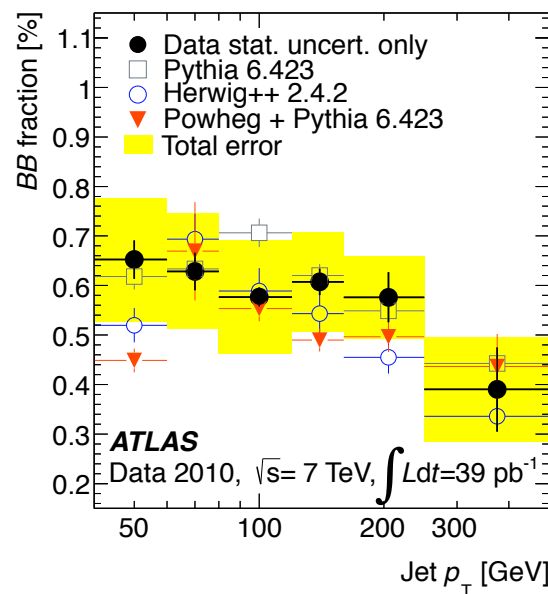


- Surprisingly good agreement
- No sign of non-Standard-Model phenomena (yet)

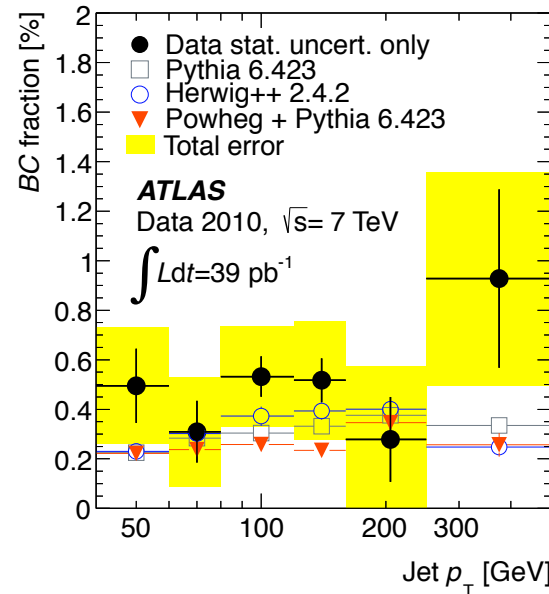
# But all is not perfect ...

- Dijet flavours versus jet  $p_T$

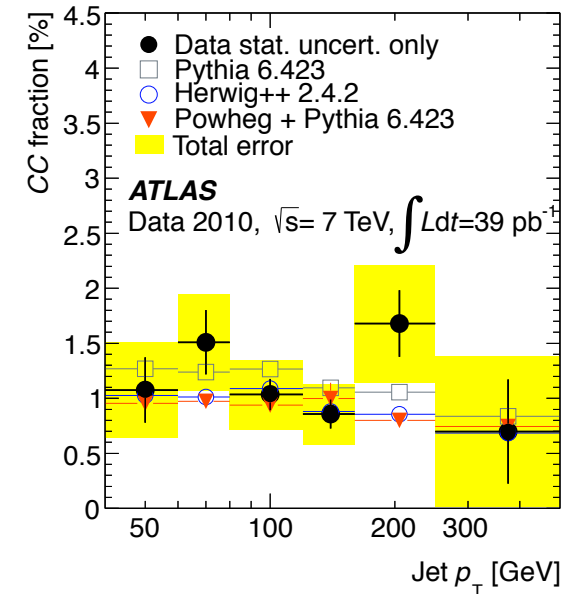
ATLAS, arXiv:1210.0441



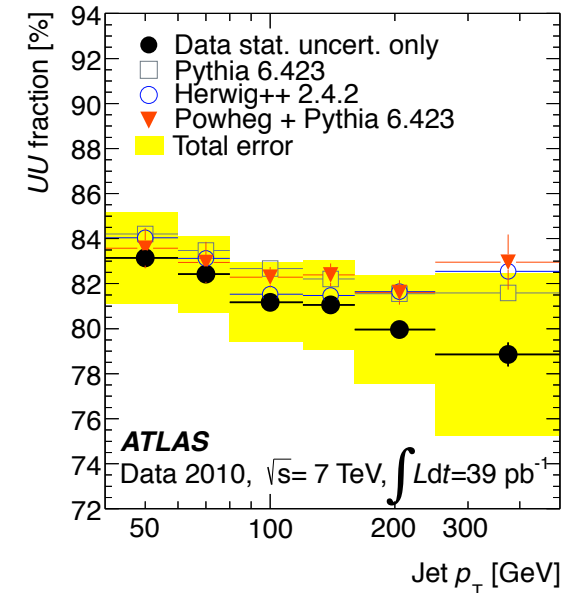
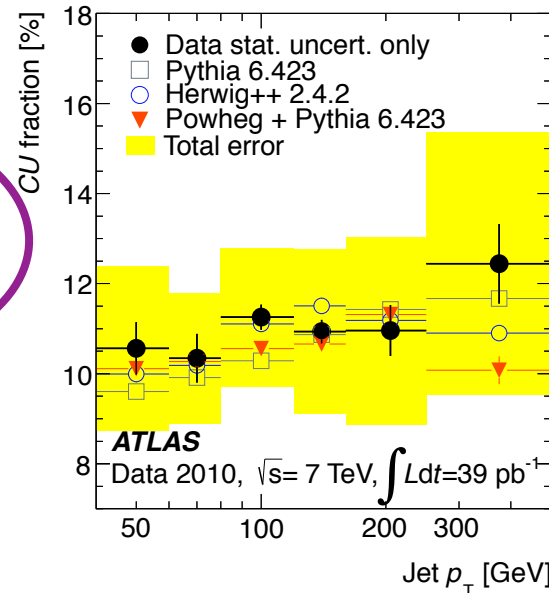
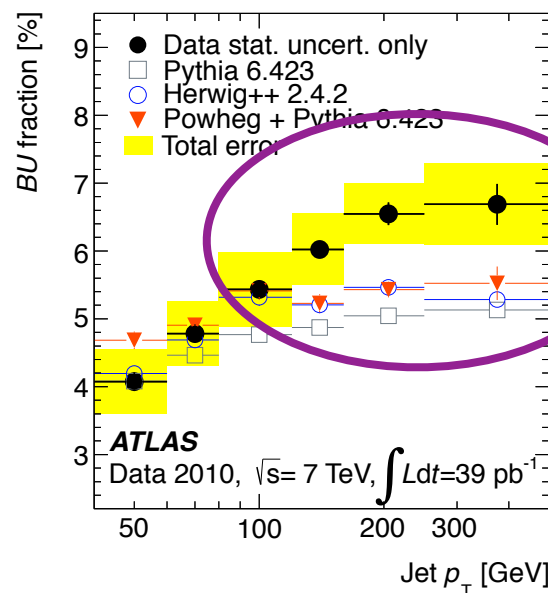
(a)



(b)



(c)



- Interesting excess of (single) b quark jets

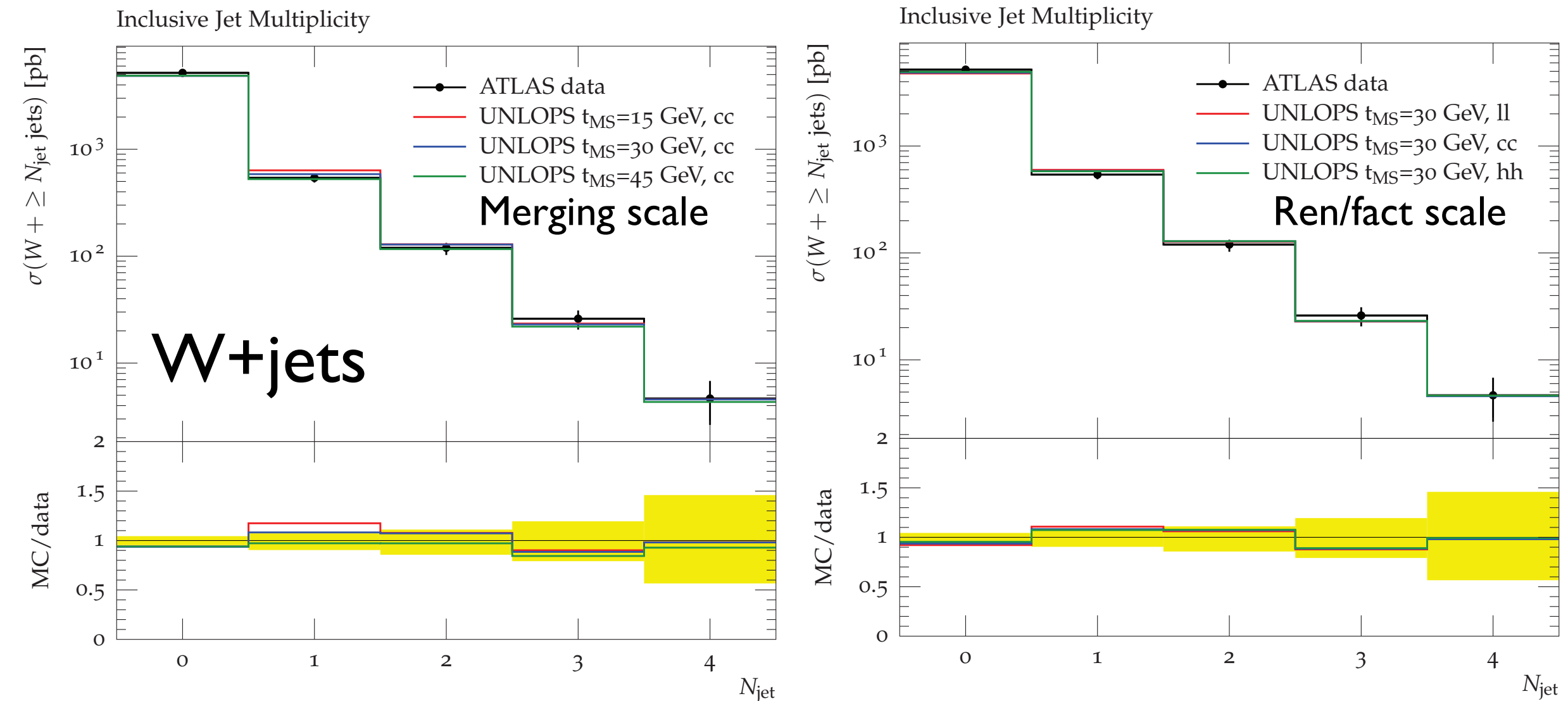
# Combined matching+merging

- NLO calculations generally refer to **inclusive** cross sections e.g.  $\sigma(W+\geq n \text{ jets})$
- Multijet merging does not preserve them, because of **mismatch** between exact real-emission and approximate (Sudakov) virtual corrections
- When correcting this mismatch, one can simultaneously upgrade them to NLO
- There remains the issue of merging scale dependence beyond NLO (large logs)

# Combined matching+merging

- Many competing schemes (pp, under development)
  - ✧ MEPS@NLO (SHERPA) Höche et al., arXiv:1207.5030
  - ✧ FxFx (aMC@NLO) Frederix & Frixione, arXiv:1209.6215
  - ✧ UNLOPS (Pythia 8) Lönnblad & Prestel, arXiv:1211.7278
  - ✧ MatchBox (Herwig++) Plätzer, arXiv:1211.5467
  - ✧ MiNLO (POWHEG) Hamilton et al., arXiv:1212.4504
  - ✧ GENEVA Alioli, Bauer et al., arXiv:1212.4504
- Some key ideas in LoopSim Rubin, Salam & Sapeta, JHEP1009, 084

# Combined matching+merging



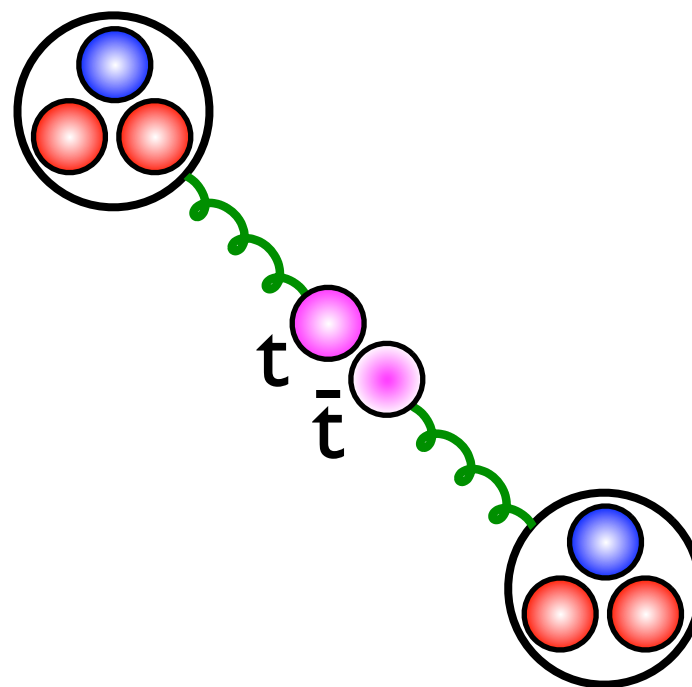
UNLOPS: Lönnblad & Prestel, arXiv:1211.7278

- Scale dependences almost eliminated

# Higgs boson production

# Higgs Production by Gluon Fusion

# Higgs Production by Gluon Fusion





# Higgs Production by Gluon Fusion

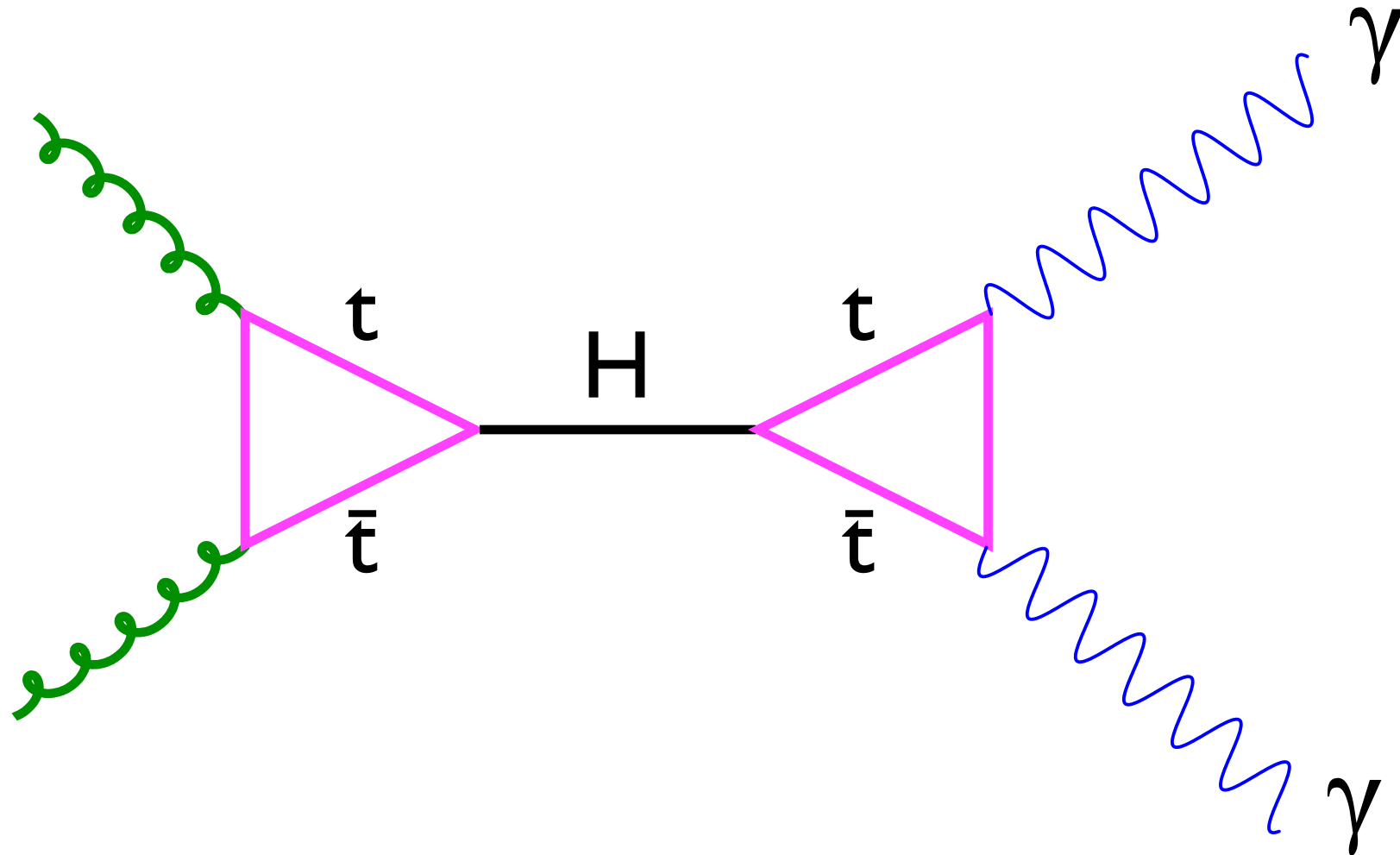


# Higgs Production by Gluon Fusion

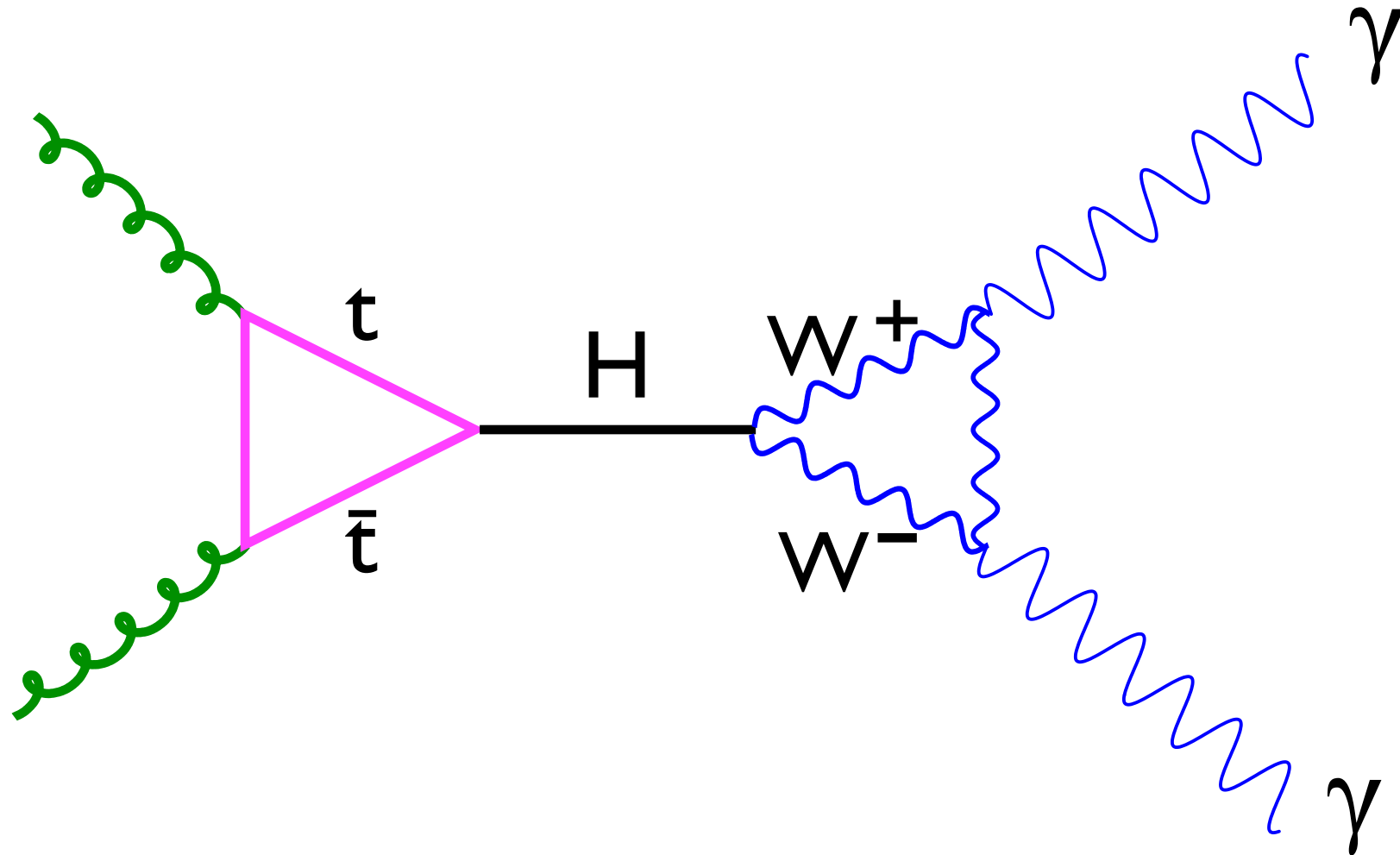


# Higgs Production by Gluon Fusion

# Higgs Production by Gluon Fusion

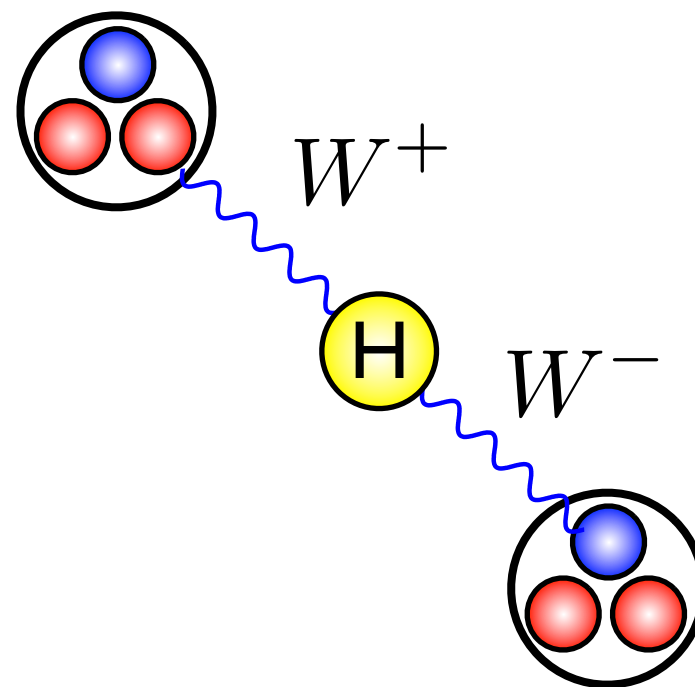


# Higgs Production by Gluon Fusion

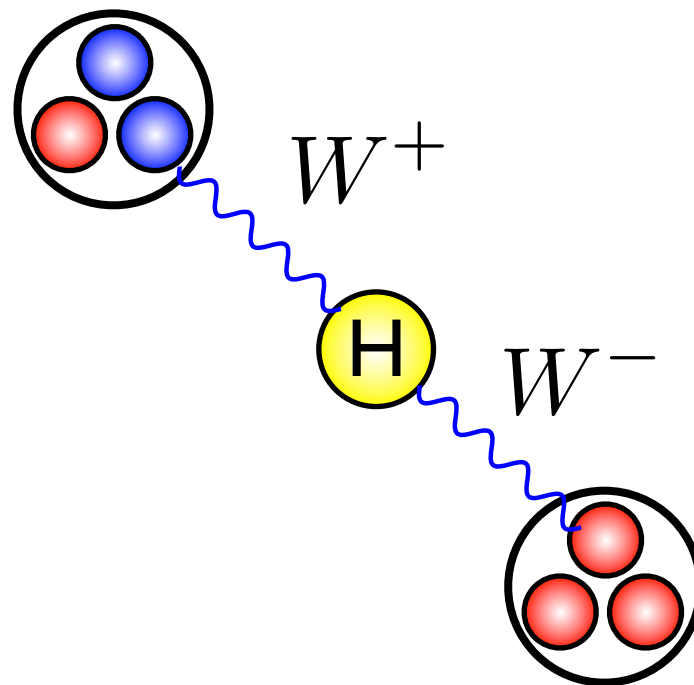


# Higgs Production by Vector Boson Fusion

# Higgs Production by Vector Boson Fusion



# Higgs Production by Vector Boson Fusion





# Higgs Production by Vector Boson Fusion



- Forward jets
- Few central jets
- Central jet veto increases S/B

# Higgs Signal and Background Simulation

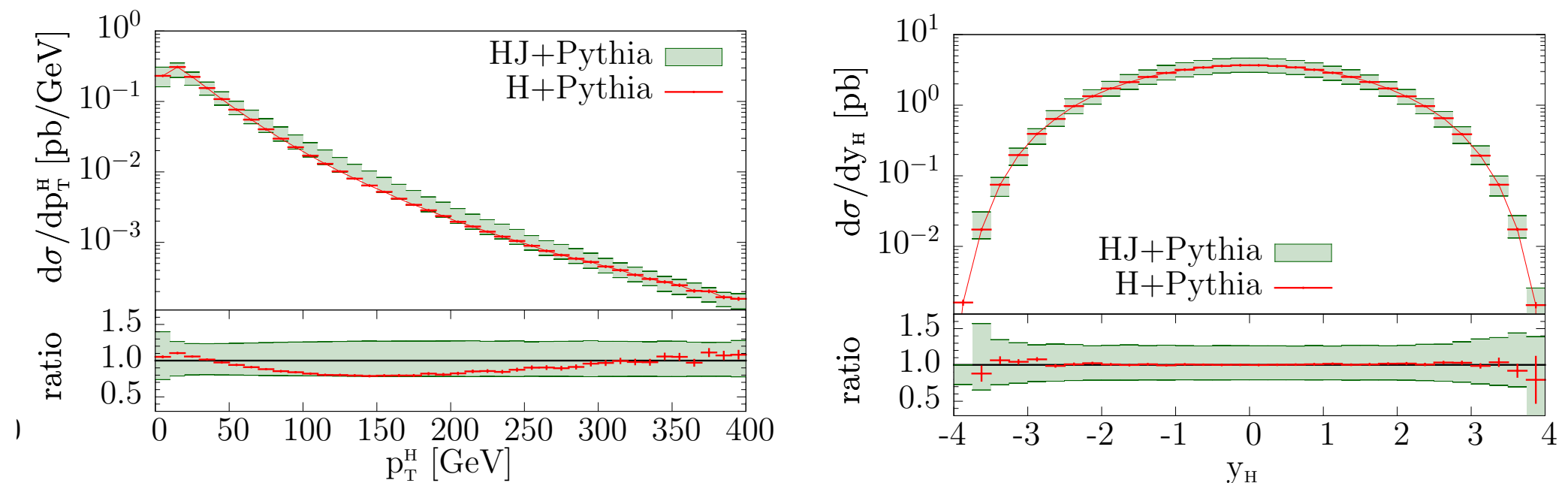
Process	Generator
ggF, VBF	POWHEG [57, 58]+PYTHIA
$WH, ZH, t\bar{t}H$	PYTHIA
$W$ +jets, $Z/\gamma^*$ +jets	ALPGEN [59]+HERWIG
$t\bar{t}, tW, tb$	MC@NLO [60]+HERWIG
$tqb$	AcerMC [61]+PYTHIA
$q\bar{q} \rightarrow WW$	MC@NLO+HERWIG
$gg \rightarrow WW$	gg2WW [62]+HERWIG
$q\bar{q} \rightarrow ZZ$	POWHEG [63]+PYTHIA
$gg \rightarrow ZZ$	gg2ZZ [64]+HERWIG
$WZ$	MadGraph+PYTHIA, HERWIG
$W\gamma$ +jets	ALPGEN+HERWIG
$W\gamma^*$ [65]	MadGraph+PYTHIA
$q\bar{q}/gg \rightarrow \gamma\gamma$	SHERPA

ATLAS, Phys.Lett.B716(2012)1

# gg → Higgs(+jet)

Higgs boson production total cross sections in pb at the LHC, 8 TeV							
$K_R, K_F$	1, 1	1, 2	2, 1	$1, \frac{1}{2}$	$\frac{1}{2}, 1$	$\frac{1}{2}, \frac{1}{2}$	2, 2
HJ-MiNLO NLO	13.33(3)	13.49(3)	<b>11.70(2)</b>	13.03(3)	<b>16.53(7)</b>	16.45(8)	11.86(2)
H NLO	13.23(1)	13.28(1)	<b>11.17(1)</b>	13.14(1)	<b>15.91(2)</b>	15.83(2)	11.22(1)
HJ-MiNLO LO	8.282(7)	8.400(7)	<b>5.880(5)</b>	7.864(6)	<b>18.28(2)</b>	17.11(2)	5.982(5)
H LO	5.741(5)	5.758(5)	<b>4.734(4)</b>	5.644(5)	<b>7.117(6)</b>	6.996(6)	4.748(4)

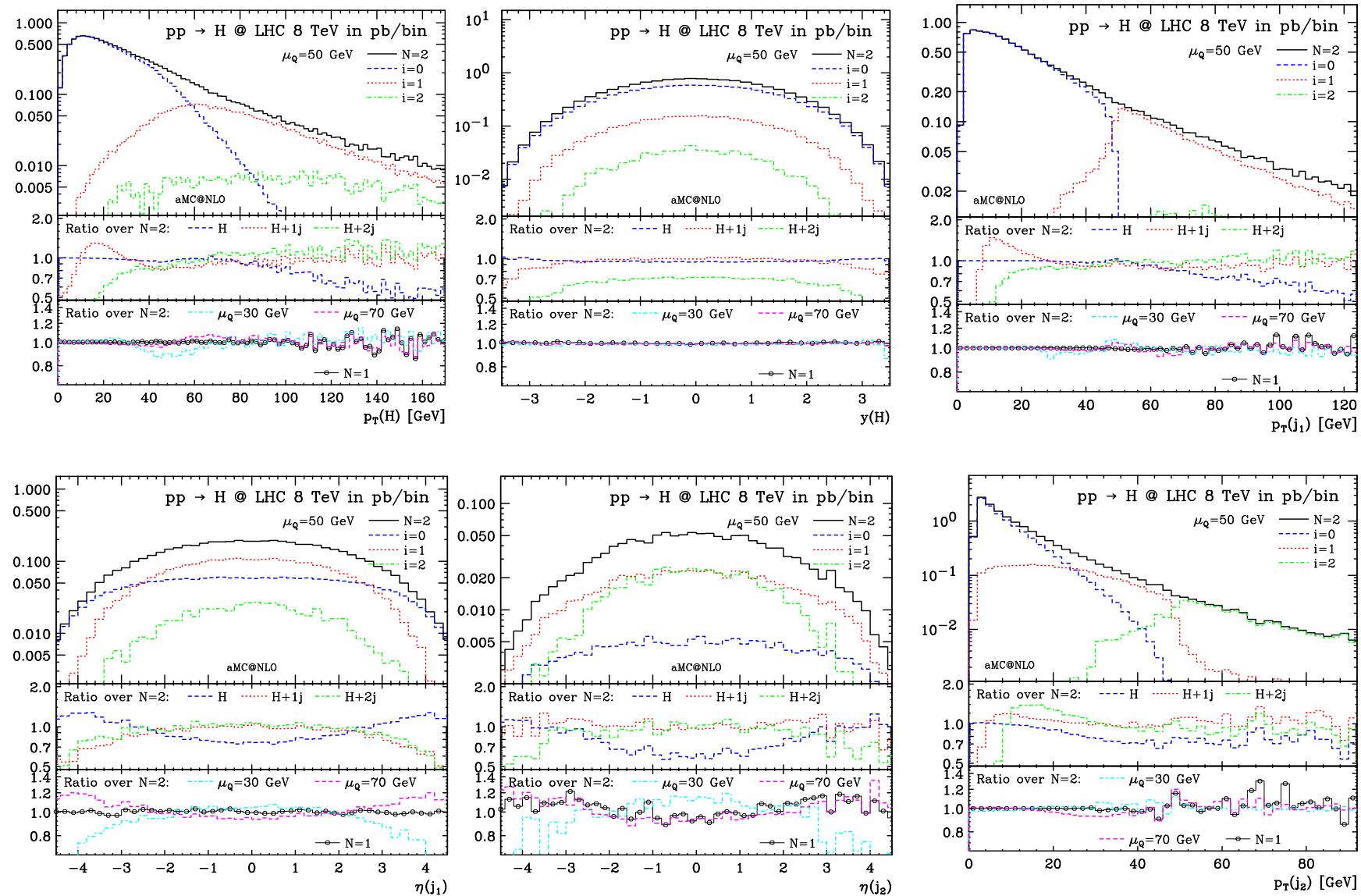
**Table 1:** Total cross section for Higgs boson production at the 8 TeV LHC, obtained with the HJ-MiNLO and the H programs, both at full NLO level and at leading order, for different scales combinations. The maximum and minimum are highlighted.



## ● Match/merge MiNLO+Pythia6

Hamilton, Nason, Oleari &  
Zanderighi, arXiv:1212.4504

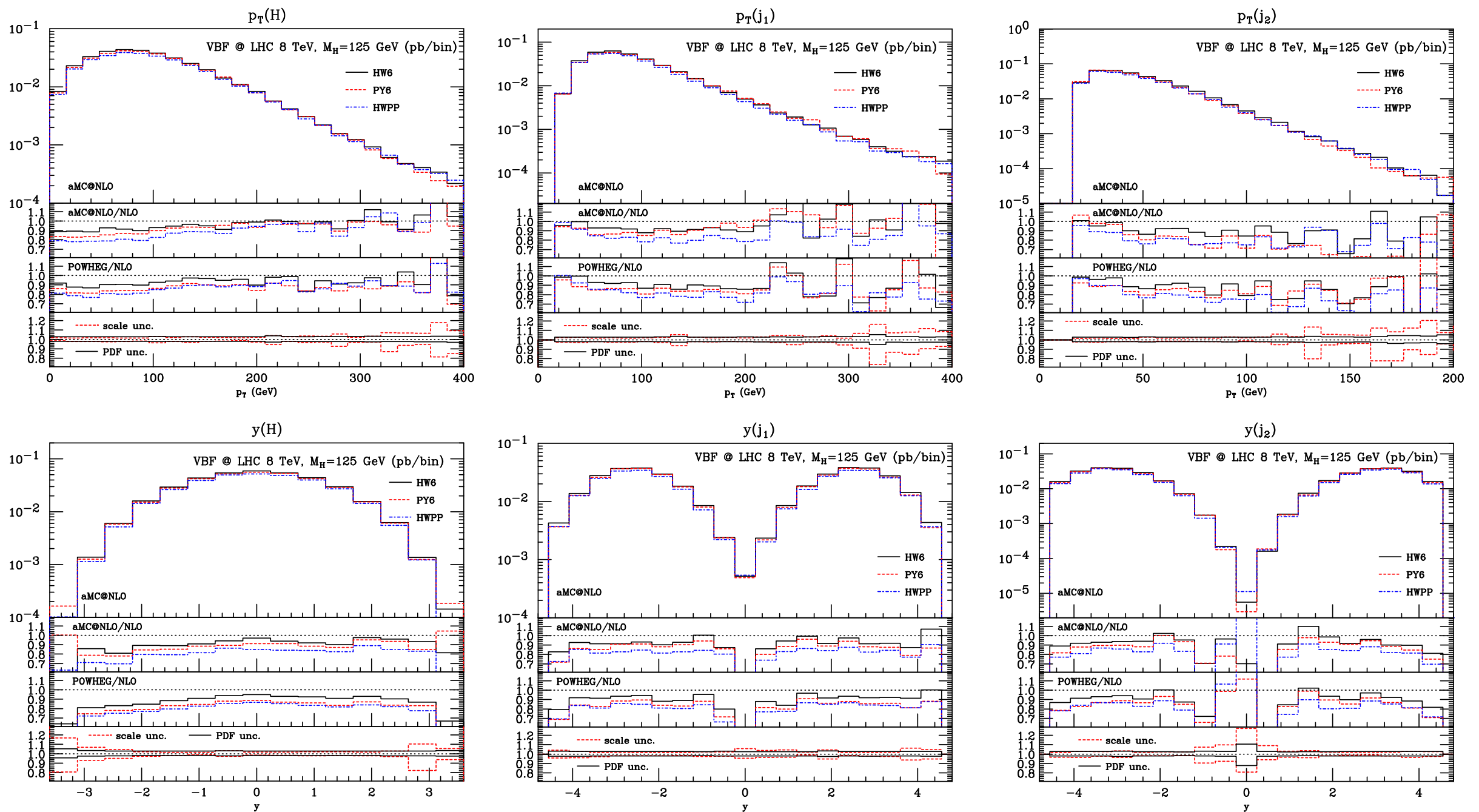
# $gg \rightarrow \text{Higgs} + \text{jets}$



- FxFx: Match/merge MC@NLO+Herwig6

Frederix & Frixione, arXiv:1209.6215

# VBF Higgs+jets



- Matched MC@NLO and POWHEG  
Frixione, Torrielli, Zaro, arXiv:1304.7927

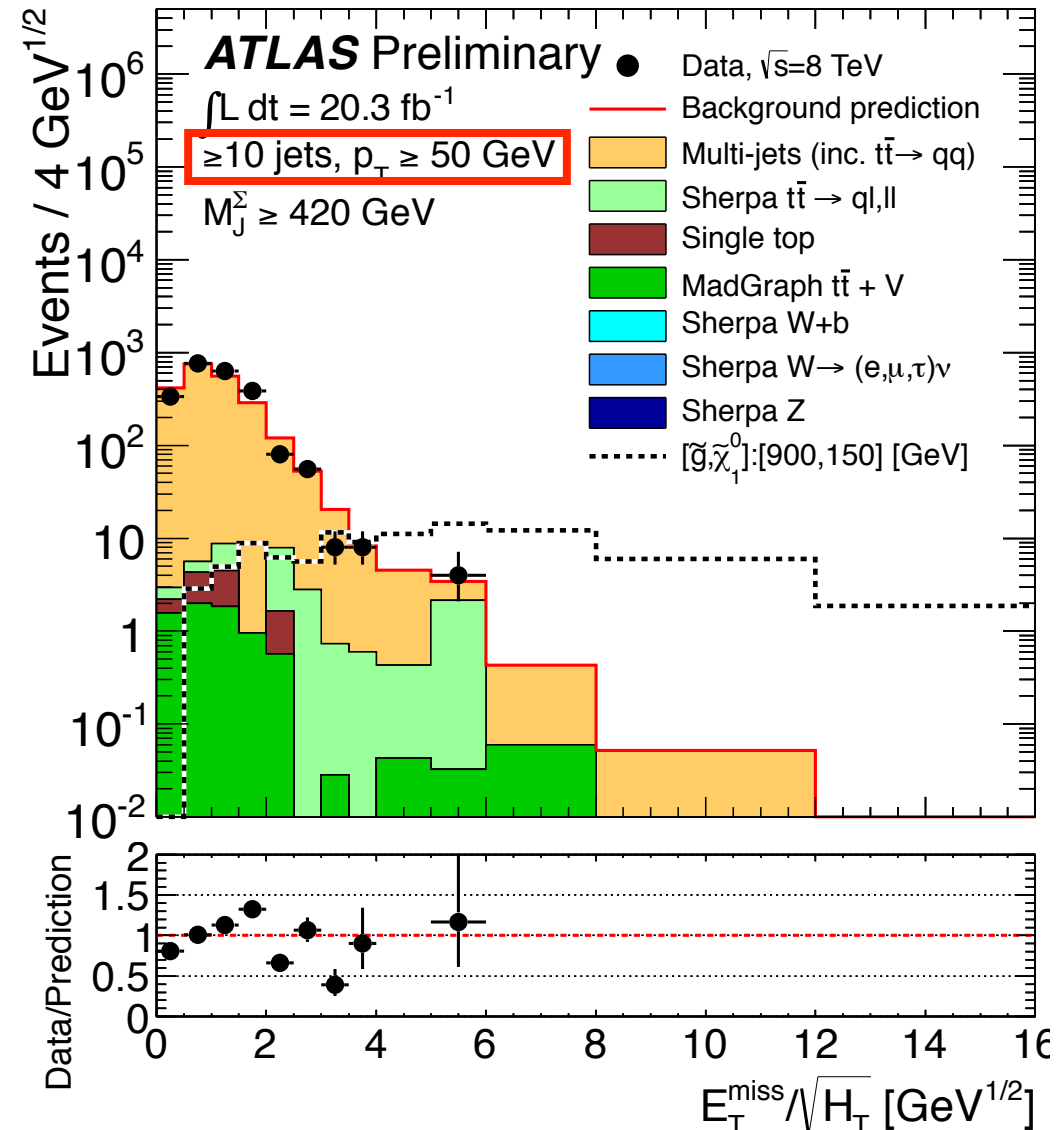
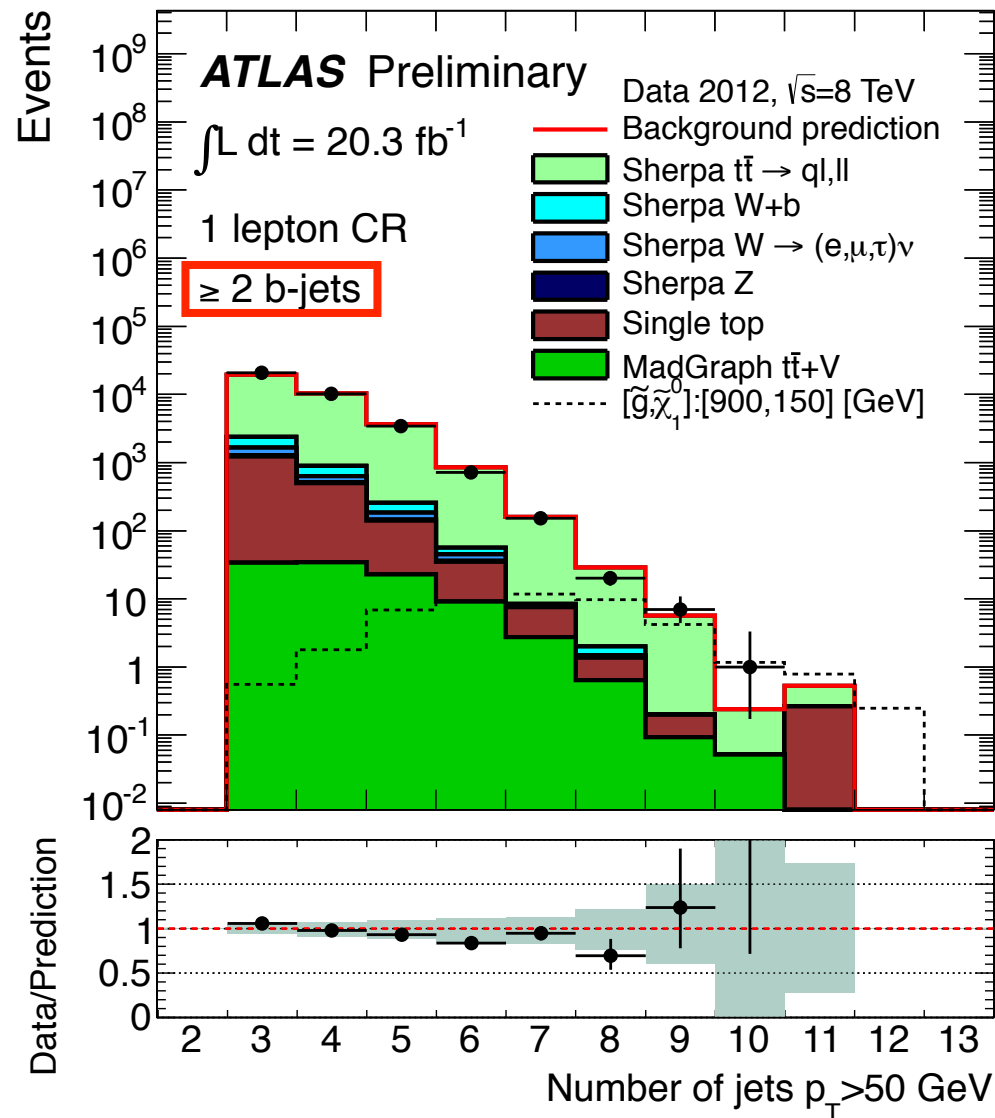
# Beyond Standard Model Simulation

# BSM Simulation

- Main generators have some BSM models built in
  - ✿ Pythia 6 has the most models
  - ✿ Herwig++ has careful treatment of SUSY spin correlations and off-shell effects
- Trend is now towards external matrix element generators: FeynRules + MadGraph, ...
- QCD corrections and matching/merging still needed

# Searching for new signals

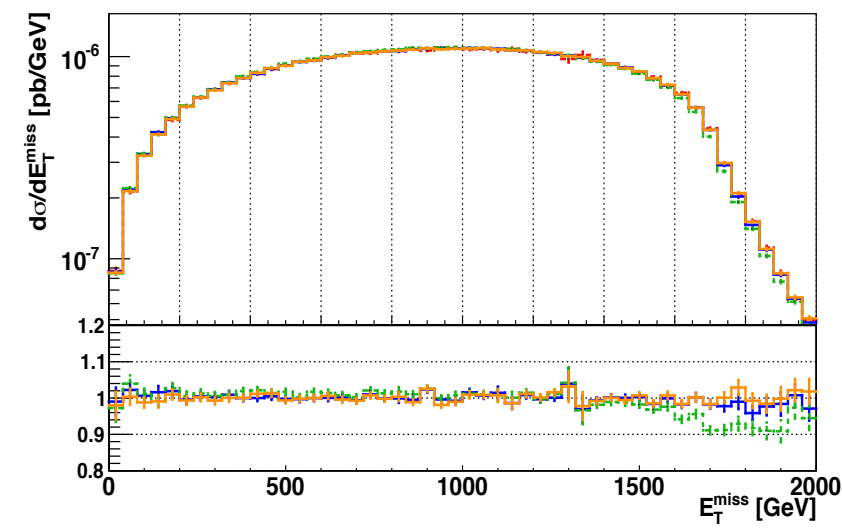
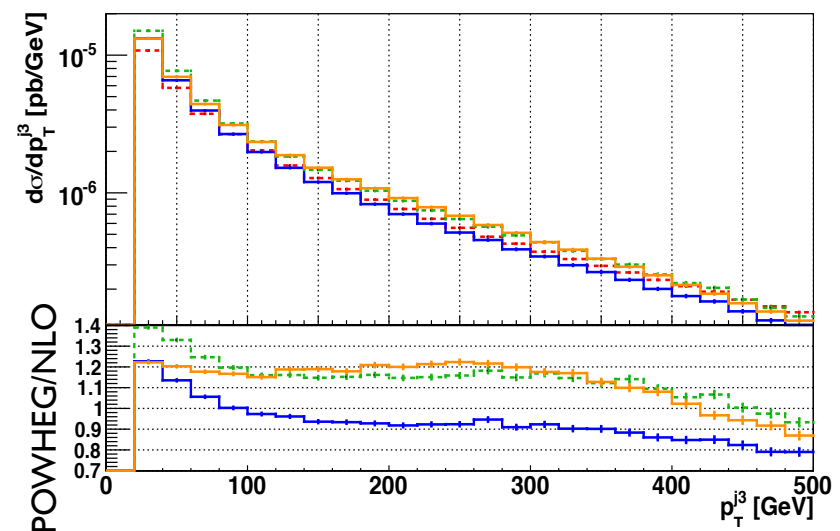
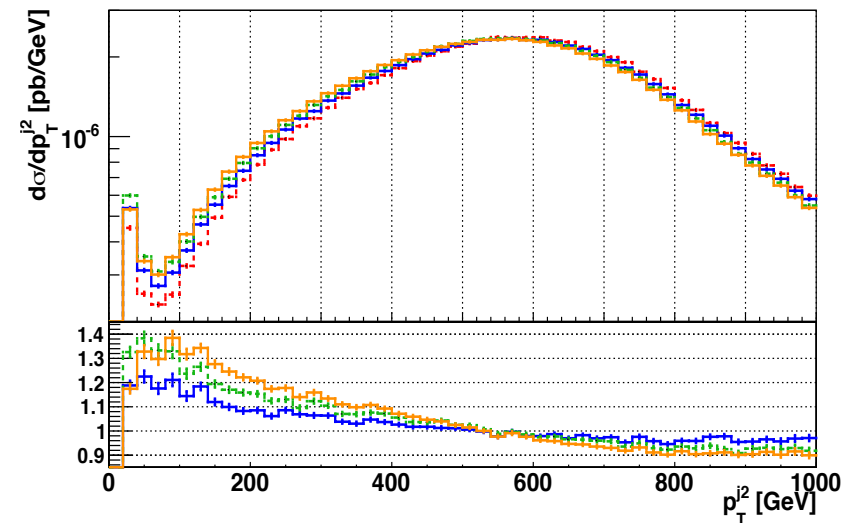
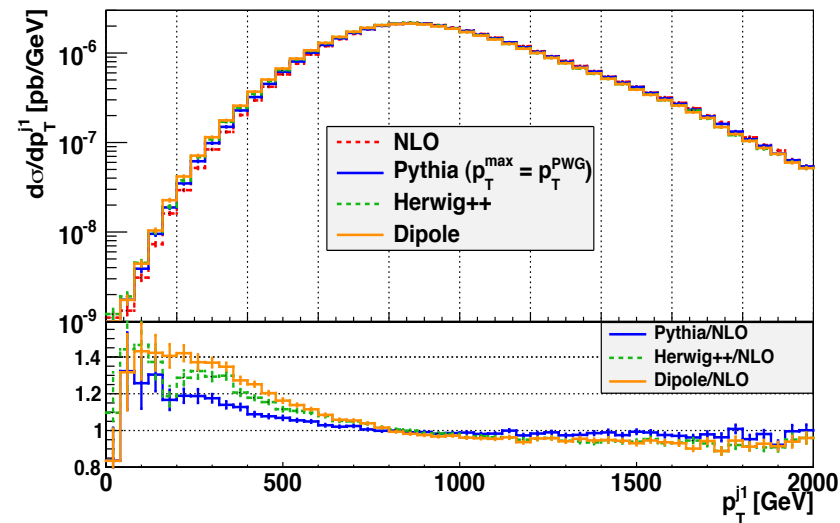
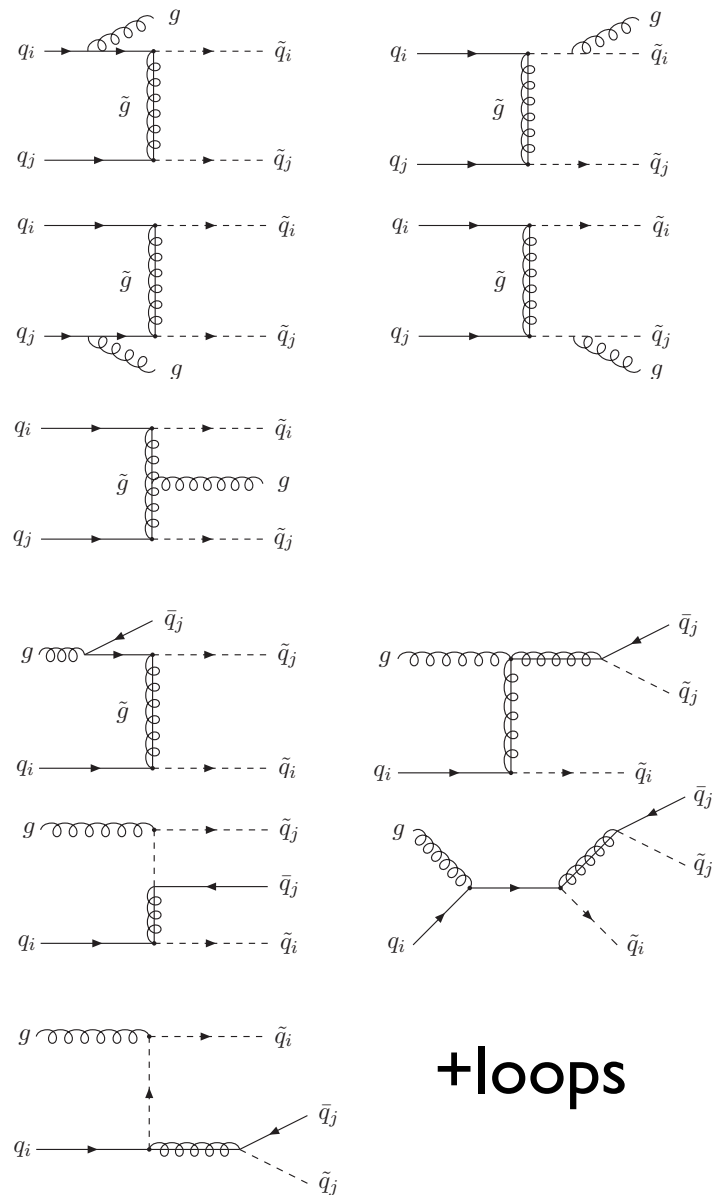
ATLAS CONF-2013-054



- Dashed = Herwig++  $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t + \bar{t} + \tilde{\chi}_1^0$
- Background: mostly Sherpa LO multijet merging



# NLO Squark Production

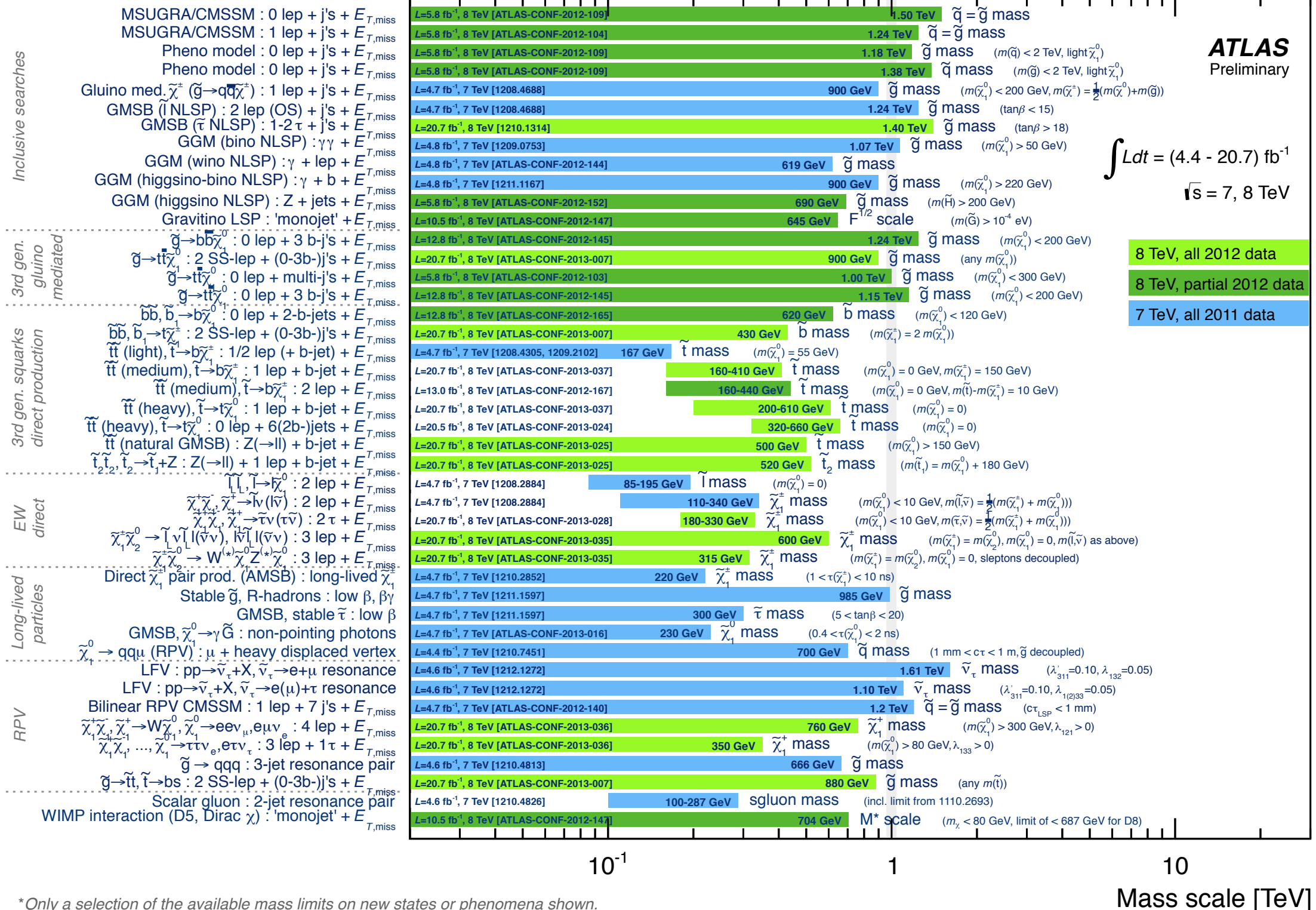


- NLO with POWHEG matching to different generators

Gavin et al., arXiv:1305.4061

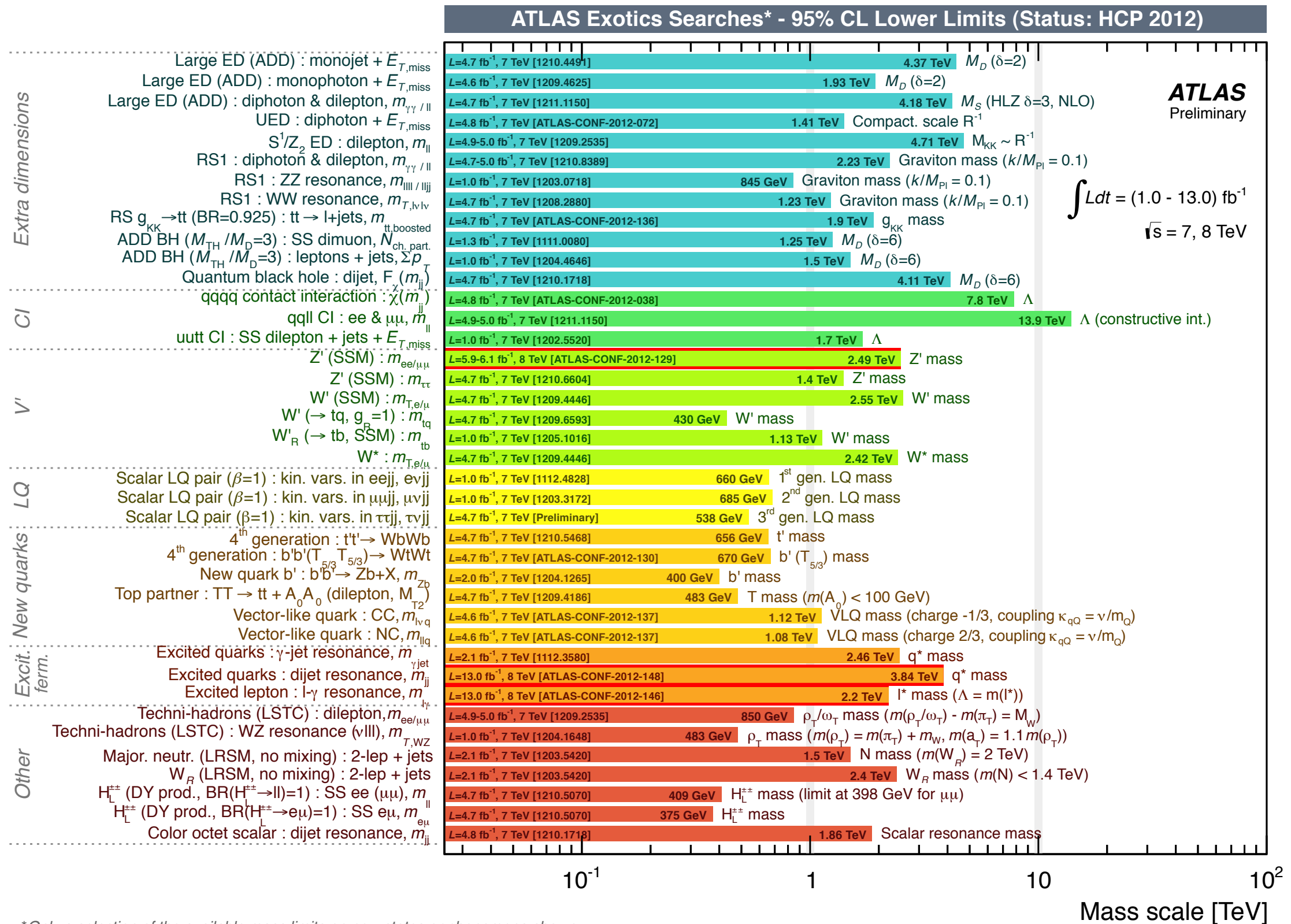
# ATLAS SUSY Search

ATLAS SUSY Searches\* - 95% CL Lower Limits (Status: March 26, 2013)



\*Only a selection of the available mass limits on new states or phenomena shown.  
 All limits quoted are observed minus  $1\sigma$  theoretical signal cross section uncertainty.

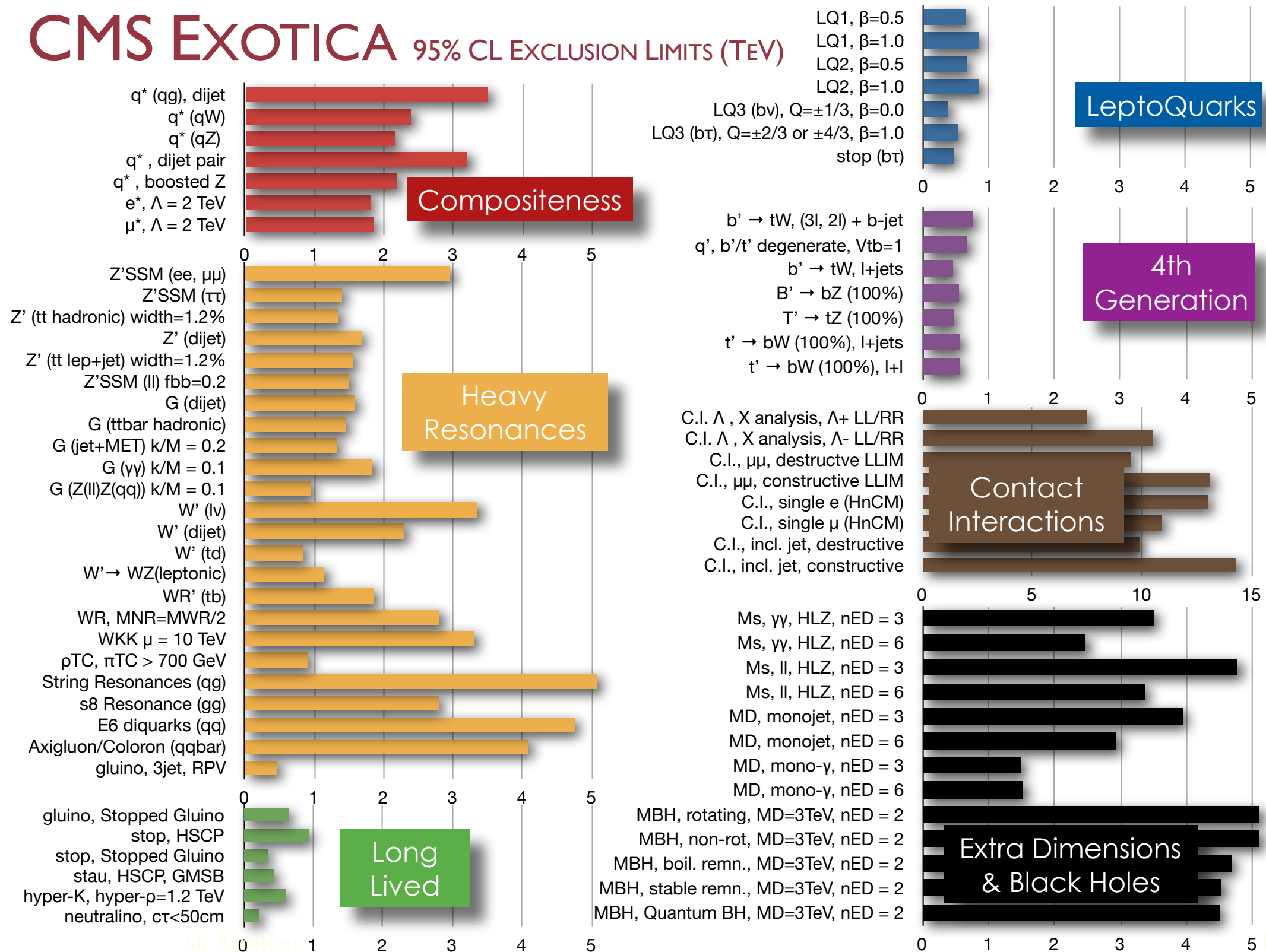
# ATLAS Exotica Search



\*Only a selection of the available mass limits on new states or phenomena shown

# CMS Exotica Search

## CMS EXOTICA 95% CL EXCLUSION LIMITS (TeV)





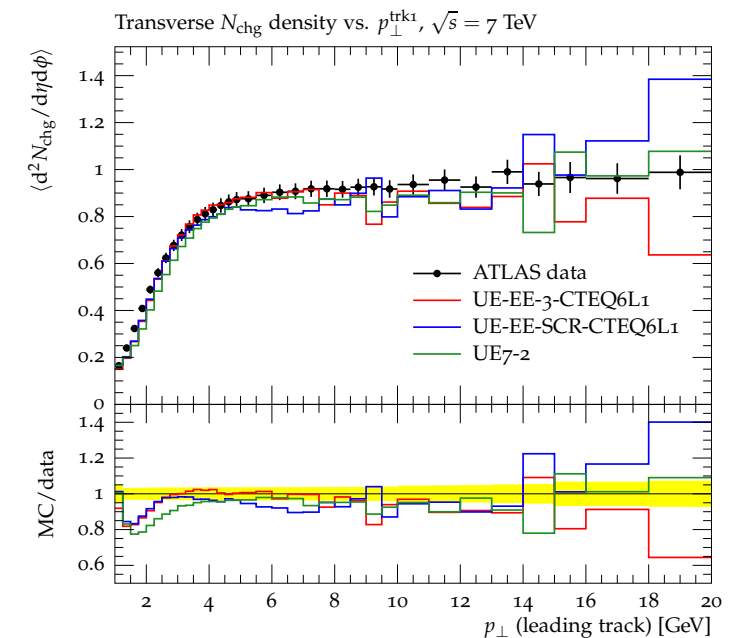
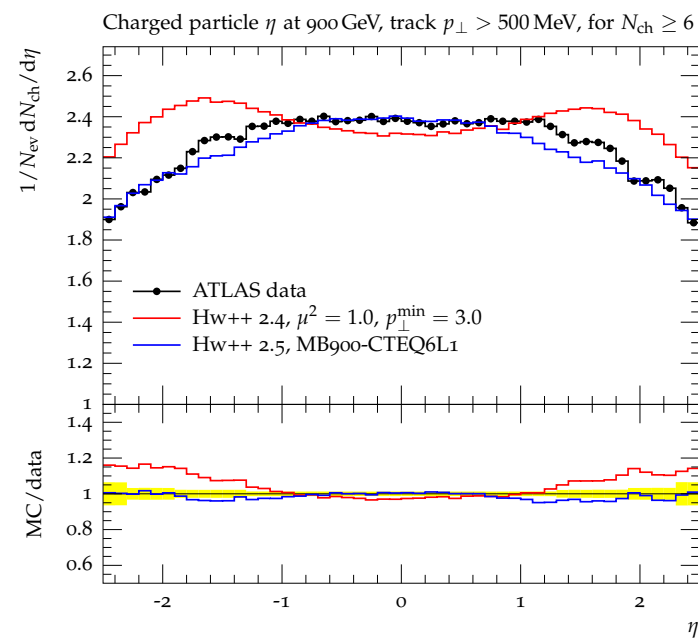
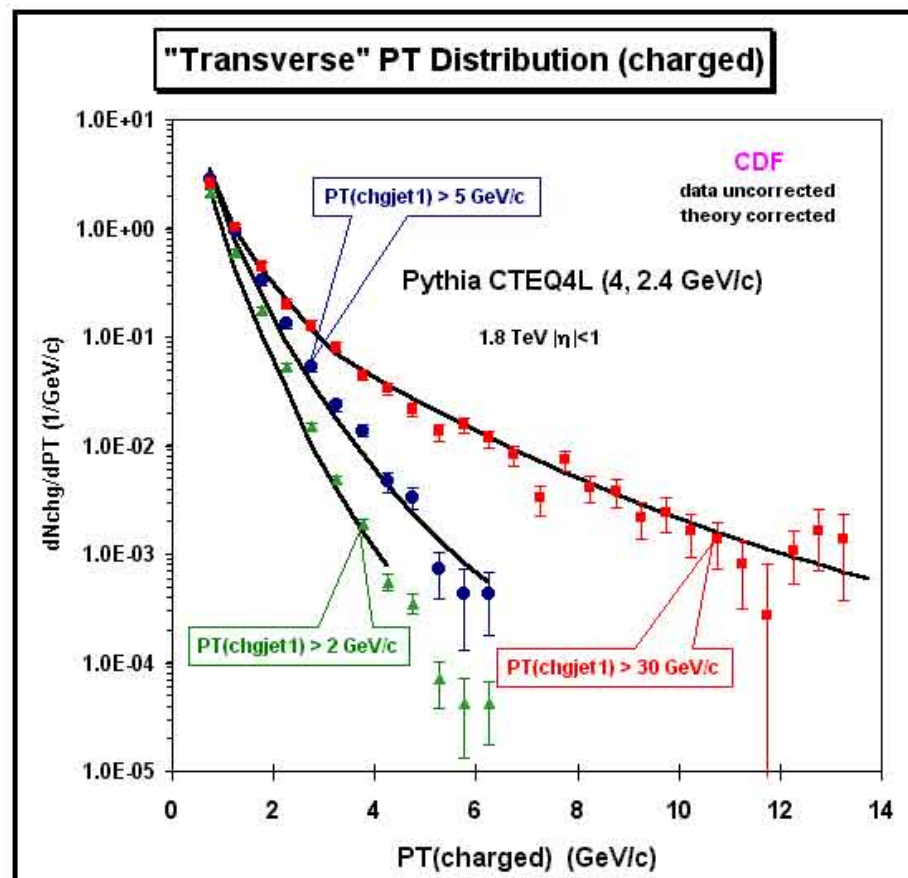
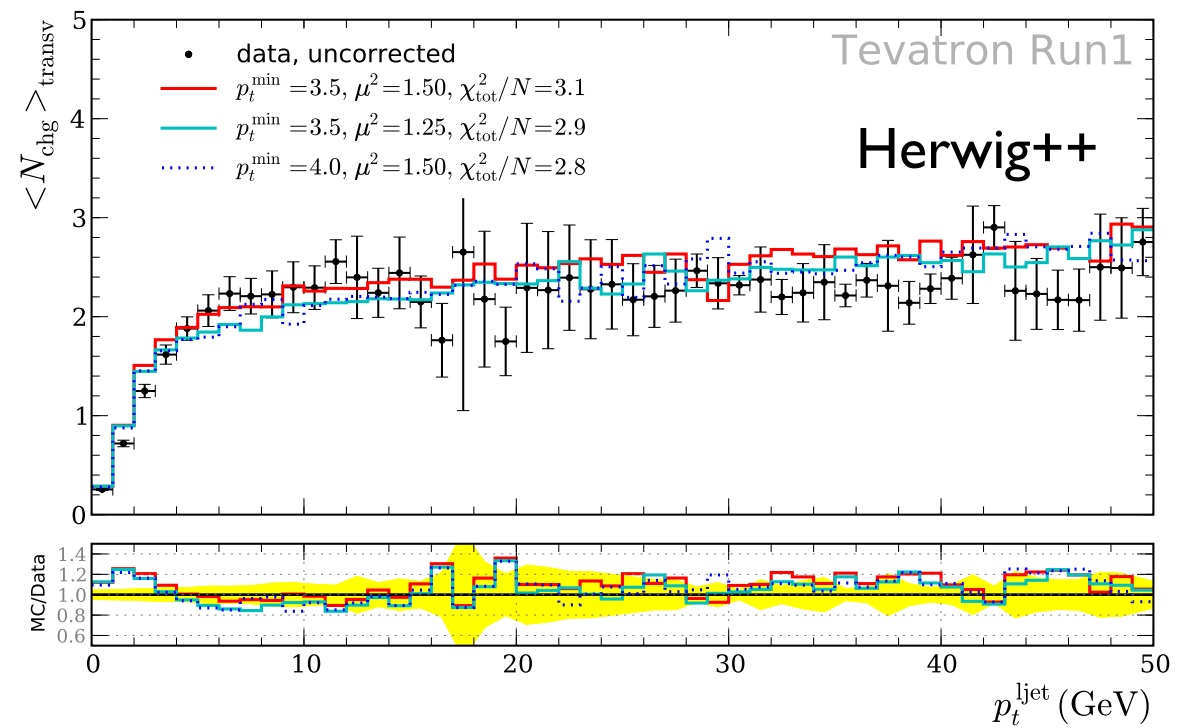
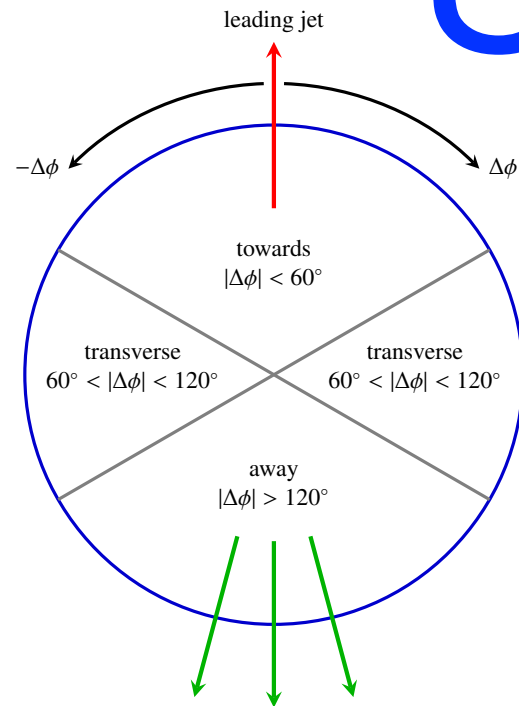
# Conclusions and Prospects

- Standard Model has (so far) been spectacularly confirmed at the LHC
- Monte Carlo event generation of (SM and BSM) signals and backgrounds plays a big part
- Matched NLO and merged multi-jet generators have proved essential
  - ✦ Automation and NLO merging in progress
  - ✦ NNLO much more challenging
- Still plenty of scope for new discoveries!

# Thanks for listening!

# Backup

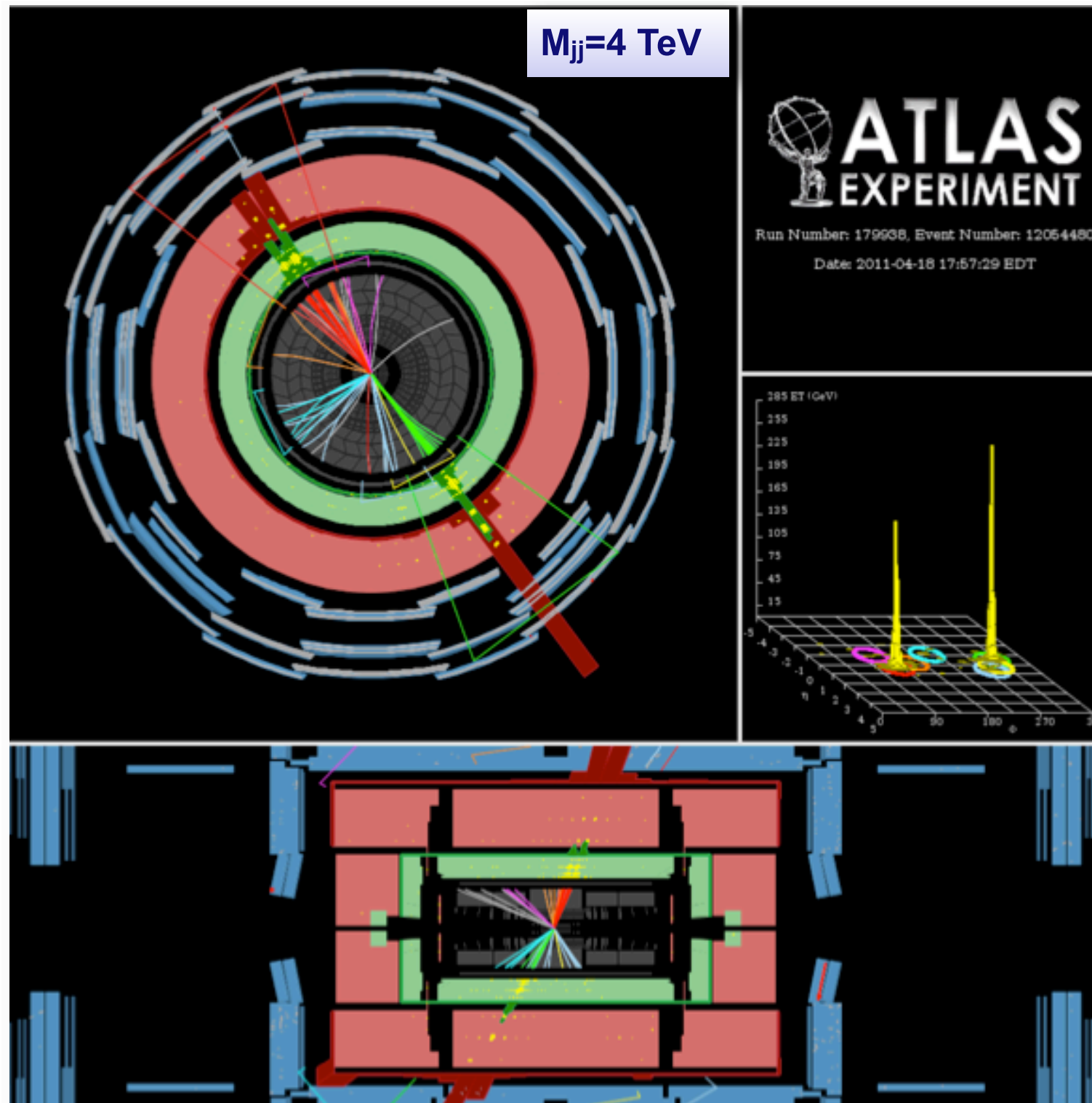
# Underlying Event



ATLAS PRD83(2011)112001  
Gieseke, Röhr, Siódmok, arXiv:1206.2205

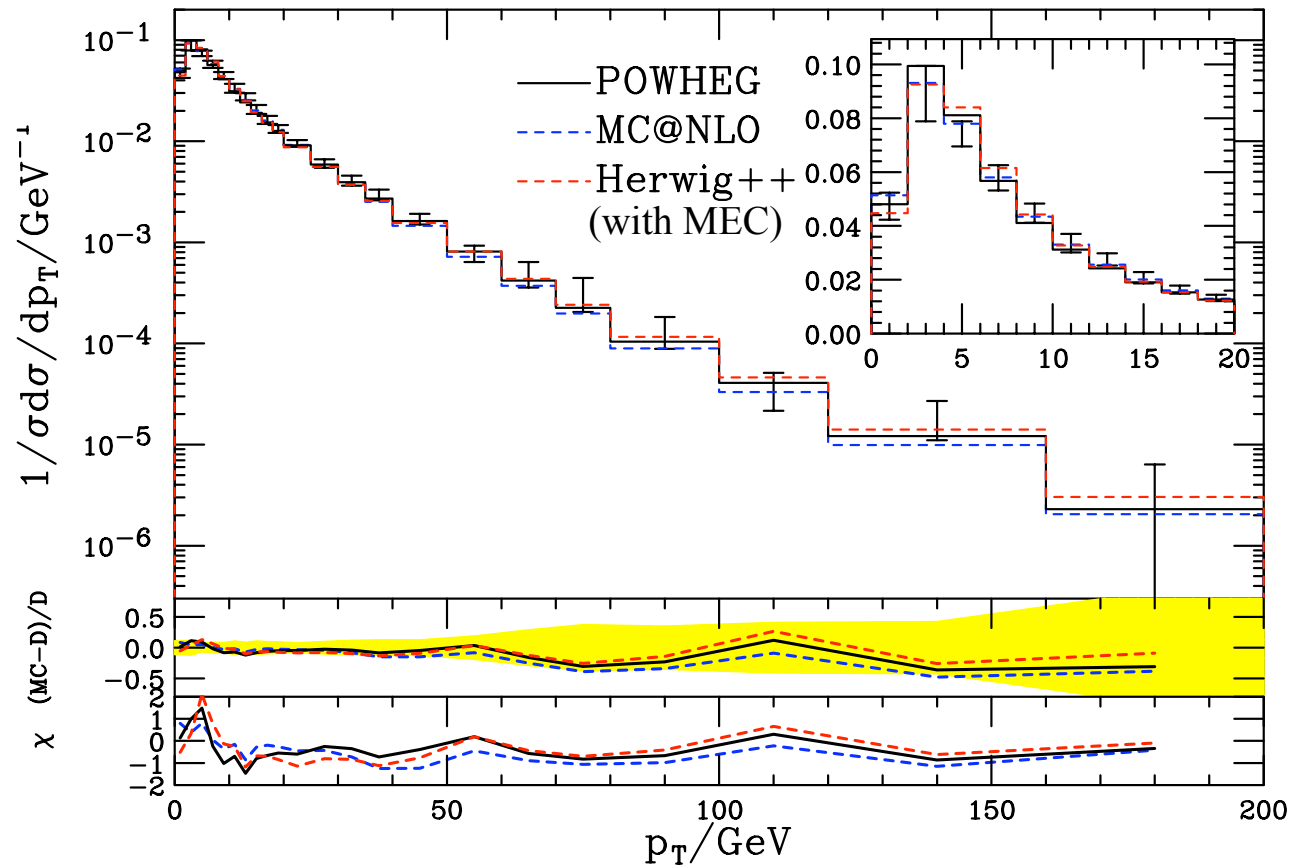


# A high-mass dijet event

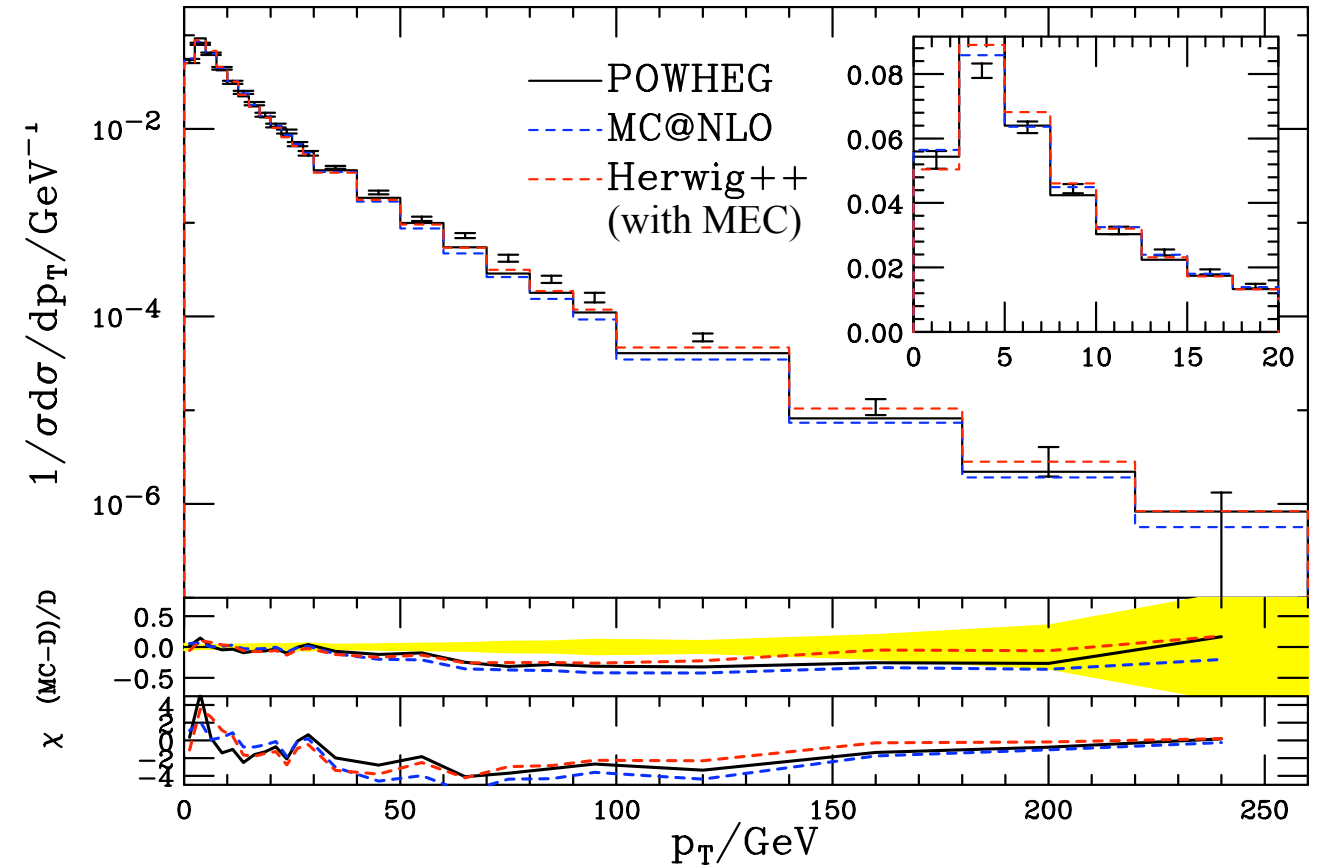


# W & Z<sup>0</sup> at Tevatron

D0 Run I: W



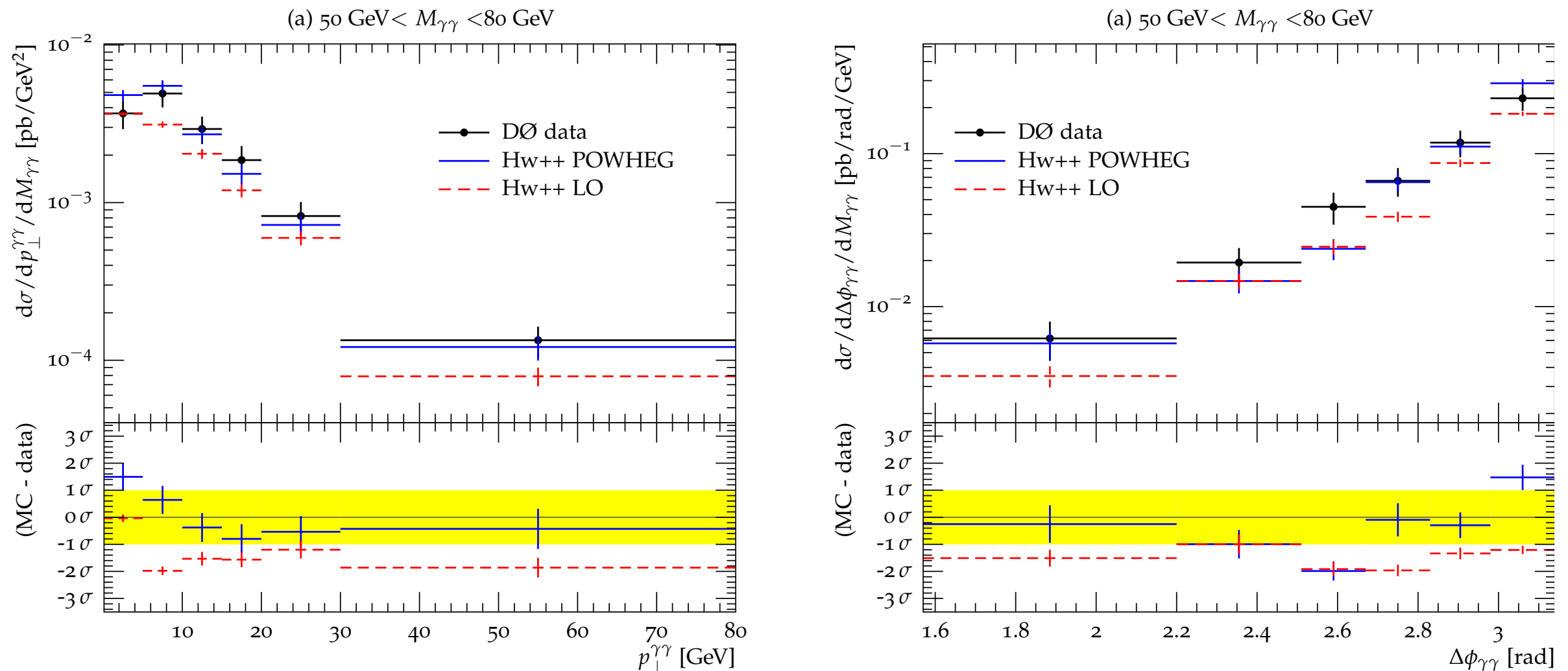
D0 Run II: Z<sup>0</sup>



- Herwig++ includes W/Z+jet (MEC)
- All agree (tuned) at Tevatron
- Normalized to data

Hamilton, Richardson, Tully JHEP10(2008)015

# $\gamma\gamma$ at Tevatron



- Absolute normalization  $\rightarrow$  LO too low
- POWHEG agrees with rate and distribution
- At LHC, important background for Higgs search

D'Errico & Richardson, JHEP02(2012)130

# To Be Confirmed

- Spin and parity  $0^+$ : correlations in  $VV^*$  decays
- Production mechanisms:  $gg, VBF, WH, ZH, ttH$
- Self-coupling ( $HH$  production): **difficult at LHC**
- Total width 4.2 MeV: **impossible?**
- Decay fractions:

$b\bar{b}$	56%	$\tau^+\tau^-$	6.2%	$\gamma\gamma$	0.23%
$WW^*$	23%	$ZZ^*$	2.9%	$\gamma Z$	0.16%
$gg$	8.5%	$c\bar{c}$	2.8%	$\mu^+\mu^-$	0.02%

# Achievable Precision?

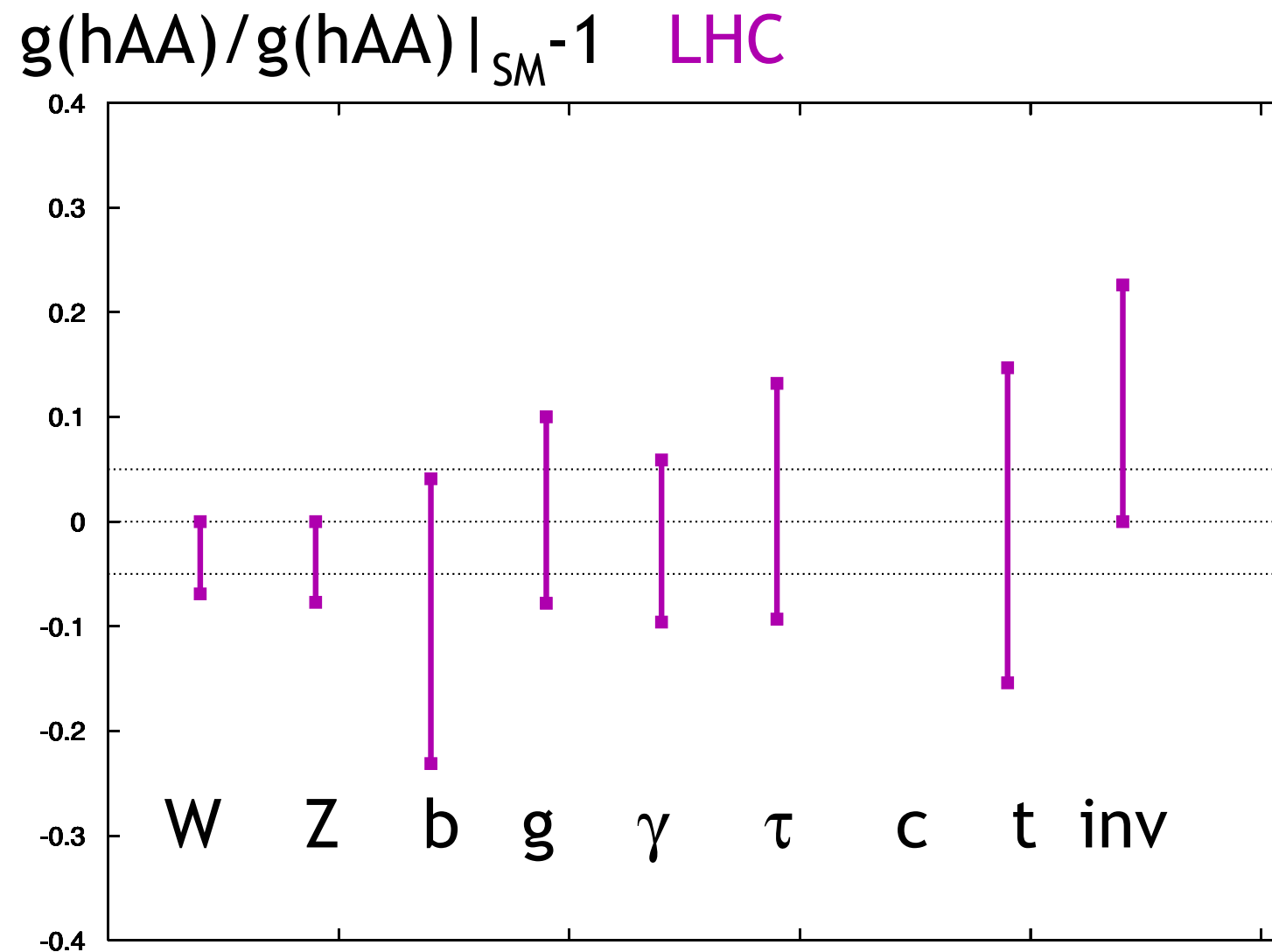


Figure 1: Capabilities of LHC for model-independent measurements of Higgs boson couplings. The plot shows  $1\sigma$  confidence intervals for LHC at 14 TeV with  $300\text{ fb}^{-1}$ . No error is estimated for  $g(hcc)$ . The marked horizontal band represents a 5% deviation from the Standard Model prediction for the coupling.

M Peskin, arXiv:1207.2516

# Achievable Precision?

$g(hAA)/g(hAA)|_{SM}^{-1}$  LHC / ILC1 / ILC / ILCTeV

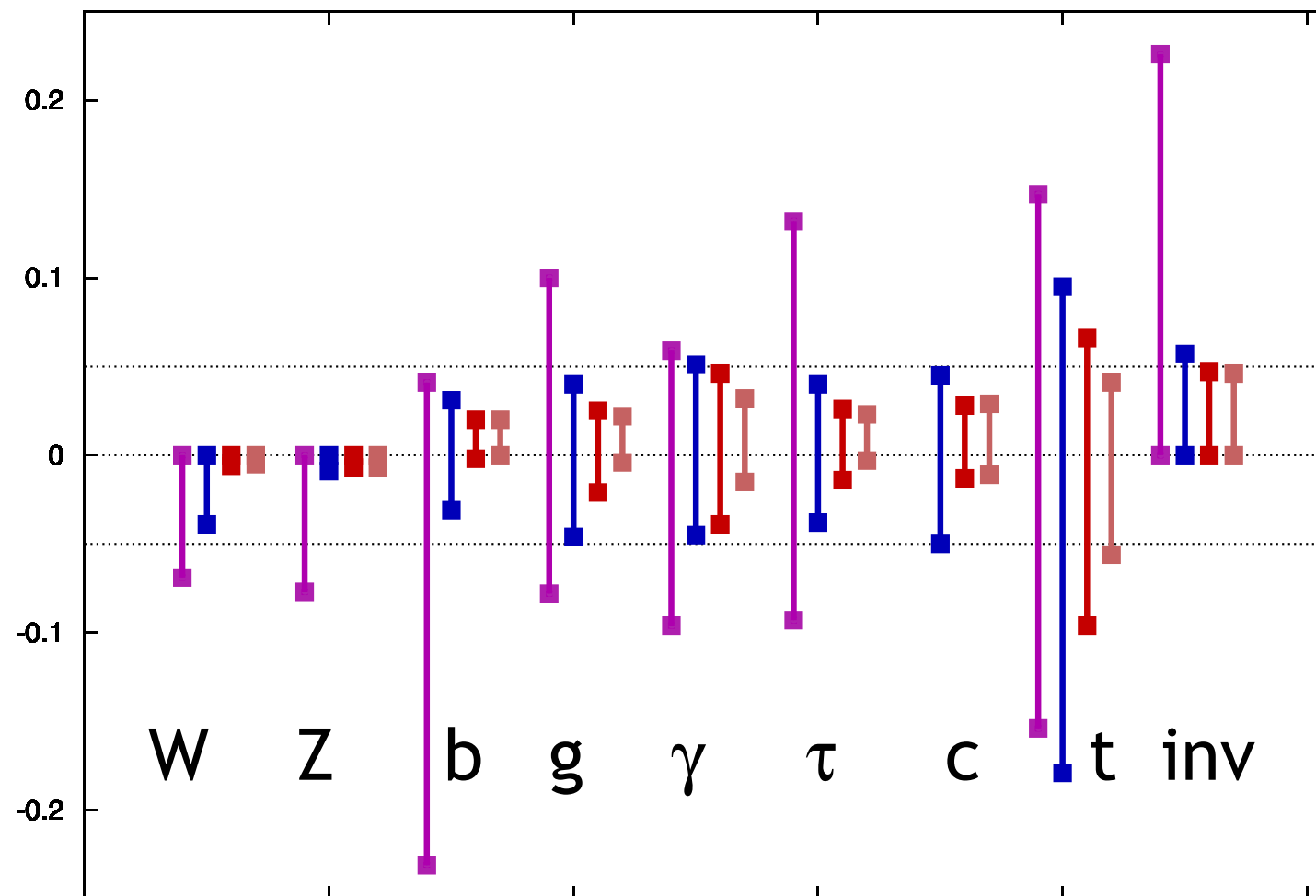


Figure 2: Comparison of the capabilities of LHC and ILC for model-independent measurements of Higgs boson couplings. The plot shows (from left to right in each set of error bars) 1  $\sigma$  confidence intervals for LHC at 14 TeV with 300 fb $^{-1}$ , for ILC at 250 GeV and 250 fb $^{-1}$  ('ILC1'), for the full ILC program up to 500 GeV with 500 fb $^{-1}$  ('ILC'), and for a program with 1000 fb $^{-1}$  for an upgraded ILC at 1 TeV ('ILCTeV'). The marked horizontal band represents a 5% deviation from the Standard Model prediction for the coupling.

M Peskin, arXiv:1207.2516